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#### Abstract

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## IIT RESEARCH INSTITUTE

Technology Center
Chicago 16, Illinois

DASA-1362
DEBRIS HAZARDS, A FUNDAMENTAL STUDY
by
Edward B. Ahlers
for
Defense Atomic Support Agency Washington 25, D. C.

Contract No. DA-49-146-XZ-097
IITRI Project No. 8231


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## FOREWORD

This is the Final Report on IIT Research Institute Project K231, "Fundamental Study of Debris Hazards," conducted for the Defense Atomic Support Agency, Washington, D. C., under Contract DA. 49..146-XZ -097 . All work done under the initial contract on the basic study of debris hazards which provided for supplemental research to analyze data and develop expressions to facilitate the formulation of debris damage predictive schemes for several specific situations done under Modification No. 1, are reported.

Work performed on "DANNY BOY TASK II .. Debris Investigations," under Contract Modification No. 2, was described in the prelimmary and final test reports, "Throwout Study of an Underground Nuclear Detonation," published by the Department of Defense -. U. S. AEC as PORO1814 (ITR-.1814) and POR-1814 (W T-1814).

Work done on "DANNY BOY TASK I - Pre-Shot Predictions," also under Contract Modification No. 2, was reported as a separate report entitled, "Hydrodynamic Analysis for a Buried Underground Nuclear Exposion." Work performed on "DANNY BOY TASK III - Post Test Crater Analysis," under Contract Modification No. 3 will also be submitted as a separate report.

The cooperation and assistance of the Armed Services Explosives Safety Board, (especially Mr. Russel G. Perkins, Chief of Explosives Branch) in allowing project engineers use of their explosions files is greatly appreciated.

IITRI personnel contributing to these studies include E. B. Ahlers, D. I. Feinstein, B. Gain, P. C. Hermann, J. Lakes and Dr. K. E. McKee. Mr. C. A. Miller conducted the analysis of horizontal motion of debris particles under the influence of the blast winds. Mr. R. L. Barnett developed the analysis of the motion of tree debris.

Respectfully submitted,


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#### Abstract

Predicting the mechanical debris, associated with nuclear detonation, stemming from several sources (blast-induced flight of structural fragments, pickup of material from the ground by blast winds, and crater throwout) is a significant problem. Designers of hardened sites are concerned with mater ials which may accumulate atop silo doors or damage vulnerable above-grade antenna systems. Military operations are concerned with the hazards to troops and equipment from the use of nuclear demolition and tactical devices. Potential users of nuclear excavating devices are concerned with hazards to personnel, utilities and equipment.

This refort describes the collection and analysis of data on various aspects of debris formation and dispersion, and examples of data utilization in estimating debris environment in several situations. The approach in the study was to collect extensive data from past experimental and analytical investigations bearing on debris formation and dispersion. By further study the most meaningful formats were summarized to be used as inputs to approximate solutions for debris environment predictions.

An extensive regression study of several hundred HE incidents, accidental and experimental, is made to relate the maximum range of debris to explosion parameters and crater dimensions. Results showing consis.tency with the limited available nuclear data relating to crater throwout are also presented to describe the nature of the debris distribution function in general terms.


Fragmentation data from HE events and laboratory experiments are used to indicate the nature of fragment.size distributions from structural demolition.

An analytical study of the motion of debris fragments caused by blast winds considers debris trajectories for various times of structural failure, fragment sizes, positive and negative phase winds, and initial elevations of the fragment.

Specific estimates are made of the debris hazards to troops of flying tree limbs in the proximity of forest stands, the vulnerability of troop personnel to throwout debris from cratering and stream bed charges, and the debris environment about hardened antenna systems.

Useful estimates of debris environment can be made for many tar.. geting situations with data contained in this report. Refinement of data is certainly essential, especially in experimental definition of fragmentation patterns of ideal structural elements, and in definition of crater lip contours near their extremities, i.e., the throwout debris within and beyond the extremities of the lip.

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## CHAPTER ONE

## INTRODUCTION

While most nuclear weapons effects -- radiation, thermal, alr blast, ground shock, and radioactive fallout -- have been investigated extensively since 1946, little study has been devoted to the creation and distribution of structural debris. This stems in part from the fact that under earlier concepts of vulnerability, yields of weapons and the accuracy of their delivery, the design hardness levels under consideration were such that structural debris was not a prime hazard. Developments in the yield of weapons, delivery accuracy, and design hardness levels make the effect of flying debris a serious consideration. With the utilization of tactical nuclear weapons by field troops, knowledge of debris environment is an important consideration in the deployment of troops. Design specifications for survival of hardened retaliatory weapons sites shortly beyond the edge of the plastic zone of the crater require a knowledge of the deposition of crater throwout material. The siting of communications systems presents particularly serious problems since they are especially vulnerable where substantal flying debris arises.

The investigation was initiated to conduct a series of analytical and experimental studies of the behavior of debris. The following tasks were included:

> Fundamental Studies of Debris Behavior
> Measures of the maximum fragment distance were developed from HE data and compared with nuclear results. The shape of debris distribution functions was studied in detall. Experimental and analytical work on fragment size-distribution was revised. An analytical model for the transport of debris under blast wind loading was developed.

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- Specific Vulnerability Studies

Several vulnerability studies were included, both as applications of the collected data, and to provide measures of vulnerability to flying debris under targeting situations of considerable interest and significance. These studies included:

- Vulnerability of field troops from branch wood within or near forest stands.
- Vulnerability of field troops to crater-emanated debris from very-low yield cratering charges.
- Vulnerability of field troops to crater-emanated debris from very-low yield stream bed charges.
- Vulnerability of antenna systems crater-emanated debris from high-yield nuclear bursts.
- Crater Throwout Study of an Underground Nuclear Detonation

This task was conducted as DANNY BOY Project 1.5 with findings reported under separate cover (Ref. 1, 2). DANNY BOY Project 1.5 involved the following activities:

- Compilation of data relating the initial and terminal positions of a series of more than 1100 "ideal" objects (steel plates, spheres, cylinders and cubes; wood cubes and boards; and common brick) emplaced on the ground surface and in drill holes in the crater zone prior to the shot.
- Compilation of data on the distribution of natural throwout debris beyond the limit of the crater lip, i. e., beyond the ground range where the ground surface is completely obscured by debris.
- Compilation of analysis and plotting of the data in various manners to describe the behavior of crater throwout for this deep underground burst.


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### 1.1 Report Organization

This report is organized into seven chapters. Each of the six chapters following the Introduction concerns a pertinent aspect of the overall debris problem. Thus, the report is a series of related, but selfcontained, studies rather than a continuous-reading document. The contents of the report are summarized by chapters as follows.

## Chapter One, Introduction

The first chapter provides a general introduction to the debris problem. The major conclusions drawn from the investigation are presented. Finally, recommendations for additional research are made.

## Chapter Two, Debris Characteristics of High Explosive

 and Nuclear DetonationsPrevious nuclear weapon effects tests have produced little data concerning the formation and distribution of structural debris, with the following exceptions:

Project 4. 5 of Operation JANGLE included measurements of the debris from airport-type runways and reinforced-concrete wall panels erected over the crater zone of an underground burst.

Project 33. 2 of Operation PLUMBBOB studied the behavior of debris (including window glass, military debris, gravel, stones, and spheres) in response to air blast at various ground ranges. Findings of this investigation had not been published at the time of this investigation.

Other Operation PLUMBBOB studies were concerned with fragments of biological interest, such as glass splinters capable of penetrating abdominal walls.

DANNY BOY Project 1.5, a study of crater
throwout, was recently completed.
SEDAN, a recent nuclear event, included an exper1mental investigation of crater throwout.

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While debris information from past nuclear tests provided only limited data for analysis, extensive measurements were found to be available in reports of planned and accidental $H E$ detonations. These data were collected, studied, and plotted in various manners to define general patterns in debris behavior. Where possible, "check points" from nuclear experience are introduced to indicate their consistency with HE results. Hundreds of explosion reports were reviewed in this task. Data from a series of more than 200 selected explosions were used to obtain expressions relating maximum debris distance to equivalent yield, and to an estimate of the explosive impulse. A smaller series of explosions was used to derive expressions relating maximum debris distance to crater dimensions. Maximum debris distance is correlated with explosion parameters and crater dimensions using the method of least squares; scaled and unscaled relationships are developed using both linear and quadratic relationships. In each regression line so obtained, the standard error and correlation coefficient have been computed as a measure of closeness of fit. Figure 1.1, which expresses the quadratic correlation between maximum debris distance and equivalent yield, and Fig. l. 2, which shows the correlation between scaled maximum debris distance and scaled crater volume, are typical of these regression results. Note particularly, the consistency of JANGLE U results with the HE findings. The much lower position of the DANNY BOY results can be explained on the basis of the very deep burial of the device in this event, which resulted in trajectories with pronounced vertical components in the ejecta. Actually, it may be argued that the HE detonations are more like the buried nuclear burst than surface nuclear bursts, because of the absence of substantial blast winds.

To define the probability of personnel or equipment being hit by missiles, and, for the shorter ground ranges, the quantity of material likely to be deposited atop a siio door, it is necessary to estimate the distribution of fragments within the maximum ground range. Several approaches to this problem were pursued. A theoretical model for fragment distribution from ground zero to maximum debris distance was derived by assuming wall and roof panels to fragment into equisized fragments upon detonation of a nearby line charge. Comparison of this model with several explosions shows

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Figure 1.1 Quadratic Regression Line: Maximum Debris Distance versus Equivalent Yield

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Figure $1.2 \frac{\text { Linear Regression Line: } W^{1 / 3}-\text { Scaled Maximum }}{\text { Debris Distance Vs. } W^{1 / 3} \text {-Scaled Crater Volume }}$
that the shape of the curves is similar (Fig. 1.3). Next, debris distribution patterns of a series of six ordnance explosions were plotted in various matters to note their similarity. Thirdly, debris distribution of an ordnance structure (involving over 30,000 recorded fragments with a total weight of about 43 tons) was plotted in detail. Contrary to expectation based on drag effects, the large fragments from this explosion did not travel as far as the s:ialler ones (Fig. 1.4). This is probably because fragments subject to forces sufficient to cause large acceleration were also subject to a greater degree of breacup.


Figure 1.3 Comparison of Theoretical and Actual Fraginent Dispersion Patterns

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## Chapter Three, Fragmentation Experimental Observations

Since energy levels of the flying fragments are a measure of their potential to penetrate shields, impart shock loading to equipment, or incapacitate personnel, it is well to define the expected size-distribution of fragments. No definitive experimental investigation of structural fragmentation has yet been made, even on idealized structural elements, by which the fragment-size distribution is related to structural strength and loading parameters. For this reason an attempt was made to collect and summarize past experimental work on fragmentation of materials and structures, and to describe the general fragment-size distribution patterns that have been observed. Five such investigations are described;

The British Coal Utilization Research Association studies on coal breakage from random forces,
Safety in Mines Research Establishment (of Great Britain) research on explosively detonated stone blocks.

Stanford Research Institute model tests on fragmentation of reactor containment structures from internal explosions,
The Pantex Ordnance Plant detailed fragment counts from the planned explosion of a reinforced concrete ordnance structure,

Project 4 . 5 of Operation JANGiLE studies of fragment size distributions from reinforced concrete wall panels erected over the crater zone.

Experimentation has shown that the higher the loading on the source material, the smaller the fragments produced, and that a wide range of fragment sizes are produced by any loading. An "Ideal Law of Breakage" which shows excellent fit to the experimental data has evovled from coal-crushing investigations. The extensive data from over 30,000 concrete fragments in the Pantex Ordnance Plant event allowed a detailed study of the fragmentsize distribution. It is interesting to note that only about 3 percent of the fragments recorded in this event (above l-ounce in size) weighed more than three pounds, but that these accounted for nearly 75 percent of the total weight of all fragments. This tends to support the hypothesis that for many problems involving impact of fragments, there may be an optimum

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fragment-size for purposes of design predictions. Figures 1.5 and 1.6 show these distributions. The only major structural fragmentation study conducted on nuclear detonations was the test of wall panels on the JANGLE $U$ event. Size distributions of the larger fractions were plotted as part of that project and it was noted that the JANGLE data did not preclude the possibility that the concrete fragment-size distribution caused by the underground nuclear shot followed the same pattern as mined coal or ore in a crusher.


Figure 1.5 Cumulative Fragment-Size Distribution for a Reinforced Concrete Ordnance Structure

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Equivalent Diameter, (in.)

 io: a Reantocod Conc: eie O-drance S Uti:

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## Chapter Four, Fragmentation, Analytical Considerations

Prediction of the number and size into which a structural component fragments, when subjected to blast loading, is totally beyond the current state-of-the-art of fracture mechanics. Thus, a relatively simple mathematical model was formulated to be used to predict debris formation, which can be adapted to include advances in fracture mechanics. The model selected for this study, is a simple extension of the single-degree-of-freedom structural model long used to analyze the elastic-plastic response of structural components subjected to blast loading. The response of the single-degree-of-freedom system with an elastic, perfectly plastic spring is determined. In particular the mass, velocity, and the time when the mass reaches a certain displacement, called the fracture displacement, are of interest. The magnitude of the fracture displacement and the description of particles formed when this displacement is reached must await further analytical and experimental developments in fracture mechanics.

## Chapter Five, Debris Transport by Blast Winds

The motion of a particle acted on by the nuclear blast winds is analyzed on the assumption that initial conditions of the particle motion are known. The model used assumes that the force acting on the particle is proportional to the square of the relative velocity between the particle and air. Blast parameters are assumed constant over the range of travel of the debris, and it is further assumed that the apparent lengthening of the positive phase duration due to the debris motion in the direction of shock propagation can be handled by a simple adjustment of positive phase duration. These assumptions are necessary to reduce the equation of motion for any debris particle to a one-parameter nonlinear differential equation. Without making these two assumptions, it would be necessary to treat each weapon yield and placement as a separate problem and no general observation could be made regarding debris behavior. The equation of motion is numerically integrated and results are obtained for a wide range of overpressures and particle sizes. It is possible to use the results directly to determine distance-time plots for any debris. The effect of negative

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phase winds on particle motion is also studied. While a complete description of the negative phase wind is not avallable, reasonable estimates of this wind can drastically change the motion of the debris particle.

## Chapter Six, Vulnerability of Field Troops to Tree Debris

The study of the vulnerability of engineer and field troops to hazards from tree debris is one of three specific debris problems of interest to the Office of the Chief of Engineers, U.S. Army. The objective of this study was to define a "safe distance" for positioning troops to avoid casualties from falling trees and limbs. The safe distance may actually fall inside or outside the forest and both cases are represented. Since any individual tree may fail at the base under blast loading; tree height is considered the minimum safe distance in any situation. Making certan simplifying assumptions (i.e., zero strength tree limts, plane blast wave loading, unobstructed trajectories, and that personnel struck by tree limbs are certain casualties), tree limbs are followed in their trajectories from the time of shock arrival to their impact with the ground. Results show that for the lower yields (l KT, for example) trajectories become vertical early. A uniform translation of all branches is thus obtained, the area in front of the forest up to the "safe distance" being similar in appearance to that of the forest floor after all branchwood was allowed to drop vertically. For the higher yield weapons ( 20 MT , for example) trajectories terminate before they become vertical. Results are similar with the exception that the highest branches of the first few rows of trees pile up in a lower density than those following closer-in trajectories. Results of this study are summarized in the following table: which lists safe distances in terms of tree height, weapon yield, and overpressure levels.

## Chapter Seven, Vulnerability of Field Troops to Throw Out Debris from Cratering and Stream Bed Charges

Methods for estimating safe distances for positioning troops in the proximity of very-low-yield cratering and stream-bed charges, based on debris criteria, were studied. As with the preceding tree debris vulner-

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SAFE DISTANCES TO PREVENT CASUALTIES FROM TREE DEBRIS

| Yield of Weapon | $\begin{gathered} \text { Height } \\ \text { of } \\ \text { Tree. } \end{gathered}$ | Safe Distance Falls Inside of Forest <br> Safe Distance from Ground Zero | Safe Distance F <br> Safe Distance from Ground Zero for Various Overpressures at the Front of the Forest, (yd) |  |  |  |  | Outside of Forest Safe Distance from Forest $^{\text {Ones }}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  | Safe Distance from Forest for Various Overpressures at the Front of the Forest, (yd) |  |  |  |  |
|  | feet | 1.7-2.2 psi | 5 psi | 10 psi | 15 psi | 20 psi | 30 psi | 5 psi | 10 psi | 15 psi | 20 psi | 30 psi |
| 0.05 KT | 20 | 311-376 yards | 193 | 130 | 109 | 98 | 87 | 7 | 8 | 10 | 12 | 16 |
|  | 40 |  | 200 | 136 | 113 | 100 | 87 | 14 | 14 | 14 | 14 | 1s |
|  | 60 |  | 206 | 142 | 119 | 106 | 91 | 20 | 20 | 20 | 20 | 20 |
|  | 80 |  | 213 | 149 | 126 | 113 | 98 | 27 | 27 | 27 | 27 | 27 |
|  | 100 |  | 220 | 156 | 133 | 120 | 105 | 34 | 34 | 34 | 34 | 34 |
|  | 120 |  | 226 | 162 | 139 | 126 | 1:1 | 40 | 40 | 40 | 40 | 40 |
|  | 160 |  | 239 | 175 | 152 | 139 | 124 | 53 | 53 | 53 | 53 | 53 |
|  | 200 |  | 253 | 189 | 166 | 153 | 138 | 67 | 67 | 67 | 67 | 67 |
| 0.1 KT | 20 | 392-473 yards | 241 | 164 | 137 | 123 | 108 | 7 | 10 | 12 | 15 | 19 |
|  | 40 |  | 248 | 168 | 139 | 123 | 108 | 14 | 14 | 14 | 15 | 19 |
|  | 60 |  | 254 | 174 | 145 | 128 | 109 | 20 | 20 | 20 | 20 | 20 |
|  | 80 |  | 261 | 181 | 152 | 135 | 116 | 27 | 27 | 27 | 27 | 27 |
|  | 100 |  | 268 | 188 | 159 | 142 | 123 | 34 | 34 | 34 | 34 | 34 |
|  | 120 |  | 274 | 194 | 165 | 148 | 129 | 40 | 40 | 40 | 40 | 40 |
|  | 160 |  | 287 | 207 | 178 | 161 | 142 | 53 | 53 | 53 | 53 | 53 |
|  | 200 |  | 301 | 221 | 192 | 175 | 156 | 67 | 67 | 67 | 67 | 67 |
| 0.5 KT | 20 | 670-809 yards | 410 | 279 | 235 | 210 | 186 | 10 | 16 | 21 | 26 | 33 |
|  | 40 |  | 414 | 279 | 235 | 210 | 186 | 14 | 16 | 21 | 26 | 33 |
|  | 60 |  | 420 | 283 | 235 | 210 | 186 | 20 | 20 | 21 | 26 | 33 |
|  | 80 |  | 427 | 290 | 241 | 211 | 186 | 27 | 27 | 27 | 27 | 33 |
|  | 100 |  | 434 | 297 | 248 | 218 | 187 | 34 | 34 | 34 | 34 | 34 |
|  | 120 |  | 440 | 303 | 25tr | 224 | 193 | 40 | 40 | 40 | 40 | 40 |
|  | 160 |  | 453 | 316 | 267 | 237 | 206 | 53 | 53 | 53 | 53 | 53 |
|  | 200 |  | 467 | 330 | 28.1 | 251 | 220 | 67 | 67 | 67 | 67 | 67 |
| 1 KT |  | 845 - 1021 yards |  |  |  |  |  |  |  |  |  |  |
|  | 40 |  | 519 | 351 | 295 | 264 | 233 | 14 | 20 | 26 | 32 | 41 |
|  | 60 |  | 525 | 351 | 295 | 264 | 233 | 20 | 20 | 26 | 32 | 41 |
|  | 80 |  | 532 | 358 | 296 | 264 | 233 | 27 | 27 | 27 | 32 | 11 |
|  | 100 |  | 539 | 365 | 303 | 266 | 233 | 34 | 34 | 34 | 34 | 41 |
|  | 120 |  | 545 | 371 | 309 | 272 | 233 | 40 | 40 | 40 | 40 | 41 |
|  | 160 |  | 559 | 385 | 323 | 286 | 246 | 54 | 54 | 54 | 54 | 54 |
|  | 200 |  | 572 | 398 | 336 | 299 | 259 | 67 | 67 | 67 | 67 | 67 |
| 1 MT | 20 | $8,448=10,208$ yards | 5,119 | 3,430 | 2,856 | 2,526 | 2, 182 | 68 | 121 | 165 | 203 | 265 |
|  | 40 |  | 5,137 | 3,460 | 2,895 | 2,574 | 2, 243 | 86 | 151 | 204 | 251 | 326 |
|  | 60 |  | 5,147 | 3,475 | 2,916 | 2,599 | 2, 273 | 96 | 166 | 225 | 276 | 356 |
|  | 80 |  | 5,153 | 3,485 | 2.927 | 2,612 | 2,290 | 102 | 176 | 236 | 289 | 373 |
|  | 100 |  | 5, 157 | 3,490 | 2,934 | 2,620 | 2,299 | 106 | 181 | 243 | 297 | 382 |
|  | 120 |  | 5, 160 | 3,493 | 2,937 | 2,623 | 2,302 | 109 | 184 | 246 | 300 | 385 |
|  | 160 |  | 5, 163 | 3,495 | 2,938 | 2, 623 | 2,302 | 112 | 186 | 247 | 300 | 385 |
|  | 200 |  | 5,163 | 3,495 | 2,938 | 2,623 | 2, 302 | 112 | 186 | 247 | 300 | 385 |
| 20 MT | 20 | 22,932-27,709 yards | 13,792 | 9,128 | 7,505 | 6,553 | 5,528 | 81 | 146 | 200 | 247 | 325 |
|  | 40 |  | 13,821 | 9.179 | 7.575 | 6,639 | 5,640 | 110 | 197 | 270 | 333 | 437 |
|  | 60 |  | 13,842 | 9,215 | 7.624 | 6,699 | 5,718 | 131 | 233 | 319 | 393 | 515 |
|  | 80 |  | 13,858 | 9,243 | 7,662 | 6,746 | 5,779 | 147 | 261 | 357 | 440 | 576 |
|  | 100 |  | 13,871 | 9,267 | 7.693 | 6, 786 | 5,829 | 160 | 285 | 388 | 480 | 626 |
|  | 120 |  | 13,882 | 9,287 | 7,720 | 6,818 | 5,872 | 171 | 305 | 415 | 512 | 669 |
|  | 160 |  | 13,901 | 9, 304 | 7,764 | 6,872 | 5,940 | 190 | 322 | 459 | 566 | 737 |
|  | 200 |  | 13,916 | 9,345 | 7,799 | 6,914 | 5,994 | 205 | 363 | 494 | 608 | 791 |

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ability problem, these two cases were studied to supply data needed by the Office of the Chief of Engineers, U. S. Army. Numerous reports on debris from HE cratering tests were studied to obtain measures of the relationship between debris distance, weapon yield, depth of burst, fragment size, and soil characteristics. The U. S. Geological Survey report on a series of HE cratering tests in basalt (Area 18 at Nevada Test Site) provided data on relationships between the first four of these factors (Ref. 24). Using these data, nomographic procedures were presented for estimating average debris distance for various fragment sizes in basalt, for various combinations of yield and depth of burst. An asymmetry factor is introduced to obtain maximum rather than average debris distances, based on measurements of the ray-like patterns in the USGS study. Using data from the Panama Canal series of cratering tests in various media (Ref. 29), distance ratios are plotted obtaining estimates of debris distances in media other than basalt, and for the stream-bed charge. Debris distances found by this method, using DANNY BOY explosion parameters, are consistent with photographic observations of the debris deposit from this event.

## 1. 2 Summary and Conclusions

Various aspects of the debris problem -- fragmentation characteristics of materials and structures, transport of the debris particles by the blast winds, and the ultimate distribution of material resulting from these factors -- are studied in this investigation. Both analytical and experimental studies of the phenomena are considered. In some cases, behavior is defined empirically from collected historical data, in other cases, by analytically derived expressions and exhibits.

Fundamental data are collected to provide a basis for initial estimates of the nature of the debris problem in specific targeting situations. Historical data (HE detonations) are used to derive empirical expressions correlating maximum debris distance with equivalent TNT yields, impulse, and crater dimensions. These serve to make initial approximations of the limiting ground range of debris problems. Data compiled on relative distribution of debris particles can be used to make intial estimates of fragment quantities at various ground ranges. Estimates made from these

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data are, at best, crude. For the most part they are based on HE data and nuclear - HE equivalence. Use of the latter is questionable -- expecially since extreme differences in blast winds exist between nuclear and HE. As with most HE results, the experimental findings include great degrees of scatter.

The existence of definite debris patterns is graphically demonstrated in this report. The only pehomenon lending itself to analytical treatment is the transport of debris particles by the blast winds. A relatively complete treatment of the phenomenon is made. Fragmentation of materials and structures is reviewed in detail -- but this study simply cannot be carried very far without considerable experimentation.

The following conclusions are drawn from this study:
(1) Data in this report can be used to determine debris distribution under a variety of targeting situations.
(2) In many situations initial estimates of the severity of the debris problem can be made from the data compiled.
(3) Progress has been made in providing data and methods for estimating the existence and severity of debris problems, but the approximations are crude.
(4) The various debris phenomena (fragment-size distributions for materials and structures, debris dispersion patterns, and variations with explosion parameters) follow characteristic patterns which can be established experimentally.
(5) The feasibility of predicting debris effects of nuclear explosions has been demonstrated by examples. Bercause random and/or uncontrollable (and unpredictable) factors influence debris behavior, prediction of these effects will never have a high degree of precision.
(6) Debris prediction methods can be greatly improved by experimental data which can be developed on future fullscale nuclear test programs.

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Problems exist which cannot be solved without further investigation, primarily experimental. Knowledge of fragmentation of actual structures, or even structural elements, is sub-minimal. Little is known of fragment quantities and the fragment-size distribution which can be expected from structural elements, and even less is known of the variations in the characteristics which would be caused by different levels of impulsive loading. Likewise, the time of failure of structural elements is a completely unknown factor. The analytical treatment of fragment transport by blast winds, included in this report, requires experimental verification. The problem of crater throwout debris and the resultant buildup of the crater lip is significant in terms of the mass of material likely to be accumulated atop hardened sites. The work on crater throwout debris begun on Project DANNY BOY under this contract should also be expended to cover variations in parameters.

## 1. 3 Recommendations

On the basis of the findings of work performed under this contract, it is recommended that further analytical and experimental investigation be made in the following areas.
(1) Further study of the data accumulated in this report, aimed toward codifying estimating procedures.
(2) Further analysis of the accumulated data from the crater throw out debris study from the DANNY BOY event.
(3) Inclusion of crater throwout debris studies in possible forthcoming tests, to study effects of variations in parameters -- other yields, depths of burial, and soils.
(4) Experimental investigation of fragmentation of phototype wall and roof panels -- studying time of failure, fragment-size distribution, fragment dispersion, and effects of variations in the loading pulse. This work could be performed under a combination of HE and full-scale nuclear testing.

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## CHAPTER TWO <br> DEBRIS CHARACTERISTICS OF HIGH EXPLOSIVE AND NUCLEAR DETONATIONS

Little data on formation and dispersion of structural debris have been collected from full-scale nuclear tests. Since debris hazards were not initially considered a major problem, tests of structures and structural elements did not include measurements of fragmentation or debris transport. One exception was Project 4. 5 of Operation JANGLE, in which reinforced concrete wall and runway panels were erected above the expected crater zone and debris measurements taken (Ref. 5). The vulnerability of parked aircraft was emphasized in this study, thus major interest was centered about the maximum ground range of aircraft damage from blast and from debris. A second exception was Project 33. 2 of Operation PLUMBBOB in which the behavior of missiles (including window glass, military debris, gravel, spheres and native stone) emplaced at various ground ranges was studied. Findings of Project 33. 2 Operation PLUMBBOB had not been published at the time of this investigation. Other debris studies on Operation PLUMBBOB involved such fragments as small glass fragments capable of penetrating abdominal walls, which are of biological interest but which cannot be regarded as structural debris. More recently, crater throwout debris studies included in underground nuclear test programs on DANNY BOY and SEDAN have investigated objects emplaced on the ground surface and within the expected crater, and their post-shot locations relative to original positions. These crater throwout studies involved deep-buried shots in which the debris was not influenced substantially by blast winds, as it would be with a surface or shallow-buried shot.

From the beginning of this investigation, it was apparent that fragmentation and debris dispersion data from past full-scale nuclear tests would provide only limited information for analysis. It became obvious, however, that an extensive body of data was available in explosion reports from HE events -- both planned and accidental -- and that these could be studied to define general debris behavior patterns. The limited nuclear data could

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then be introduced as "check points" on the HE studies. Several hundred selected reports on HE detonations were studied to determine the relationship between the maximum debris distance and various explosion parameters; these were then used to define the expected limit of the debris hazard. The least-squares method was used to correlate maximum debris distance with equivalent yield, computed values of impulse, and crater dimensions. It is of particular interest to note that comparisons of maximum debris distance from the JANGLE and DANNY BOY studies are consistent with the HE results.

In addition to defining the maximum debris distance, it was necessary to estimate the distribution of debris fragments at distances shorter than maximum. Quantitative data are needed for any estimate of probability of equipment or personnel being hit by flying fragments in the debris zone. Several approaches to this problem were taken. A theoretical model for the distribution of fragments from ground zero to maximum debris distance was developed by assuming structural wall and roof panels to fail into equi-sized fragments, with radial initial velocity vectors, upon detonation of a line-charge placed near them. Comparison with debris-distribution patterns of several explosions show that curves are of similar shapes. Next, the debris-distribution patterns of a series of six HE detonations involving ordnance structures were plotted together in various manners to note their similarity. Thirdly, debris distribution from a reinforced concrete ordnance structure was plotted in various ways. The availability of data, from this event, on the weights and terminal locations of over 30,000 fragments with a total weight of about 43 tons permitted highly detailed treatment of debris distribution. An interesting observation of this study is that the large fragments did not travel as far as the smaller fragments. This seemingly contradicts expectations based on air drag effects, but the comparison is inappropriate since the particles subject to forces sufficient to cause large motions are also subject to a greater degree of fragmentation.

## 2. 1 Application of HE Data from Planned and Accidental Explosions

Past HE detonations, both planned and accidental, are an extensive source of data on the behavior of structural debris. Explosion reports customarily include some or all of the following information.

> Total weight of explosive involved Weight of explosive exploding at one time
> Kind of explosive
> Source structure
> Crater dimensions
> Maximum debris distance
> Distances at which individual fragments were found Distances at which various degrees of structural damage occurred

Several detailed compilations of explosion reports have been made (Ref. 6, 7, 8). Explosion reports make it possible to plot debris distances -- frequently only the maximums -- against various explosion parameters. Statistical measures of the consistency of these data can also be made.

Although the validity of debris data from individual explosion reports is limited in certain respects, the availability of many reports does permit averaging the data and constructing graphical and statistical measures of debris beharior. Plots of maximum debris distance in terms of equivalent explosive weight which permit estimates of expected debris ranges are included here. While scaling from HE to nuclear is generally a questionable procedure, these plots can be, and in fact have been, applied in determining whether or not debris problems are likely to exist at various ground ranges.

Certain limitations in the use of data taken mostly from reports on accidental explosions are apparent. Extreme variations exist in the range of explosion environments. Explosion source structures include steamers, freight cars, munitions plants, warehouses and igloos. Some of the explosions occurred in the open field. The relative degree of confinement afforded by these structures certainly affects the loading on the flying fragments. The amount of material available for fragmentation is a function of the source structure and, for larger explosions, of the surrounding neighborhood as well. In some cases, explosion data may be questioned as to the accuracy of the weight of explosives involved in accidental explosions, whether high or low order detonation took place, and the quanity of explosives involved in individual detonations of a series of consecutive explosions. Also, the

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trajectories of fragments from HE detonations will range from the horizontal to the vertical, the farthest-flung may have initial trajectory angles of around $45^{\circ}$ from the horizontal. This behavior is unlike fragment trajectories resulting from a structure failing under loading from the plane blast wave of a nuclear explosion.

## 2. 2 Measures of Maximum Debris Distance

A series of 206 accidental and planned explosions was selected from those tabulated in references 6 and 8 and subjected to a regression study to develop estimating expressions for maximum debris distance -- a measure of maximum range of vulnerability to debris hazards. Selected explosions are listed in Appendix A. which tabulates maximum debris distances and the major explosion parameters. Maximum debris distance was correlated with equivalent weight of explosives and impulse; various scaled relationships between these functions were also studied. A total of 40 linear and 40 quadratic regression lines was determined by computer methods, together with standard errors and correlation coefficients for each line. (Results of this correlation study are included in Table 2.1 and Fig. 2. 3 through 2. 10.)

## 2. 2.1 Details of Regression Analysis

The plots of maximum debris distancu und explosion parameters, (Fig. 2. 3 through 2.10) show a great degree of scatter among data points. It is quite apparent that this dispersion is, in all cases,far too great to permit visual location of any average line through the data points. Statistical devices must be used to describe the plotted relationships satisfactorily. Simple (linear) and quadratic correlations were therefore made, using least-squares methods.

It should be noted that the computation of a regression line does not necessarily assure a functional relationship between the factors correlated; in this case it merely describes what has historically occurred. Thus, the computation and drawing of a regression line describing the correlation between maximum debris distance and equivalent weight of explosives does not inean that the two factors are functionally related by some physical law to the exclusion of such other considerations as configuration and strength

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of structures and impulse distance. If, however, the association between maximum debris distance and some explosion parameter were found to be sufficiently close, it could be possible to estimate, with some calculable degree of accuracy, the maximum debris distance for other explosions on the basis of the observed relationship found in this sample of 206 explosions.

The computed regression lines were derived by the method of least squares -- one of the most widely used methods of curve fitting - which yields the equation of the line from which the summation of the squares of the deviations of all the plotted points is a minimum. Application of the least-squares linear regression line is demonstrated in Fig. 2. 1.

Derivations of equations for least-squares regression lines can be found in most engineering statistics texts. The linear least squares line of the form

$$
x_{12}=a_{12}+b_{12} x_{2}
$$

is obtained by solving the simultaneous equations

$$
\begin{aligned}
& \mathrm{Na}_{12}+\mathrm{b}_{12} \Sigma \mathrm{X}_{2}=\Sigma \mathrm{X}_{1} \\
& \mathrm{a}_{12} \Sigma \mathrm{X}_{2}+\mathrm{b}_{12} \Sigma \mathrm{X}_{2}^{2}=\Sigma \mathrm{X}_{1} \mathrm{X}_{2}
\end{aligned}
$$

to obtain the coefficients $\mathrm{a}_{12}$ and $\mathrm{b}_{12}$.
Similarly, the quadratic least-squares regression line of the form

$$
x_{12}=a_{12}^{\prime}+b^{\prime} X_{2}+c^{\prime} x_{2}^{2}
$$

is obtained by solving the set of three simultaneous equations:

$$
\begin{align*}
& \mathrm{Na}_{12}^{\prime}+\mathrm{b}_{12}^{\prime} \Sigma \mathrm{X}_{2}+\mathrm{c}_{12}^{\prime} \Sigma \mathrm{X}_{2}^{2}=\Sigma \mathrm{X}_{1} \\
& \mathrm{a}_{12}^{\prime} \Sigma \mathrm{X}_{2}+\mathrm{b}_{12}^{\prime} \Sigma \mathrm{X}_{2}^{2}+\mathrm{c}_{12}^{\prime} \Sigma \mathrm{X}_{2}^{3}=\Sigma \mathrm{x}_{1} \mathrm{X}_{2}  \tag{2,1}\\
& \mathrm{a}_{12}^{\prime} \Sigma \mathrm{X}_{2}^{2}+\mathrm{b}_{12}^{\prime} \Sigma \mathrm{X}_{2}^{3}+\mathrm{c}_{12}^{\prime} \Sigma \mathrm{X}_{2}^{4}=\Sigma \mathrm{x}_{1} \mathrm{X}_{2}^{2}
\end{align*}
$$

The advantages of deriving least-squares lines in this study are that (1) reproducible associative relationships are developed in accordance with accepted statistical procedures, and (2) measures of probable error and closeness of fit can be computed. The standard error is a measure of

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Figure 2.1 Graphical Representation of Least-Squares
Linear Regression Line

$$
\begin{aligned}
\mathbf{X}_{1}, \mathbf{X}_{2} & =\text { a set of bivariate data } \\
\overline{\mathbf{X}}_{1} & =\text { arithmetic mean of } \mathbf{X}_{1} \\
\mathbf{X}_{12} & =\text { least-squares regression line values of } X_{1} \\
\mathbf{x}_{12} & =\text { deviations explained by regression of } X_{1} \text { on } \mathbf{X}_{2} \\
\mathbf{x}_{1,2} & =\mathbf{X}_{1}-X_{12}=\text { residual or unexplained deviations } \\
N & =\text { number of items in sample. }
\end{aligned}
$$

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the distribution of data points about the regression line, and is defined as


The standard error about the least-squares regression line may be interpreted in a manner analogous to the standard deviation of a frequency distribution as represented by the normal curve. It is a measure of the distribution of data points about the least-squares regression line, and its interpretation is demonstrated graphically in Fig. 2. 2. The upper chart in Fig. 2.2 shows the relative frequency of events producing data points within various multiples of the standard error about the regression line if the distribution is normal. About two-thirds of all values would be within $\mathrm{X}_{12} \pm \mathrm{s}_{1.2}$, and all but 27 out 10,000 would fall within $\mathrm{X}_{12} \pm 3 \mathrm{~s} 1.2^{\cdot}$ The lower chart is a more appropriate expression of this distribution for vulnerability analysis. This chart shows the frequency of events occurring below the various multiples of standard error about the regression line. Its use is best shown by example.

If we let $D_{1}$ equal maximum debris distance and $W_{2}$, and $W_{3}$ the various equivalent yields, then the probability of a piece of debris being thrown as far as $D_{1}$ or beyond is 0.14 percent for $W_{1}, 2.9$ percent for $W_{2}$ and 15.9 percent for $W_{3}$. Similar$l y$, for a weapon of yield $W_{3}$ the probability of a piece of debris being thrown a distance equal to or beyond $D_{1}$ is 15.9 percent; and is 2.9 percent for $D_{2}$ and 0.14 percent for $D_{3}$.

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Figure 2.2 Normal Distribution of Events about Least-Squares Regression Line

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Standard error is expressed in units of the dependent variable and, as such, is not amenable to comparison with unlike quantities -- or moreover, with any general standard of closeness-of-fit. The correlation coefficient is a measure of scatter in dimensionless terms, and is defined as the square root of the ratio of the explained variation $\Sigma x_{12}^{2}$ to the total variation $\Sigma \mathbf{x}_{1}^{2}$,

$$
r=\sqrt{\frac{\Sigma \mathbf{x}_{12}^{2}}{\Sigma \mathbf{x}_{1}^{2}}}
$$

The correlation coefficient can be computed for the linear case by the expression

$$
r=\sqrt{\frac{\Sigma x_{12}^{2}-X \Sigma x_{1}}{\Sigma x_{1}^{2}-X \Sigma x_{1}}}
$$

and for the quadratic case by

$$
r=\sqrt{1-\frac{N-1}{N-3} \frac{\left(\Sigma x_{1}^{2}-\frac{\left(\Sigma x_{1}\right)^{2}}{N}\right)-b_{12}^{\prime}\left(\Sigma x_{1} x_{2}-\frac{\Sigma x_{1} \Sigma x_{12}}{N}\right)-c_{12}^{\prime}\left(\Sigma x_{1} x_{2}-\frac{\Sigma x_{1} \Sigma x_{2}^{2}}{N}\right)}{\Sigma x_{1}^{2}-\frac{\left(\Sigma x_{1}\right)^{2}}{N}}} .
$$

The sign of $r$ is positive for a regression line of positive slope, and negative for a regression line with a negative slope.

In general a correlation coefficient of +0.90 or greater indicates high positive correlation, and between zero and +0.10 indicates a low positive correlation. Positive correlation coefficients below +0.90 generally do not engender high degrees of confidence. A correlation coefficient of +0.50 or thereabouts is decidedly marginal.

## 2. 2. 2 Regression Study of Maximum Debris Distance

The quantities of explosives involved in the 206 explosions studied ranged from 8 lb to $9,000,000 \mathrm{lb}$, providing more than seven cycles of data. The total weight of explosive materials involved was about $50,000,000 \mathrm{lb}$. Since the largest explosion in this compilation involved 9,000, 000 lb (4.5 kilotons of ammonium nitrate), extension of the regression lines to the magnitudes of nuclear weapon yields involves an extrapolation of one or two

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orders of magnitude beyond the limits of plotted data. Their use in estimating maximum debris distance in the range of high-yield nuclear weapons requires extrapolation of three or four orders of magnitude beyond the limits of the plotted data. Thus, while there is no specific interest in the very small HE detonations, they are included in the total sample to extend the range of basic data -- while extrapolating several orders of magnitude beyond the iimits of the data can be considered questionable under many circumstances, it is better to do so with seven cycles of data input than with a lesser amount.

A total of 40 linear and 40 quadratic regression lines, in logarithmic terms, waミ derived by computer methods. Each of the debris distance parameters,
$\log _{10}$ maximum debris distance ( ft )
$\log _{10}$ cube-root maximum debris distance ( $\mathrm{ft}^{1 / 3}$ )
$\log _{10}$ 2/3-power maximum debris distance ( $\mathrm{ft}^{2 / 3}$ )
$\log _{10} \mathrm{~W}^{1 / 3}$-scaled maximum debris distance $\left(\frac{\mathrm{ft}}{\frac{\mathrm{ft}}{\mathrm{tons}} \mathrm{TNT}}\right)$
$\log _{10} \mathrm{w}^{2 / 3}$-scaled maximum debris distance $\left(\frac{\mathrm{ft}}{\operatorname{tons}_{\mathrm{TNT}}^{2 / 3}}\right)$
was correlated against each of the explosion parameters:
$\log _{10}$ equivalent yield (tons ${ }_{\mathrm{TNT}}$ )
$\log _{10}$ impulse $\left(\frac{\mathrm{lb}-\mathrm{msec}}{\text { sq in. }}\right)$
$\log _{10}$ cube-root equivalent yield (tons $\frac{1 / 3}{\mathrm{TNT}}$ )
$\log _{10}$ cube-root impulse ( $\left(\frac{\mathrm{lb}-\mathrm{msec}}{\text { sq in. }}\right)^{1 / 3}$,
$\log _{10} 2 / 3$-powerequivalent yield (tons $\frac{2 / 3}{2 / 3}$ )
$\log _{10} 2 / 3$-power impulse $\left(\left(\frac{l b-m s e c}{s q i n .}\right)^{2 / 3}\right.$,
$\log W^{1 / 3}$-scaled impulse $\left(\frac{\text { lb-msec }}{\mathrm{in}^{2}-\operatorname{tons} \frac{1 / 3}{\text { TNT }}}\right)$
$\log \mathrm{w}^{2 / 3}$-scaled impulse $\left(\frac{\mathrm{lb}-\mathrm{msec}}{\mathrm{in}^{2}-\operatorname{tons} \mathrm{TNT}}\right\rangle$

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The selection of associative relationships to be correlated was arbitrary. No analytical basis for a fixed relationship existed, and no justification for scaled relationships was apparent. By computer methods, the computation of a large number of correlations required little additional effort beyond that for the one or two obvious correlations, and provided the opportunity to check the relative closeness of fit for various approaches. The computer program outlined in Appendix B was written to type out coefficients for the regression lines, the standard error and the correlation coefficient. Regression study results are tabulated in Table 2.1. As this table shows, most correlation coefficients were between 0.51 and 0.69 , which is not a high positive correlation. No single correlation appeared significantly better than all others. The quadratic correlations were only marginally better than the linear correlations. The improved correlation coefficients for the quadratic cases can sometimes be regarded as suspect, in fact, especially if they yield expressions indicating a continually increasing slope with increasing explosion size.

Several correlations were selected for plotting, together with their respective data points. These are presented in the following figures.

Fig. 2. 3 Maximum Debris Distance Vs Equivalent Yield (linear log-log scale)
Fig. 2. 4 Maximum Debris Distance Vs Equivalent Yield
(quadratic log-log scale)
Fig. 2. $5 \mathrm{~W}^{1 / 3}$-Scaled Maximum Debris Distance Vs Equivalent Yield
Fig. $2.6 \mathrm{~W}^{1 / 3}$-Scaled Maximum Debris Distance Vs Equivalent Yield
Fig. 2. 7 Maximum Debris Distance Vs Impulse (linear log - log scale)
Fig. 2. 8 Maximum Debris Distance Vs Impulse (quadratic log-log scale)
Fig. 2.9 $\mathrm{w}^{1 / 3}$-Scaled Maximum Debris Distance Vs $\mathrm{w}^{1 / 3}$-Scaled Impulse (linear log-log scale)
Fig. 2.10 $\mathrm{w}^{1 / 3}$-Scaled Maximum Debris Distance Vs $\mathrm{W}^{1 / 3}$-Scaled Impulse (quadratic log-log scale)
Text follows on page 51


Table 2.1
S FOR MAXIMUM DEBRIS DISTANCE IN TERMS OF EXPLOSION PARAMETERS

| Logarithmic Form |  | Exponential Form | Correlation Coefficient | Figure Number |
| :---: | :---: | :---: | :---: | :---: |
| egression Line | $\begin{gathered} \text { Standard } \\ \text { Error } \end{gathered}$ | Regression Line ${ }^{\text {a }}$ ( $\begin{gathered}\text { Standard } \\ \text { Error }\end{gathered}$ |  |  |
| $13=0.982+0.107 \log _{10} \mathrm{~W}$ | $\pm 0.131$ | $\mathrm{D}_{\mathrm{M}}^{1 / 3}=9.60 \mathrm{w}^{0.107}$ | 0.67 |  |
| $13=0.986+0.116 \log _{10} W-0.00536\left(\log _{10} W\right)^{2}$ | $\pm 0.131$ |  | 0.67 |  |
| $/ 3=0.525+0.139 \log _{10} I$ | $\pm 0.128$ |  | 0.68 |  |
| $i^{\prime}=0.457+0.177 \log _{10} \mathrm{I}-0.00503\left(\log _{10} \mathrm{I}\right)^{2}$ | $\pm 0.128$ | $\mathrm{D}_{\mathrm{M}}^{1 / 3}=2.871^{0.177}(10)^{-0.00503\left(\log _{10} 1\right)^{2}} \quad \stackrel{\mathbf{x}}{ \pm} 1.34$ | 0.68 |  |
| $1 / 3=0.982+0.322 \log _{10} w^{1 / 3}$ | $\pm 0.131$ |  | 0.67 |  |
| $13=0.986+0.347 \log _{10} w^{1 / 3}-0.0484\left(\log _{10} w^{1 / 3}\right)^{2}$ | $\pm 0.131$ | $\begin{array}{l\|l\|l} \left.\hline D_{M}^{1 / 3}=9.68\left(w^{1 / 3}\right)^{0.347}(10)^{-0.0484\left(\log _{10}\right.} w^{1 / 3}\right)^{2} & \underset{\div}{\div} 1.35 \end{array}$ | 0.67 |  |
| $1 / 3=0.525+0.416 \log _{10} 1^{1 / 3}$ | $\pm 0.128$ |  | 0.68 |  |
| /3 $=0.457+0.531 \log _{10} \mathrm{I}^{1 / 3}-0.0454\left(\log _{10} \mathrm{I}^{1 / 3}\right)^{2}$ | $\pm 0.128$ | $\left.D_{M}^{1 / 3}=2.87\left(1^{1 / 3}\right)^{0.531}(10)^{-0.0454\left(\log _{10}\right.} 1^{1 / 3}\right)^{2}-\underset{\div}{\div 1.34}$ | 0.68 |  |
| $1 / 3=0.982+0.161 \log _{10} w^{2 / 3}$ | $\pm 0.131$ | $\mathrm{D}_{\mathrm{M}}^{1 / 3}=9.60\left(\mathrm{~W}^{2 / 3}\right)^{0.161} \mathrm{~m}^{2 / 32} \ldots$ | 0.67 |  |
| $1 / 3=0.986+0.173 \log _{10} w^{2 / 3}-0.0121\left(\log _{10} w^{2 / 3}\right)^{2} \pm 0.131$ |  | $\left.D_{M}^{1 / 3}=9.68\left(\mathrm{w}^{2 / 3}\right)^{0.173}(10)^{-0.0121\left(\log _{10}\right.} \mathrm{w}^{2 / 3}\right)^{2}-\left\lvert\, \begin{array}{\|c} \div \\ \hline \end{array}\right.$ | 0.66 |  |
| $1 / 3=0.525+0.208 \log _{10} 1^{2 / 3}$ | $\pm 0.128$ |  | 0.68 |  |
| ${ }^{1 / 3}=0.457+0.265 \log _{10} 1^{2 / 3}-0.0113\left(\log _{10} I^{2 / 3}\right)^{2}$ | $\pm 0.128$ | $\left.\mathrm{D}_{\mathrm{M}}^{1 / 3}=2.87\left(\mathrm{I}^{2 / 3}\right)^{0.265}(10)^{-0.0113\left(\log _{10}\right.} 1^{2 / 3}\right)^{2}-\left\lvert\, \begin{aligned} & \dot{x} 1.34 \\ & \hline \end{aligned}\right.$ | 0.68 |  |
| $1 / 3=-0.779+0.544 \log _{10}\left(\frac{1}{W^{1 / 3}}\right)$ | $\pm 0.145$ | $\mathrm{D}_{\mathrm{M}}^{1 / 3}=0.166\left(\frac{\mathrm{I}}{\mathrm{~W}^{1 / 3}}\right)^{0.544}$ | 0.57 |  |
| ${ }^{3}=4.689-2.893 \log _{10}\left(\frac{1}{w^{1 / 3}}\right)+0.537\left[\log _{10}\left(\frac{1}{w^{1 / 3}}\right)\right]^{2}$ | $\pm 0.142$ | $D_{M}^{1 / 3}=4.89 \times 10^{4}\left(\frac{1}{W^{1 / 3}}\right)^{-2.893}(10)^{-0.537}\left(\left.\log _{10}\left(\frac{1}{W^{1 / 3}}\right)^{2} \right\rvert\, \frac{x}{1.39}\right.$ | 0.59 |  |
| $1 / 3=2.104-0.333 \log _{10}\left(\frac{1}{W^{2 / 3}}\right)$ | $\pm 0.151$ | $\mathrm{D}_{\mathrm{M}}^{1 / 3}=1.270 \times 10^{2}\left(\frac{1}{\mathrm{w}^{2 / 3}}\right)^{-0.333}{ }^{\text {a }}$ ( $]_{2}^{ \pm 1.42}$ | 0.51 |  |
| ${ }^{3}=3.388-1.153 \log _{10}\left(\frac{1}{w^{2 / 3}}\right)+0.131\left[\log _{10}\left(\frac{I}{w^{2 / 3}}\right)\right]^{2}$ | $\pm 0.150$ | $D_{M}^{1 / 3}=2.444 \times 10^{3}\left(\frac{I}{W^{2 / 3}}\right)^{-1.154}(10)^{0.130}\left[\left.\log _{10}\left(\frac{1}{W^{2 / 3}}\right)\right\|^{\frac{x}{\div} 1.41}\right.$ | 0.51 |  |


| Factors Correlated | Type of Correlation | Logarithmic Form |
| :---: | :---: | :---: |
|  |  | Regression Line |
| $\log _{10}$ Maximum Debris Distance ( ft ) <br> Versus <br> $\log _{10}$ Equivalent Yield ( ${ }^{\text {(ons }}{ }^{\text {TNT }}$ ) | Linear | $\log _{10} D_{M}=2.950+0.322 \log _{10} W$ |
|  | Quadratic | $\log _{10} D_{M}=2.960+0.347 \log _{10} W-0.0161\left(\log _{10} W\right)^{2}$ |
| $\begin{aligned} & \log _{10} \text { Maximum Debris Distance (ft) } \\ & \text { Versus } \\ & \log _{10} \text { Impulse (lb-msec/in. }{ }^{2} \text { ) } \end{aligned}$ | Linear | $\log _{10} \mathrm{D}_{\mathrm{M}}=1.578+0.416 \log _{10} \mathrm{I}$ |
|  | Quadratic | $\log _{10}{ }^{D}{ }_{M}=1.373+0.531 \log _{10} 1-0.0151\left(\log _{10} I\right)^{2}$ |
| $\log _{10}$ Maximum Debris Distance (ft) <br> Versus $\log _{10} \text { Cube-Root Equivalent Yield (tons } \frac{1 / 3}{\text { TNT }} \text { ) }$ | Linear | $\log _{10} D_{M}=2.950+0.968 \log _{10} w^{1 / 3}$ |
|  | Quadratic | $\log _{10} D_{M}=2.960+1.042 \log _{10} W^{1 / 3}-0.145\left(\log _{10} W^{1 / 3}\right)^{2}$ |
| $\log _{10}$ Maximum Debris Distance ( ft ) <br> Versus $\left.\operatorname{Iog}_{10} \text { Cube-Root Impulse (lb-msec/in. } 2\right)^{1 / 3}$ | Linear | $\log _{10} D_{M}=1.578+1.250 \log _{10} I^{1 / 3}$ |
|  | Quadratic | $\log _{10} D_{M}=1.373+1.593 \log _{10} I^{1 / 3}-0.136\left(\log _{10} I^{1 / 3}\right)^{2}$ |
| $\begin{gathered} \log _{10} \text { Maximum Debria Distance (ft) } \\ \text { Versus } \\ \log _{10} \frac{2}{3}-\text { Power Equivalent Yield (tons } \frac{2 / 3}{\text { TNT }} \text { ) } \end{gathered}$ | Linear | $\log _{10} D_{M}=2.950+0.483 \log _{10} \mathrm{w}^{2 / 3}$ |
|  | Quadratic | $\log _{10} D_{M}=2.960+0.520 \log _{10} w^{2 / 3}-0.0362\left(\log _{10} w^{2 / 3}\right)^{2}$ |
| $\log _{10}$ Maximum Debris Distance (ft) Versus $\log _{10} \frac{2}{3}$-Power Impulse (lb-msec/in. ${ }^{2}$ ) ${ }^{2 / 3}$ | Linear | $\log _{10} D_{M}=1.578+0.624 \log _{10} I^{2 / 3}$ |
|  | Quadratic | $\log _{10} D_{M}=1.373+0.795\left(\log _{10} I^{2 / 3}-0.0340\left(\log _{10} I^{2 / 3}\right)^{2}\right.$ |
| Log $_{10}$ Maximum Debris Distance (ft) <br> Versus $\log _{10} \mathrm{w}^{1 / 3} \text { - Scaled Impulse (lb-msec/in. }{ }^{2} \text {-tons } \frac{1 / 3}{\mathrm{TNT}} \text { ) }$ | Linear | $\log _{10} D_{M}=-2.338+1.634 \log _{10}\left(\frac{\mathrm{I}}{\mathrm{W}^{1 / 3}}\right)$ |
|  | Quadratic | $\log _{10} D_{M}=14.08-8.69 \log _{10}\left(\frac{1}{W^{1 / 3}}\right)+1.612\left[\log _{10}\left(\frac{1}{W^{1 / 3}}\right)\right]^{2}$ |
| $\begin{aligned} & \log _{10} \text { Maximum Debris Distance (ft) } \\ & \text { Versus } \\ & \left.\log _{10} \mathrm{w}^{2 / 3}-\text { Scaled Impulse (lb-msec/in. } .^{2} \text {-tons } \frac{2 / 3}{T N T}\right) \end{aligned}$ | Linear | $\log _{10} D_{M}=6.317-1.001 \log _{10}\left(\frac{\mathrm{I}}{\mathrm{w}^{2 / 3}}\right)$ |
|  | Quadratic | $\log _{10} D_{M}=10.17-3.462 \log _{10}\left(\frac{\mathrm{I}}{W^{2 / 3}}\right)+0.390\left[\log _{10}\left(\frac{1}{w^{2 / 3}}\right)\right]^{2}$ |

Table 2.1 (Continued)


| Logarithmic Form |  | Exponential Form |  | Corrolation Coefficient | Figure <br> Number |
| :---: | :---: | :---: | :---: | :---: | :---: |
| zgression Line | Standard Error | Regression Line | $\begin{aligned} & \text { Standard } \\ & \text { Error } \end{aligned}$ |  |  |
| $=2.950+0.322 \log _{10} \mathrm{~W}$ | $\pm 0.392$ | $\mathrm{D}_{\mathrm{M}}=892 \mathrm{~W}^{0.322}$ | $\stackrel{\times}{\div} 2.47$ | 0.67 | 2.3 |
| $=2.960+0.347 \log _{10} W-0.0161\left(\log _{10} W\right)^{2}$ | $\pm 0.392$ | $\mathrm{D}_{\mathrm{M}}=912 \mathrm{~W}^{0.347}(10)^{-0.0161(\mathrm{Log} \mathrm{W})^{2}}$ | $\stackrel{\mathrm{x}}{\div} 2.47$ | 0.67 | 2.4 |
| $=1.578+0.416 \log _{10} \mathrm{I}$ | $\pm 0.385$ | $\mathrm{D}_{\mathrm{M}}=38 \mathrm{I}^{0.416}$ | $\stackrel{\times}{\square} 2.43$ | 0.68 | 2.7 |
| $=1.373+0.531 \log _{10} 1-0.0151\left(\log _{10} 1\right)^{2}$ | $\pm 0.386$ | $\left.D_{M}=241^{0.531}(10)^{-0.0151(L o g} 10 \mathrm{I}\right)^{2}$ | $\stackrel{\times}{¢} 2.43$ | 0.68 | 2.8 |
| $=2.950+0.968 \log _{10} \mathrm{w}^{1 / 3}$ | $\pm 0.392$ | $\mathrm{D}_{\mathrm{M}}=892\left(\mathrm{~W}^{1 / 3}\right)^{0.968}$ | $\stackrel{x}{\div} 2.47$ | 0.67 |  |
| $=2.960+1.042 \log _{10} w^{1 / 3}-0.145\left(\log _{10} w^{1 / 3}\right)^{2}$ | $\pm 0.392$ | $\left.D_{M}=912\left(w^{1 / 3}\right)^{1.042}(10)^{-0.145\left(\log _{10}\right.} w^{1 / 3}\right)^{2}$ | $\stackrel{\times}{\div} 2.47$ | 0.67 |  |
| $=1.578+1.250 \log _{10} I^{1 / 3}$ | $\pm 0.385$ | $D_{M}=38\left(1^{1 / 3}\right)^{1.250}$ | $\stackrel{\mathbf{x}}{\stackrel{1}{+} 2.43}$ | 0.68 |  |
| $=1.373+1.593 \log _{10} 1^{1 / 3}-0.136\left(\log _{10} I^{1 / 3}\right)^{2}$ | $\pm 0.386$ | $D_{M}=24\left(1^{1 / 3}\right)^{1.593}(10)^{-0.136\left(\log _{10} I^{1 / 3}\right)^{2}}$ | $\stackrel{\times}{\times} 2.43$ | 0.68 |  |
| $=2.950+0.483 \log _{10} \mathrm{w}^{2 / 3}$ | $\pm 0.392$ | $\mathrm{D}_{\mathrm{M}}=892\left(\mathrm{w}^{2 / 3}\right)^{0.484}$ | $\stackrel{\mathbf{x}}{\div} 2.47$ | 0.67 |  |
| $=2.960+0.520 \log _{10} w^{2 / 3}-0.0362\left(\log _{10} w^{2 / 3}\right)^{2}$ | $\pm 0.392$ | $\left.D_{M}=912\left(w^{2 / 3}\right)^{0.520}(10)^{-0.0362\left(\log _{10}\right.} w^{2 / 3}\right)^{2}$ | $\stackrel{x}{\div} 2.47$ | 0.67 |  |
| $=1.578+0.624 \log _{10} I^{2 / 3}$ | $\pm 0.385$ | $D_{M}=38\left(1^{2 / 3}\right){ }^{0.624}$ | $\stackrel{\mathbf{x}}{\div} 2.43$ | 0.68 |  |
| $=1.373+0.795\left(\log _{10} \mathrm{I}^{2 / 3}-0.0340\left(\log _{10} \mathrm{I}^{2 / 3}\right)^{2}\right.$ | $\pm 0.386$ | $\left.D_{M}=24\left(I^{2 / 3}\right)^{0.795}(10)^{-0.0340\left(\log _{10}\right.} \mathrm{I}^{2 / 3}\right)^{2}$ | $\stackrel{\mathbf{x}}{\div} 2.43$ | 0.68 |  |
| $=-2.338+1.634 \log _{10}\left(\frac{\mathrm{I}}{\mathrm{W}^{1 / 3}}\right)$ | $\pm 0.434$ | $D_{M}=0.00459\left(\frac{1}{W^{1 / 3}}\right)^{1.634}$ | $\stackrel{\mathrm{x}}{\div} 2.72$ | 0.57 |  |
| 14.08-8.69 $\log _{10}\left(\frac{\mathrm{I}}{\mathrm{w}^{1 / 3}}\right)+1.612\left[\log _{10}\left(\frac{\mathrm{I}}{\mathrm{w}^{1 / 3}}\right)\right]^{2}$ | $\pm 0.426$ | $D_{M}=1.211 \times 10^{14}\left(\frac{1}{W^{1 / 3}}\right)^{-8.69}(10)^{-1612}\left[\log _{10}\left(\frac{1}{W^{1 / 3}}\right)\right]$ | $\stackrel{\mathrm{x}}{\div} \mathbf{2 . 6 6}$ | 0.59 |  |
| $=6.317-1.001 \log _{10}\left(\frac{I}{W^{2 / 3}}\right)$ | $\pm 0.456$ | $D_{M}=2.078 \times 10^{6}\left(\frac{1}{W^{2 / 3}}\right)^{-1.001}$ | $\stackrel{\mathbf{x}}{\div} \mathrm{C} 2.84$ | 0.51 |  |
| $=10.17-3.462 \log _{10}\left(\frac{1}{w^{2 / 3}}\right)+0.390\left[\log _{10}\left(\frac{\mathrm{I}}{W^{2 / 3}}\right)\right]^{2}$ | $\pm 0.451$ | $D_{M}=1.483 \times 10^{10}\left(\frac{1}{W^{2 / 3}}\right)^{-3.462}(10)^{0.390}\left[\log _{10}\left(\frac{1}{W^{2 / 3}}\right)\right]$ | $\stackrel{\mathrm{x}}{\div} \mathrm{C} 2.83$ | 0.51 |  |


| Factors Correlated | Type of Correlation | Logarithmic Form |
| :---: | :---: | :---: |
|  |  | Regression Line |
| $\begin{gathered} \log _{10} \frac{2}{3} \text { - Power Maximum Debria Distance }\left(\mathrm{ft}^{2 / 3}\right) \\ \text { Versus } \\ \log _{10} \text { Equivalent Yield (tons } \mathrm{TNT} \text { ) } \end{gathered}$ | Linear | $\log _{10} \mathrm{D}_{\mathrm{M}^{2 / 3}}=1.968+0.215 \log _{10} \mathrm{~W}$ |
|  | Quadratic | $\log _{10} \mathrm{D}_{\mathrm{M}^{2 / 3}}=1.974+0.231 \log _{10} W-0.0108\left(\log _{10} \mathrm{~W}\right)^{2}$ |
| $\begin{gathered} \log _{10} \frac{2}{3} \text { - Power Maximum Debris Distance }\left(\mathrm{ft}^{2 / 3}\right) \\ \text { Versus } \\ \left.\log _{10} \text { Impulse (lb-msec/in. }{ }^{2}\right) \end{gathered}$ | Linear | $\begin{array}{\|l} \log _{10} D_{M}^{2 / 3}=1.052+0.278 \log _{10} I \\ \log _{10} D_{M}^{2 / 3}=0.916+0.354 \log _{10} I-0.0101\left(\log _{10} I\right)^{2} \end{array}$ |
|  | Quadratic |  |
| $\log _{10} \frac{2}{3}$ - Power Maximum Debris Distance $\left(\mathrm{ft}^{2 / 3}\right)$ Versus <br> $\log _{10}$ Cube-Root Equivalent Yield (tond $1 / \mathbf{3}$ ) | Linear | $\frac{\log _{10} D_{M}^{2 / 3}=1.968+0.646 \log _{10} W^{1 / 3}}{\log _{10} D_{M}^{2 / 3}=1.974+0.695 \log _{10} W^{1 / 3}-0.0969\left(\log _{10} W^{1 / 3}\right)}$ |
|  | Quadratic |  |
| $\log _{10} \frac{2}{3}$ - Power Maximum Debris Distance ( $\mathrm{ft}^{2 / 3}$ ) Versus <br> Log $_{10}$ Cube-Root Impulse (lb-msec/in. $\left.{ }^{2}\right)^{1 / 3}$ | Linear | $\log _{10} \mathrm{D}_{\mathrm{M}}^{2 / 3}=1.052+0.834 \log _{10} I^{1 / 3}$ |
|  | Quadratic | $\log _{10} D_{M}^{2 / 3}=0.916+1.062 \log _{10} I^{1 / 3}-0.0909\left(\log _{10} I^{1 / 3}\right)^{2}$ |
| $\log _{10} \frac{2}{3}$ - Power Maximum Debris Distance ( $\mathrm{ft}{ }^{2 / 3}$ ) Versus <br> $\log _{10} \frac{2}{3}$-Power Equivalent Yield (tons $\frac{2 / 3}{}$ TNT) | Linear | $\log _{10} D_{M}^{2 / 3}=1.968+0.322 \log _{10} w^{2 / 3} 1 \log _{10} D_{M}^{2 / 3}=1.974+0.347 \log _{10} w^{2 / 3}-0.0242\left(\log _{10} w^{2 / 3}\right)^{2}$ |
|  | Quadratic |  |
| $\log _{10} \frac{2}{3}$ - Power Maximum Debris Distance ( $\mathrm{ft}^{2 / 3}$ ) <br> Versus <br> $\log _{10} \frac{2}{3}$-Power Impulse (lb-msec/in. $\left.{ }^{2}\right)^{2 / 3}$ | Linear | $\log _{10} D_{M}^{2 / 3}=1.052 \div 0.10 \log _{10} I^{2 / 3} 1 \log _{10} D_{M}^{2 / 3}=0.916+0.530 \log _{10} I^{2 / 3}-0.0227\left(\log _{10} I^{2 / 3}\right)^{2}$ |
|  | Quadratic |  |
| $\log _{10} \frac{2}{3}$ - Power Maximum Debris Distance ( $f t^{2 / 3}$ ) Versus | Linear | $\log _{10} D_{M}^{2 / 3}=1.560+1.090 \log _{10}\left(\frac{I}{w^{1 / 3}}\right)$ |
| $\log _{10} \mathrm{w}^{1 / 3} \text { - Scaled Impulse (lb-msec/in. }{ }^{2} \text {-tons } \frac{1 / 3}{} \text { TNT }$ | Quadratic | $\log _{10} \mathrm{D}_{\mathrm{M}}^{2 / 3}=9.405-5.803 \log _{10}\left(\frac{1}{w^{1 / 3}}\right)+1.077\left[\log _{10}\left(\frac{1}{w^{1 / 3}}\right)\right]$ |
| $\log _{10} \frac{2}{3}$ - Power Maximum Debris Distance ( $\mathrm{ft}{ }^{2 / 3}$ ) | Linear | $\log _{10} D_{M}^{2 / 3}=4.214-0.0667 \log _{10}\left(\frac{I}{W^{2 / 3}}\right)$ |
| $\log _{10} \mathrm{w}^{2 / 3} \text { - Scaled Impulse (Ib-msec/in. }{ }^{2} \text {-tons } \frac{2 / 3}{\mathrm{TNT}} \text { ) }$ | Quadratic | $\log _{10} D_{M}^{2 / 3}=6.786-2.310 \log _{10}\left(\frac{1}{w^{2 / 3}}\right)+0.0260\left[\log _{10}\left(\frac{1}{w^{2 / 3}}\right)\right]$. |


| Logarithmic form |  | Exponential Form |  | Correlation Confficient | Figure Number |
| :---: | :---: | :---: | :---: | :---: | :---: |
| iegression Line | Standard Error | Regression Line | $\begin{aligned} & \text { Standard } \\ & \text { Error } \\ & \hline \end{aligned}$ |  |  |
| $M^{2 / 3}=1.968+0.215 \log _{10} W$ | $\pm 0.262$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=92.9 \mathrm{w}^{0.215}$ | $\underset{\sim}{\times} 1.83$ | 0.67 |  |
| $u^{2 / 3}=1.974+0.231 \log _{10} W-0.0108\left(\log _{10} w\right)^{2}$ | $\pm 0.262$ | $D_{M}^{2 / 3}=94.3 \mathrm{w}^{0.231}(10)^{-0.0107\left(\log _{10} W\right)^{2}}$ | $\underset{\sim}{\times}$ | 0.67 |  |
| $M^{2 / 3}=1.052+0.278 \log _{10} I$ | $\pm 0.257$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=11.28 \mathrm{I}^{0.278}$ | $\underset{\sim}{\square} 1.81$ | 0.68 |  |
| $u^{2 / 3}=0.916+0.354 \log _{10} 1-0.0101\left(\log _{10} 1\right)^{2}$ | $\pm 0.257$ | $\mathrm{D}_{\mathrm{M}^{2 / 3}}=8.24 \mathrm{I}^{0.354}(10)^{-0.0101\left(\log _{10}\right)^{2}}$ | $\begin{array}{cc}\times 1.81 \\ \div & 1.81\end{array}$ | 0.68 |  |
| $M^{2 / 3}=1.968+0.646 \log _{10} w^{1 / 3}$ | $\pm 0.262$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=92.86\left(\mathrm{w}^{1 / 3}\right)^{0.646}$ | x $\div$ $\div$ | 0.67 |  |
| $M^{2 / 3}=1.974+0.695 \log _{10} W^{1 / 3}-0.0069\left(\log _{10} w^{1 / 3}\right)^{2}$ | $\pm 0.262$ | $\left.D_{M}^{2 / 3}=94.27\left(w^{1 / 3}\right)^{0.695}(10)^{-0.0969\left(\log _{10}\right.} w^{1 / 3}\right)^{2}$ | $\times 1.83$ $\div$ | 0.67 |  |
| $M_{M}^{2 / 3}=1.052+0.834 \log _{10} I^{1 / 3}$ | $\pm 0.257$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=11.28\left(\mathrm{I}^{1 / 3}\right)^{0.834}$ | $\underset{\sim}{\times} 1.81$ | 0.68 |  |
| $\mathbf{A}^{2 / 3}=0.916+1.062 \log _{10} I^{1 / 3}-0.0909\left(\log _{10} I^{1 / 3}\right)^{2}$ | $\pm 0.257$ | $\left.D_{M}^{2 / 3}=8.24\left(I^{1 / 3}\right)^{1.062}(10)^{-0.0909\left(\log _{10}\right.} I^{1 / 3}\right)^{2}$ | $\stackrel{\times}{\div} \div 1.81$ | 0.68 |  |
| $M^{2 / 3}=1.968+0.322 \log _{10} W^{2 / 3}$ | $\pm 0.262$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=92.86\left(\mathrm{w}^{2 / 3}\right)^{0.322}$ | $\stackrel{\times}{\square} 1.83$ | 0.67 |  |
| $\mathrm{M}^{2 / 3}=1.974+0.347 \log _{10} \mathrm{w}^{2 / 3}-0.0242\left(\log _{10} w^{2 / 3}\right)^{2}$ | $\pm 0.262$ | $D_{M}^{2 / 3}=94.27\left(w^{2 / 3}\right)^{0.347}(10)^{-0.0242\left(\log _{10}\right.} w^{2 / 3)^{2}}$ | $\stackrel{\times}{\div} \times 1.83$ | 0.67 |  |
| $M_{M}^{2 / 3}=1.052+0.416 \log _{10} I^{2 / 3}$ | $\pm 0.257$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=11.28\left(\mathrm{I}^{2 / 3}\right)^{0.416}$ | $\stackrel{\times}{\div} 1.81$ | 0.68 |  |
| $u^{2 / 3}=0.916+0.530 \log _{10} I^{2 / 3}-0.0227\left(\log _{10} I^{2 / 3}\right)^{2}$ | $\pm 0.257$ | $\left.D_{M}^{2 / 3}=8.24\left(I^{2 / 3}\right) 0.530(10)^{-0.0227\left(\log _{10}\right.} I^{2 / 3}\right)^{2}$ | $\underset{\square}{\div} 1.81$ | 0.68 |  |
| $M_{M}^{2 / 3}=1.560+1.090 \log _{10}\left(\frac{I}{W^{l / 3}}\right)$ | $\pm 0.290$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=2.756 \times 10^{-2}\left(\frac{\mathrm{I}}{\mathrm{~W}^{1 / 3}}\right)^{1.090}$ | $\times 1.95$ | 0.57 |  |
| $u^{2 / 3}=9.405-5.803 \log _{10}\left(\frac{1}{w^{1 / 3}}\right)+1.077\left[\log _{10}\left(\frac{1}{w^{1 / 3}}\right)\right]^{2}$ | $\pm 0.284$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=2.54 \times 10^{9}\left(\frac{\mathrm{I}}{\mathrm{w}^{2 / 3}}\right)^{-5.80}(10)^{1.077}\left[\log _{10}\left(\frac{I}{w^{\frac{1}{2} / 3}}\right)\right]$ | x $\div 1.92$ | 0.59 |  |
| $M^{2 / 3}=4.214-0.0667 \log _{10}\left(\frac{I}{W^{2 / 3}}\right)$ | $\pm 0.303$ | $\mathrm{D}_{\mathrm{M}}^{2 / 3}=1.636 \times 10^{4}\left(\frac{1}{W^{2 / 3}}\right)^{-0.0667}$ | $\stackrel{x}{\div} 2.01$ | 0.51 |  |
| $\left.M^{2 / 3}=6.786-2.310 \log _{10}\left(\frac{I}{W^{2 / 3}}\right)+0.0260!\log _{10}\left(\frac{1}{W^{2 / 3}}\right)\right]^{2}$ | $\pm 0.301$ | $D_{M}^{2 / 3}=6.107 \times 10^{6}\left(\frac{1}{W^{2 / 3}}\right)^{-2.310}(10)^{0.0262}\left[\log _{10}\left(\frac{I}{W^{2} / 3}\right)\right]$ | $\times 2.00$ $\div$ | 0.51 |  |


| Factors Correlated | $\begin{array}{\|c} \text { Type } \\ \text { of } \\ \text { Correlation } \end{array}$ | Logarithmic Form |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Regression Line |  |  |  |
| $\begin{array}{\|l} \log _{10} \mathrm{~W}^{1 / 3}-\text { Scaled Maximum Debris Distance (f/tons } \mathrm{T} \mathrm{NT} \text { ) } \\ \begin{array}{c} \text { Versus } \end{array} \\ \log _{10} \text { Equivalent Yield (tons } \mathrm{TNT} \text { ) } \end{array}$ | Linea | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.950-0.0105 \log _{10} \mathrm{~W}$ |  |  |  |
|  | Quadratic | $\log _{10}\left(\frac{D_{M}}{w^{1 / 3}}\right)=2.960+0.0140 \log _{10} w-0.0161\left(\log _{10} w\right\}^{2}$ |  |  |  |
| $\begin{array}{\|l} \log _{10} W^{1 / 3}-\text { Scaled Maximum Debris Distance (ft/tons } \mathrm{TNT}^{1 / 3} \text { ) } \\ \text { Versus } \\ \log _{10} \text { Impulse (lb-msec/in. }{ }^{2} \text { ) } \end{array}$ | Linea | $\log _{10}\left(\frac{D_{M}}{W^{W / 3}}\right)=2.941+0.00048 \log _{10} 1$ |  |  |  |
|  | Quadrati | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.401+0.302 \log _{10} 1-0.399\left(\log _{10} 1\right)^{2}$ |  |  |  |
| $\begin{aligned} & \log _{10} W^{1 / 3}-\text { Scaled Maximum Debris Distance ( } \mathrm{f} / \text { /tons } \frac{1 / 3}{} \mathrm{TNT} \text { ) } \\ & \begin{array}{c} \text { Versus } \end{array} \\ & \log _{10} \text { Cube-Root Equivalent Yield (tons } \frac{1 / 3}{\mathrm{TNT}} \text { ) } \end{aligned}$ | Linear | $\log _{10}\left(\frac{D_{M}}{w^{1 / 3}}\right)=2.950-0.0316 \log _{10} w^{1 / 3}$ |  |  |  |
|  | Quadrati | $\log _{10}\left(\frac{S_{M}}{w^{1 / 3}}\right)=2.960+0.042 \log _{10} w^{1 / 3}-0.145\left(\log _{10} w^{1 / 3}\right)^{2}$ |  |  |  |
| $\begin{aligned} & \log _{10} \mathrm{~W}^{1 / 3}-\text { Scaled Maximum Debris Distance }(\mathrm{f} / \text { tons } \mathrm{TNT}) \\ & \text { Versus } \\ & \left.\log _{10} \text { Cube-Root Impulse (lb-msec/in. }{ }^{2}\right)^{1 / 3} \end{aligned}$ | Linear | $\log _{10}\left(\frac{D_{M}}{W^{7 / 3}}\right)=2.941+0.00144 \log _{10} 1^{1 / 3}$ |  |  |  |
|  | Quadratic | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.401+0.907 \log _{10} 1^{1 / 3}-0.360\left(\log ^{1 / 3}\right)^{2}$ |  |  |  |
|  | Linear | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.950-0.0158 \log _{10} w^{2 / 3}$ |  |  |  |
| $\log _{10} \begin{gathered}\text { Versus } \\ \frac{2}{3} \text {-Power Equivalent Yield (tons }\end{gathered} \frac{2 / 3}{}{ }^{2 / 3}$ ) | Quadrati | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.960+0.0210 \log _{10} w^{2 / 3}-0.0362\left(\log _{10} w^{2 / 3}\right)^{2}$ |  |  |  |
| $\begin{aligned} & \log _{10} \mathrm{w}^{1 / 3}-\text { Scaled Maximum Debris Distance }\left(\mathrm{ft} / \mathrm{tons}_{\mathrm{TNT}} \mathrm{I} / 3\right. \text { ) } \\ & \text { Versus } \\ & \log _{10} \frac{2}{3} \text { - Power Impulse }\left(\mathrm{lb}-\mathrm{msec} / \mathrm{in.}^{2}\right)^{2 / 3} \end{aligned}$ | Linea | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.941+0.00072 \log _{10} \mathrm{I}^{2 / 3}$ |  |  |  |
|  | Quadrat | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.401+0.453 \log _{10} 1^{2 / 3}-0.0898\left(\log _{10} 1^{2 / 3}\right)^{2}$ |  |  |  |
|  | Linear | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.067+0.260 \log _{10}\left(\frac{1}{w^{1 / 3}}\right)$ |  |  |  |
|  | Quadratic | $\log _{10}\left(\frac{D}{W^{1 / 3}}\right)=2.504-0.0150 \log _{10}\left(\frac{1}{w^{1 / 3}}\right)+0.0429[$ |  |  | $\log _{10}\left(\frac{1}{w^{1 / 3}}\right)^{\text {( }}$ |
| $\left\lvert\, \begin{gathered} \log _{10} \mathrm{w}^{1 / 3}-\text { Scaled Maximum Debris Distance }\left(\mathrm{ft} / \operatorname{tons} \frac{1 / 3}{\mathrm{TNT}}\right. \text { ) } \\ \text { Versus } \\ \log _{10} \mathrm{w}^{2 / 3}-\text { Scaled Impulse ( }\left(\mathrm{lb}-\mathrm{msec} / \mathrm{in}^{2}{ }^{2}-\text { tons }_{\mathrm{TNT}}^{2 / 3}\right) \end{gathered}\right.$ | Linea | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.382+0.178 \log _{10}\left(\frac{\mathrm{I}}{\mathrm{w}^{2 / 3}}\right)$ |  |  |  |
|  | Quadratic | $\log _{10}\left(\frac{\mathrm{D}}{\mathrm{~W}^{1 / 3}}\right)=3.369-0.452 \log _{10}\left(\frac{\mathrm{I}}{\mathrm{w}^{2 / 3}}\right)+0.100\left[\log _{10}\left(\frac{\mathrm{I}}{\mathrm{w}^{2 / 3}}\right)\right]^{2}$ |  |  |  |

Table 2.1 (Continued)

| Logarithmic Form |  | Exponential Form |  | Correlation Coefficient | Figure Number |
| :---: | :---: | :---: | :---: | :---: | :---: |
| gression Line | Standard Error | Regression Line | $\begin{aligned} & \text { Standard } \\ & \text { Error } \end{aligned}$ |  |  |
|  | $\pm 0.392$ | $\frac{D_{M}}{\mathbf{w}^{1 / 3}}=892 \mathrm{w}^{-0.0105}$ | $\underset{+}{\underset{+}{+}} 2.47$ | 0.029 | 2.5 |
| $\bar{M})=2.960+0.0140 \log _{10} w-0.0161\left(\log _{10} w\right)^{2}$ | $\pm 0.3 n 2$ | $\left.\frac{D_{M}}{W^{1 / 3}}=912 \mathrm{w}^{0.0140}(10)^{-0.0161\left(\log _{10}\right.} \mathrm{w}\right)^{2}$ | $\stackrel{x}{\div} \mathbf{~} 2.47$ | 0.059 | 2.6 |
| $\left.\begin{array}{l}\frac{M}{M} \\ \underline{1 / 3}\end{array}\right)=2.941+0.00048 \log _{10} 1$ | $\pm 0.392$ | $\frac{D_{M}}{W^{1 / 3}}=8741^{0.00048}$ | $\stackrel{\mathrm{x}}{\div} \mathbf{+} 2.47$ | 0.0012 |  |
| $\left(\frac{M}{1 / 3}\right)=2.401+0.302 \log _{10} 1-0.399\left(\log _{10} 1\right)^{2}$ | $\pm 0.391$ | $\frac{\mathrm{D}_{\mathrm{M}}}{\mathrm{~W}^{1 / 3}}=252 \mathrm{I}^{0.302(10)^{-0.0399\left(\log _{10} 1\right)^{2}}}$ | $\stackrel{\mathrm{x}}{\div} \mathbf{+} 2.46$ | 0.053 |  |
| $\bar{M}(7 / 3)=2.950-0.0316 \log _{10} w^{1 / 3}$ | $\pm 0.392$ | $\frac{D_{M}}{W^{1 / 3}}=892\left(w^{1 / 3}\right)^{-0.0316}$ | $\underset{ }{\times} 2.47$ | 0.029 |  |
| $\bar{M})=2.960+0.042 \log _{10} w^{1 / 3}-0.145\left(\log _{10} w^{1 / 3}\right)^{2}$ | $\pm 0.392$ | $\frac{D_{M}}{w^{1 / 3}}=912\left(w^{1 / 3}\right)^{0.0}$ <br> (10) | $\stackrel{\times}{\div} 2.47$ | 0.059 |  |
| $\left.\frac{M}{\nabla / 3}\right)=2.941+0.00144 \log _{10} 1^{1 / 3}$ | $\pm 0.392$ | $\frac{D_{M}}{W^{1 / 3}}=874\left(I^{1 / 3}\right)^{0.00144}$ | $\stackrel{x}{\div} 2.47$ | 0.0010 |  |
| $M(\bar{M})=2.401+0.907 \log _{10} 1^{1 / 3}-0.360\left(\operatorname{Log~I}^{1 / 3}\right)^{2}$ | $\pm 0.391$ | $\frac{D_{M}}{W^{1 / 3}}=252\left(I^{1 / 3}\right)^{0.907}(10)^{-0.360\left(\log _{10}\right.}$ | $\underset{\sim}{\div} \mathbf{\square} 2.46$ | 0.053 |  |
| $\left.\frac{M}{1 / 3}\right)=2.950-0.0158 \log _{10} w^{2 / 3}$ | $\pm 0.392$ | $\frac{\mathrm{D}_{\mathrm{M}}}{\mathrm{w}^{1 / 3}}=892\left(\mathrm{w}^{2 / 3}\right)^{-0.0158}$ | $\stackrel{x}{\div} 2.47$ | 0.029 |  |
| $(\bar{M})=2.960+0.0210 \log _{10} w^{2 / 3}-0.0362\left(\log _{10} w^{2 / 3}\right)^{2}$ | $\pm 0.392$ | $\frac{D_{M}}{w^{1 / 3}}=912\left(w^{2 / 3}\right)^{0.0210}(10)^{-0.0362\left(\log _{10}\right.}$ | $\stackrel{x}{\div} 2.47$ | 0.059 |  |
| $\left(\frac{D_{M}}{1 / 3}\right)=2.941+0.00072 \log _{10} I^{2 / 3}$ | $\pm 0.392$ | $\frac{D_{M}}{W^{1 / 3}}=873\left(\mathrm{I}^{2 / 3}\right)^{0.00072}$ | $\stackrel{x}{\div} 2.47$ | 0.0012 |  |
| $\left.\frac{M}{I / 3}\right)=2.401+0.453 \log _{10} \mathrm{I}^{2 / 3}-0.0898\left(\log _{10} \mathrm{I}^{2 / 3}\right)^{2}$ | $\pm 0.391$ | $\frac{D_{M}}{W^{1 / 3}}=252\left(1^{2 / 3}\right)_{(10)^{0.453}}=0.0898\left(\log _{10} I^{2 / 3}\right)$ | $\stackrel{x}{\div} \mathbf{2} .46$ | 0.053 |  |
| $\left.\frac{5_{M}}{1 / 3}\right)=2.067+0.260 \log _{10}\left(\frac{1}{W^{L / 3}}\right)$ | $\pm 0.390$ | $\frac{D_{M}}{W^{I / 3}}=117\left(\frac{1}{W^{1 / 3}}\right)^{0.260}$ | $\stackrel{\mathrm{x}}{\div} \mathbf{} 2.45$ | 0.12 | 2.9 |
| $\left.\frac{M}{I / 3}\right)=2.504-0.0150 \log _{10}\left(\frac{1}{w^{1 / 3}}\right)+0.0429\left[\log _{10}\left(\frac{1}{w^{1 / 3}}\right)\right.$ | ${ }^{2} \pm 0.391$ | $\frac{D_{M}}{W^{1 / 3}}=320\left(\frac{\left.\left.\left.\left.\frac{1}{W^{1 / 3}}\right)^{-0.0150}(10)^{0.0429}\left[\log _{10}\left(\frac{1}{W^{1 / 3}}\right)\right] .\right] .\right] .\right] . ~}{}\right.$ | $\stackrel{\mathrm{x}}{\div} \mathbf{+} 2.46$ | 0.072 | 2.10 |
| $\left.\frac{M}{1 / 3}\right)=2.382+0.178 \log _{10}\left(\frac{I}{W^{2 / 3}}\right)$ | $\pm 0.390$ | $\frac{D_{M}}{W^{I / 3}}=241\left(\frac{1}{w^{2 / 3}}\right)^{0.178}$ | $\stackrel{\mathrm{x}}{\div} \mathrm{C} 2.45$ | 0.12 |  |
| $\left.\frac{M}{13}\right)=3.369-0.452 \log _{10}\left(\frac{1}{W^{2 / 3}}\right)+0.100\left[\log _{10}\left(\frac{1}{W^{2 / 3}}\right)\right]^{2}$ | $\pm 0.390$ | $\frac{D_{M}}{W^{1 / 3}}=234\left(\frac{1}{W^{2 / 3}}\right)^{-0.453}(10)^{0.100}\left[\log _{10}\left(\frac{1}{W^{2 / 3}}\right)\right]^{2}$ | $\stackrel{x}{\div} 2.46$ | 0.081 |  |


| Factors Correlated | yp | Logarithmic form |
| :---: | :---: | :---: |
|  | Correlation | Regression Line |
| $\begin{gathered} \log _{10} \mathrm{w}^{2 / 3} \text { - Scaled Maximum Debris Distance (ft/ons } \frac{2 / 3}{\mathrm{TNT}} \text { ) } \\ \text { Versus } \\ \log _{10} \text { Equivalent Yield (tons } \mathrm{TNT}^{\text {E }} \text { ) } \end{gathered}$ | Linear | $\log _{10}\left(\frac{D_{M}}{W^{2 / 3}}\right)=2.950-0.315 \mathrm{Iog}_{10}$ |
|  | Quadrati | $\log _{10}\left(\frac{D_{M}}{W^{2 / 3}}\right)=2.960-0.320 \mathrm{Lug}_{10} \mathrm{w} \cdot 9.0151\left(\log _{10} \mathrm{~W}\right)^{2}$ |
| $\log _{10} \mathrm{w}^{2 / 3}$ - Scaled Maximum Debris Distance (ft/tons $\frac{2 / 3}{\mathrm{TNT}}$ ) | Linear | $\log _{10}\left(\frac{D_{M}}{W^{2 / 3}}\right)=4.309-0.417 \log _{1}$ |
| $\log _{10} \begin{gathered}\text { Versus } \\ \text { Impulse (lb-msac/in. }\end{gathered}$ | Quadrati | $\log _{10}\left(\frac{D_{M 1}}{W^{2 / 3}}\right)=3.432+0.0731 \log _{10^{1}} 1-0.0648\left(\log _{10} I\right)^{2}$ |
| $\log _{10} \mathrm{~W}^{2 / 3}$ - Scaled Maximum Debris Distance (ft/tons | Linear | $\log _{10}\left(\frac{D_{M}}{u^{2 / 3}}\right)=2.905-1.035 \log _{10} \mathrm{w}^{1 / 3}$ |
| Versus <br> $\log _{10}$ Cube-Root Equivalent Yield (tuns $\frac{1 / 3}{}$ ) | Quadrati | $\log _{10}\left(\frac{D_{M}}{W^{2 / 3}}\right)=2.960-0.961 \log _{10} w^{1 / 3}-1.453\left(\log _{10} w^{1 / 3}\right)^{2}$ |
| $\log _{10} \mathrm{w}^{2 / 3}$ - Scaled Maximum Debris Distance (ft/tons $\mathrm{TN}^{2 / \mathrm{T}}$ ) | Linear | $\log _{10}\left(\frac{\bar{D}}{W^{2} \sqrt{3}}\right)=4.309-1.25!1 . \int_{10} 1^{1 / 3}$ |
| $\left.\log _{10} \text { Cube-Root Impulse (lb-msec/in. }{ }^{2}\right)^{1 / 3}$ | Quadr | $\log _{10}\left(\frac{-\frac{M}{W^{2 / 3}}}{)}\right)=3.452+0.2 \cdot 9 \log _{10} I^{1 / 3}-0.585\left(\log _{10} I^{1 / 3}\right)^{2}$ |
| $\log _{10} \mathrm{w}^{2 / 3}$ - Scaled Maximum Debris Distance (ft/tons $\frac{2 / 3}{\mathrm{TNT}}$ ) | Linear | $\log _{10}\left(\frac{D_{M}}{W^{3} 3}\right)=2.950-0.517 \log _{10} \mathrm{w}^{2 / 3}$ |
| Versus $\log _{10} \frac{2}{3} \text { - Power Equivalent Yield (tons } \frac{2 / 3}{\mathrm{TNT}} \text { ) }$ | Quadrati | $\log _{10}\left(\frac{M}{W^{2 / 3}}\right)=2.960-0.480 \log _{10} w^{2 / 3}-0.0362\left(\log _{10} w^{2 / 3}\right)^{2}$ |
| $\log _{10} \mathrm{w}^{2 / 3}$ - Scaled Maximum Debris Distance (ft/tons ${ }_{\mathrm{T}}{ }_{\mathrm{NT}}{ }^{2 / 3}$ ) | Linear | $\log _{10}\left(\frac{D_{M}}{W^{2 / 3}}\right)=4.309-0.625 \log _{10} \mathrm{I}^{2 / 3}$ |
| $\left.\log _{10} \frac{2}{3} \text { - Power Impulse (lb-msec/in. }{ }^{2}\right)^{2 / 3}$ | Quadrati | $\log _{10}\left(\frac{D_{M}}{W^{2 / 3}}\right)=3.432+0.1096 \log _{10} 1^{2 / 3}-0.1457\left(\log _{10} I^{2 / 3}\right)^{2}$ |
|  | Linear | $\log _{10}\left(\frac{D_{M}}{W^{2 / 3}}\right)=6.486-1.118 \log _{10}\left(\frac{1}{W^{1 / 3}}\right)$ |
|  | Quadratic | $\log _{10}\left(\frac{\mathrm{D}}{W^{2 / 3}}\right)=-9.111+8.69 \log _{10}\left(\frac{1}{W^{1 / 3}}\right)-1.531\left[\log _{10}\left(\frac{1}{W^{1 / 3}}\right)\right]^{2}$ |
| $\log _{10} \mathrm{~W}^{2 / 3}$ - Scaled Maximum Debris Distance (ft/tons ${ }_{\mathrm{TNT}}{ }^{2 / 3}$ ) | Linear | $\log _{10}\left(\frac{D_{M}}{W^{2 / 3}}\right)=1.565+1.361 \log _{10}\left(\frac{1}{W^{2 / 3}}\right)$ |
| Versub $\left.\log _{10} \mathrm{w}^{2 / 3} \text { - Scaled Impulse (lb-msec/in. } .^{2} \text {-tons } \frac{2 / 3}{\mathrm{TNT}}\right)$ | Quadrati | $\log _{10}\left(\frac{D_{M}}{w^{2 / 3}}\right)=-3.451+2.565 \log _{10}\left(\frac{1}{w^{2 / 3}}\right)-1.910\left[\log _{10}\left(\frac{I}{W^{2 / 3}}\right)\right]$. |

Table 2. 1 (:omtinued)

| Logarithmic form |  | Exponertial Form |  | Correlation Coefficient | Figure Number |
| :---: | :---: | :---: | :---: | :---: | :---: |
| ssion Line | $\begin{gathered} \text { Standard } \\ \text { Error } \end{gathered}$ | Regression Line | Standard Error |  |  |
| $=2.950-0.34 \% \mathrm{I}^{\left(0 g_{10}\right.} \mathrm{W}$ | $\pm 0.392$ | $\frac{D_{M}}{W^{2 / 3}}=892 \mathrm{w}^{-0.345}$ | $\stackrel{\text { ¢ }}{+} \mathbf{+} 2.47$ | 0.69 |  |
| =2.960-0.320 $\left.\log _{10} \mathrm{~W}-1\right) .0191\left(\log _{10} \mathrm{~W}\right)^{2}$ | $\pm 0.392$ | $\frac{\mathrm{D}_{M}}{\mathrm{~W}^{2 / 3}}=912 \mathrm{~W}^{-0.320}(10)^{-0.0161\left(\log _{10} W\right)^{2}}$ | $\stackrel{x}{\div} 2.47$ | 0.69 |  |
| $5_{3}=4.309-0.417 \log _{15} \mathrm{y}$ | $\pm 0.407$ | $\frac{D_{M}}{W^{2 / 3}}=2.04 \times 10^{4} I^{-0.417}$ | $\stackrel{\times}{+} 2.55$ | 0.66 |  |
| $=3.432+0.0731 \log _{10} 1-0.06+0\left(\log _{10} 1\right)^{2}$ | $\pm 0.402$ | $\frac{D_{M}}{W^{2 / 3}}=2.71 \times 10^{3} I^{0.0731}(10)^{-0.0648\left(\log _{10} 1\right)^{2}}$ | $\stackrel{x}{+} 2.52$ | 0.67 |  |
| $=2.905-1.035 \log _{10} \mathrm{w}^{1 / 3}$ | $\pm 0.392$ | $\frac{D_{M}}{\mathrm{w}^{2 / 3}}=892\left(\mathrm{w}^{1 / 3}\right)^{-1.035}$ | $\stackrel{\times}{+} 2.47$ | 0.69 |  |
| $=2.960-0.961 \log _{10} \mathrm{w}^{1 / 3}-1.453\left(\log _{10} \mathrm{w}^{1 / 3}\right)^{2}$ | $\pm 0.392$ | $\left.\frac{D_{M}}{W^{2 / 3}}=912\left(\mathrm{w}^{1 / 3}\right)^{-0.961}(10)^{-1.453\left(\log _{10} W^{1 / 3}\right.}\right)^{0}$ | $\stackrel{\times}{\div} \mathbf{\square} 2.47$ | 0.69 |  |
| $=4.309-1.25!1 . c J_{10} 1^{1 / 3}$ | $\pm 0.407$ | $\frac{\bar{D}_{M}}{w^{2 / 3}}=2.04 \times 10^{4}\left(1^{1 / 3}\right)^{-1.251}$ | $\stackrel{x}{\div}$ | 0.66 |  |
| $=3.452+0.2 \cdot 9 \log _{10} \mathrm{I}^{1 / 3}-0.585\left(\log _{10} \mathrm{I}^{1 / 3}\right)^{2}$ | $\pm 0.402$ | $\frac{\bar{D}_{M}}{W^{2 / 3}}=2.71 \times 10^{3}\left(\mathrm{I}^{1 / 3}\right)^{0.219}(10)^{-0.585\left(\log _{10}\right.}$ | $\stackrel{\mathrm{x}}{\div} \mathbf{-} 2.52$ | 0.67 |  |
| $=2.950-0.517 \log _{10} \mathrm{w}^{2 / 3}$ | $\pm 0.392$ | $\frac{D_{M}}{W^{2 / 3}}=892\left(\mathrm{w}^{2 / 3}\right)^{-0.517}$ | $\stackrel{x}{\square} 2.47$ | 0.69 |  |
| $=2.960-0.480 \mathrm{Lotg}_{10} \mathrm{w}^{2 / 3}-0.0362\left(\log _{10} \mathrm{w}^{2 / 3}\right)^{2}$ | $\pm 0.392$ | $\left.\frac{\mathrm{D}_{\mathrm{M}}}{\mathrm{w}^{2 / 3}}=912\left(\mathrm{w}^{2 / 3}\right)^{-0.480}(10)^{-0.0362\left(\log _{10}\right.} \mathrm{w}^{2 / 3}\right)^{2}$ | $\stackrel{x}{\div} \mathbf{+} 2.47$ | 0.69 |  |
| $\overline{3})=4.309-0.625 \log _{10} \mathrm{I}^{2 / 3}$ | $\pm 0.407$ | $\frac{D_{M}}{W^{2 / 3}}=2.04 \times 10^{4}\left(\mathrm{I}^{2 / 3}\right)^{-0.625}$ | $\stackrel{\mathrm{x}}{\div} \mathbf{\square} 2.55$ | 0.66 |  |
| $\begin{equation*} \overline{3}=3.432 \div 0.1096 \log _{10} 1^{2 / 3}-0.1457\left(\log _{10} \mathrm{I}^{2 / 3}\right)^{2} \tag{10} \end{equation*}$ | $\pm 0.402$ | $\frac{\mathrm{D}_{\mathrm{M}}}{\mathrm{w}^{2 / 3}}=2.71 \times 10^{3}\left(\mathrm{I}^{2 / 3^{0.1096}}\right.$ | $\stackrel{\mathrm{x}}{\div} \mathrm{\square} 2.52$ | 0.67 |  |
| $\left.\frac{T}{3}\right)=6.486-1.118 \log _{10}\left(\frac{I}{W^{1 / 3}}\right)$ | $\pm 0.504$ | $\frac{\mathrm{D}_{\mathrm{M}}}{\mathrm{w}^{2 / 3}}=3.06 \times 10^{6}\left(\frac{\mathrm{I}}{\mathrm{w}^{1 / 3}}\right)^{-1.118}$ | $\stackrel{\mathrm{x}}{\div} 3.19$ | 0.38 |  |
| $\left.\frac{1}{3}\right)=-9.111+8.69 \log _{10}\left(\frac{1}{w^{1 / 3}}\right)-1.531\left[\log _{10}\left(\frac{1}{w^{1 / 3}}\right)\right]^{2}$ | $\pm 0.497$ | $\frac{\bar{D}_{M}}{W^{2 / 3}}=7.74 \times 10^{-10}\left(\frac{1}{W^{1 / 3}}\right)^{8.69}(10)^{-1.531}\left[\log _{10}\left(\frac{1}{W^{1 / 3}}\right)\right]$ | $\stackrel{\mathrm{x}}{\div} 3.14$ | 0.40 |  |
| $\left.r_{3}\right)=1.565+1.361 \log _{10}\left(\frac{I}{W^{2 / 3}}\right.$ | $\pm 0.403$ | $\frac{D_{M}}{\mathrm{w}^{2 / 3}}=2.72 \times 10^{-2}\left(\frac{\mathrm{I}}{\mathrm{w}^{2 / 3}}\right)^{1.361}$ | $\stackrel{x}{\div} 2.53$ | 0.67 |  |
| $=-3.451+2.565 \log _{10}\left(\frac{I}{W^{2 / 3}}\right)-1.910\left[\log _{10}\left(\frac{I}{W^{2 / 3}}\right)\right]^{1} \pm 0.403$ |  | $\frac{D_{M}}{W^{2 / 3}}=3.54 \times 10^{-4}\left(\frac{I}{w^{2 / 3}}\right)^{2.57}(10)^{-1.910}\left[\log _{10}\left(\frac{1}{W^{2 / 3}}\right)\right]^{2}$ | $\stackrel{x}{+} 2.53$ | 0.67 |  |





Figure 2. 4 Quadratic Regression Line: Maximum Debris Distance versus Equivalent Yield

n Debris Distance









Figtife 2.7 Linear Regzession Line: Maximum Debris Distance versus !mpulse

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Figure 2.8 Quadratic Regression Line: Maximum Debis Distance versus Impulse

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Figure 2. 10 Quadratic Regression Line: $\mathrm{W}^{1 / 3}$-Scaled Maximum Debris Distance versus $\mathrm{W}^{1 / 3}$-Scaled Impulse

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Maximum debris distance data from the JANGLE $U$ shot and the DANNY BOY event are included in Fig. 2. 3 to 2.6 to show their consistency with the HE data. Maximum debris distance for the DANNY BOY event was taken at 750 ft , the distance to a lobe of detached dust observed in post-shot aerial photographic observation (Ref. 9). The DANNY BOY event involved a $0.430-\mathrm{KT}$ device buried at 110 ft . Maximum debris distance for JANGLE $U$ was taken at 550 ft -- the farthest-thrown recorded fragment of the reinforced concrete runways and wall panels built within the crater zone. The JANGLE U shot involved a $1.2-\mathrm{KI}$ device buried at 17 ft . These points are consistent with the HE results, especially when considered in relation to the regression line and the one-standard-error limit. The DANNY BOY event data compare less favorably with the HE regression lines than do the JANGLE U findings. The much greater depth of burial in DANNY BOY, where the device was placed at "optimum" depth, accounts for this. The resultant trajectories observed in DANNY BOY had pronounced vertical components. The favorable comparison of the underground nuclear events with the HE results is not surprising. Because of the absence of extreme blast winds, the HE detonations may be more closely akin to the underground nuclear burst than to surface or above-grade nuclear bursts, as far as debris behavior is concerned.

Lack of similarity between contained HE explosions and nuclear explosions under various siting and target conditions is obvious. The application of these curves to specific targeting situations can certainly be questioned. The dearth of actual debris measurements from full-scale nuclear tests is also realized. In view of these factors, the necessity of using available data such as this, tempered by engineering judgment, is apparent. Use of these charts will permit the making of order-of-magnitude estimates of maximum debris distance. Such estimates can at least indicate whether debris is a problem under various targeting conditions. A first-order estimate of maximum debris range measured from ground zero can be obtained from the correlations of maximum debris distance and equivalent yield. The correlations of maximum debris distance with impulse can be applied to the problem of an individual structure located intermediate between ground zero and the target whose vulnerability is questioned. Consideration of the

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standard error provides an indication of the likelihood of the target being within maximum debris distance.

These charts, in themselves, do not provide any means for estimating the probability of the target being hit by debris, or estimating the number of pieces of debris likely to hit the target. They estimate maximum debris distance only-- the location of the furthest-thrown fragment. The problem of areal density of debris (the number of fragments per unit area) will be considered separately.

Considerable difficulty was experienced in keeping within the accurate capacity of the electronic computer in these regression studies. The matrix solution to the simultaneous equations for the least-squares lines involves taking small differences between very large numbers which are summations of products or summations of powers of numbers. As these summations increase in size, the process taxes the accurate capacity of the computer. Although where necessary, recourse was made to double precision ( 16 digits) on the IBM $7090 / 401$ computer, it was still not possible to ascribe a high degree of mathematical precision to the constants in the regression results. Thus we see that mathematically the correlation coefficient for the quadratic case should not be lower than that in the linear case, or else the quadratic expression should have a zero coefficient to the squared independent variable. Regardless of this, we have observed that computed values of correlation coefficients were sometimes lower for the quadratic case, and that computed values of standard errors were frequently higher for the quadratic case. In these instances the differences in values of correlation coefficients and standard errors for the linear and quadratic cases were small -- actually marginal. The first digit of the standard error and correlation coefficient is good, the second digit probably merits some confidence. The constants in the regression lines are of unknown precision. The relative position of the lines among the scattered data points and the ratio of numbers of data points within and without the plus-and-minus-one standard error bands are reasonable.

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## 2. 2. 3 Correlation of Maximum Debris Distance with Crater Dimensions

Maximum debris distance was correlated with crater dimensions using data from the series of thirty-six explosions tabulated in Appendix $C$ (Ref. 6). Results of the regression study are tabulated in Table 2. 2. Linear and quadratic plots of maximum missile distance in terms of crater dimensions are presented in Fig. 2. 11 through 2.16 , while $W^{1 / 3}$-scaled relationships are presented in Fig. 2. 17 through 2. 22. JANGLE U AND DANNY BOY results are also plotted on these figures. Again the considerable depth of burial in the DANNY BOY event accounts for the position of these points well below the lower one-standard-error band.

## 2. 3 Debris Dispersion Patterns

Theoretical models for distribution of debris about line charges are developed and compared with actual results from semi-contained explosions. Measured debris dispersion from HE detonations is studied. Comparisons of dispersion patterns for a series of six HE events of different magnitudes are made. Dispersion patterns for a single HE event are also studied in considerable detail.

### 2.3.1 Theoretical Models

## Debris Dispersion for Line-Source Explosions

Theoretical models of the dispersion of structural fragments in terms of the maximum missile distance were derived for flat roots, walls, and arches. These models were prepared for the purpose of comparing theoretical distributions with those from experimental and accidental explosions.

Models developed to explain the dispersion of structural fragments, in terms of terminal ground range, are based on the following assumptions:

Complete fragmentation of the structure
into equal-sized fragments occurs
Fragments follow ballistic trajectories, and
air resistance can be neglected

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Figure 2.13 Linear Regression Line: Maximum Debris Distance versus Crater Diameter


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Figurc $2.14 \frac{\text { Quadiatıc Regression Line: Maximum Debris Distance }}{\text { versus Crate = Diameter }}$
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Figure 2.15 Linear Regression Line: Maximum Debris Distance versus Crater Depth
(1y) a.mpisid kisqad unumerw

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| Fectors Correlated | $\begin{gathered} \text { Type } \\ \text { of } \\ \text { Corytelation } \end{gathered}$ | Lokarlil:, Form <br> Rogression Line | E |
| :---: | :---: | :---: | :---: |
| $\begin{gathered} \log _{10} \begin{array}{c} \text { Maximum Debris Distance (ft) } \\ \text { Versums } \end{array} \\ \log _{10} \text { Crater Volume }\left(\mathrm{ft}^{3}\right) \\ \hline \end{gathered}$ | Linear | $\mathrm{Log}_{10} \mathrm{D}_{\mathrm{MM}}=\therefore .187 \div 0.262 \mathrm{Log}_{10}$ |  |
|  | Quadratic | $\log _{10} \mathrm{D}_{\mathrm{M}}=3.391-0.389 \log _{16}{ }^{\text {c }}$ c $+0.0835\left(\log _{10} \mathrm{~V}_{6}\right)^{\prime}$ |  |
| $\log _{10} \underset{\text { Versus }}{\text { Maximum Dis }}$ Debris ( ft ) $\log _{10}$ Crater Diameter (ft) | L.inear | $\log _{10} D_{M}=2.261+0.577 \log _{10} 0^{\circ}$ |  |
|  | Quadratic | $\log _{10} \mathrm{D}_{M}=2.946-0.258 \log _{10} 0+0.245\left(\log _{10} \phi\right)^{2}$ |  |
| $\log _{10}$ Maxımum Debris Distance (ft) <br> Versus <br> $\log _{10}$ Crater Depth ( ft ) | Linear | $\mathrm{Los}_{10} \mathrm{D}_{\mathrm{L} 1} \times 2.428+0.770 \mathrm{Log}_{10} \mathrm{D}_{5}$ |  |
|  | Quadratic |  |  |
| $\begin{aligned} & \log _{10} \mathrm{~W}^{1 / 3} \text { - Scaled Maximum Debris Distance }\left(\frac{\mathrm{ft}}{\mathrm{tons}^{\mathrm{TNT}}} \mathrm{~V}{ }^{2 / 3}\right) \\ & \quad \text { Versus } \end{aligned}$ | Lincar | $\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.822+0.0467 \mathrm{Len}_{10} V^{\prime}$ | $\pm$ |
|  | Quadratic | $\log _{10}\left(\frac{D_{M 1}}{\chi^{1 / 3}}\right)=3.867-0.012 \log _{10}{ }^{\prime}{ }^{(1)}+6.072+\left(\log _{10} \mathrm{~V}_{\mathrm{c}}\right)^{2}$ | $\pm$ |
| $\left.\begin{array}{l} \log _{10} \mathrm{w}^{1 / 3}-\text { Scaled Maximum Debris Distance }\left(\frac{\mathrm{tt}}{\mathrm{tons}_{\mathrm{TN}}} 1 / 3\right) \\ \quad \text { Versus } \end{array}\right)$ | Linear | $\log _{10}\left(\frac{5{ }^{1}}{1 / 3}\right)=3.026-0.0122 \log _{10} \prime$ | + |
|  | Quadratic | $\log _{10}\left(\frac{\sigma^{M}}{1 / 3}\right)=3.526-0.6 \therefore:\left(\log _{10} \cdot 1+0.179\left(\log _{10} \not\right)^{2}\right.$ | $\pm$ |
| $\begin{aligned} & \log _{10} \mathrm{w}^{1 / 3} \text { - Scaled Maximum Debris Distance }\left(\frac{\mathrm{ft}}{\left.\operatorname{tons}_{\mathrm{TNT}}{ }^{1 / 3}\right)}\right. \\ & \quad \text { Versus } \\ & \log _{10} \text { Crater Depth (ft) } \end{aligned}$ | Linear |  | $\pm$ |
|  | Quadratic | $\log _{10}\left(\frac{D}{W^{1 / 3}}\right)=3.273-1.077 \mathrm{~L}^{1} \varepsilon_{10} i_{c}+0.827\left(\log _{10} \mathrm{D}_{\mathrm{c}}\right)^{2}$ | $\pm$ |
| $\begin{gathered} \log _{10} \mathrm{w}^{1 / 3}-\text { Scaled Maxinum Debris Distance }\left(\frac{\mathrm{ft}}{\text { tons } \mathrm{TNT}^{1 / 3}}\right) \\ \text { Versus } \\ \log _{10} \mathrm{w}^{1 / 3}-\text { Scaled Crater Volume }\left(\frac{\mathrm{ft}^{3}}{\text { tons }_{\mathrm{TNT}}}\right) \end{gathered}$ | Linear | $\log _{10}\left(\frac{\mathrm{D}^{\mathrm{M}}}{\mathrm{W} / 3}\right)=2.569+11.11 ; \log _{10}\left(\frac{\because}{\mathrm{w}^{1 / 3}}\right)$ | $\pm$ |
|  | Quadratic | $\log _{10}\left(\frac{D_{M}}{\mathbf{w}^{I / 3}}\right)=3.379-0.339 \log _{10}\left(\frac{c}{w^{1 / 3}}\right)+0.0618\left[\log _{10}\left(\frac{v_{c}}{w^{1 / 3}}\right)\right]^{2}$ | $\pm$ |
| $\begin{aligned} & \log _{10} \mathrm{w}^{1 / 3}-\text { Scaled Maximum Debris Distance }\left(\frac{\mathrm{ft}}{\operatorname{Cons} \mathrm{TNT}^{1 / 3}}\right) \\ & \log _{10} \mathrm{w}^{1 / 3} \text { - Scaled Crater Diameter }\left(\frac{\mathrm{ft}}{\text { tons }_{\mathrm{TNT}}} \mathrm{~T}^{1 / 3}:\right. \end{aligned}$ | Linear | $\left.\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.450+0.385 \log _{10} \frac{\phi}{w^{1 / 3}}\right)$ | $\pm$ |
|  | Quadratic | $\log _{10}\left(\frac{D_{M}}{\frac{w^{1 / 3}}{1 / 3}}\right)=0.366+2.64+\operatorname{Lo}_{10}\left(\frac{6}{w^{i / 3}}-0.781\left[\log _{10} \frac{\phi}{w^{1 / 3}}\right]^{2}\right.$ | + |
| $\begin{aligned} & \log _{10} \mathrm{w}^{1 / 3} \text { - Scaled Maximum Debris Distance }\left(\frac{\mathrm{ft}}{\operatorname{tons} \mathrm{TNT}^{1 / 3}}\right) \\ & \quad \text { Versus } \\ & \log _{10} \mathrm{w}^{1 / 2}-\text { Scalcd Crater Depth }\left(\frac{\mathrm{ft}}{\operatorname{tons} \mathrm{TNT}^{1 / 3}}\right. \end{aligned}$ | Linear | $\left.\log _{10}\left(\frac{D_{M}}{W^{1 / 3}}\right)=2.506+0.610 \log _{10} \frac{D}{w^{1 / 3}}\right)$ | $\pm$ |
|  | Quadratic | $\log _{10}\left(\frac{\mathrm{D}_{\mathrm{M}}}{\mathrm{w}^{1 / 3}}\right)=2.299+1.164 \log _{10}\left(\frac{\mathrm{D}}{\mathrm{w}^{1 / \%}} \cdot 0.313\left[\log _{10} \frac{\mathrm{D}_{\mathrm{c}}}{\mathrm{w}^{1 / 3}}\right)^{2}\right.$ | + |



Figure 2.17 Linear Regression Linc:

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Figure 2.18 Quadratic Regression Line: $w^{1 / 3}$-Scaled Maximum Debris Distance


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Figure 221 Linear Regression Line: $W^{1 / 3}$-Scaled Maximurn Debris Distance versus $\mathrm{W}^{1 / 3}$-Scaled Crater Depth

Figure 2. $22 \frac{\text { Quadratic Regression Line: } W^{1 / 3} \text {-Scaled Míaxiinum Debris Distance }}{\text { (/3 }}$
versus $\mathrm{W}^{1 / 3}$-Scaled Crater Depth


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All fragments have equal initial velocities, though other assumptions as to the initial velocity field may be made.

These assumptions are obviously not true, but refined expressions either cannot be obtained for the general case or would unduly complicate the model. These assumptions are justified in development of a model for order-of-magnitude estimates of the relative densities of debris at various ground ranges.

Figure 2. 23 represents a wall subjected to an explosive impulse from a line-source charge.


Figure 2. 23 Wall Panel under Explosive Impulse from Line Source Charge

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In this figure,

$$
\left.\begin{array}{l}
\mathrm{r}_{1}=\frac{\mathrm{r}}{\cos \theta} \\
\mathrm{dx}=\frac{\mathrm{r}_{1} \mathrm{~d} \theta}{\cos \theta}=\frac{\mathrm{rd} \mathrm{\theta}}{\cos ^{2} \theta}  \tag{2,3}\\
\frac{d x}{d \theta}=\frac{\mathrm{r}}{\cos ^{2} \theta} \text { in which the limit on } \theta \leqslant \tan ^{-1} \frac{\mathrm{~h}^{\prime}}{\mathrm{r}},
\end{array}\right\}
$$

and the fragment distribution in the wall becomes:

$$
\frac{\text { Number of fragments }}{\text { Unit elevation angle }} \propto \frac{1}{\cos ^{2} \theta} \text { with limit } \theta \leqslant \tan ^{-1} \frac{h^{\prime}}{r}
$$

Similarly, the geometry of the roof is shown in Fig. 2. 24.


Figure 2. 24 Roof Panel Fragmented by Line-Source Charge

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Here,

$$
\begin{align*}
& d x=\frac{r_{2} d \theta}{\cos \theta} \\
& r_{2}=\frac{h^{\prime}}{\sin \theta}  \tag{2.4}\\
& \frac{d x}{d \theta}=\frac{h^{\prime}}{\sin \theta \cos \theta} \text { in which the limit on } \theta \geqslant \tan ^{-1} \frac{h^{\prime}}{r}
\end{align*}
$$

and the fragment distribution in the roof becomes
$\frac{\text { Number of fragments }}{\text { Unit elevation angle }} \propto \frac{\mathrm{J}}{\sin \theta \cos \theta}$ with limit $\theta \geqslant \tan ^{-1} \frac{\mathrm{~h}^{\prime}}{\mathrm{r}}$.

For the circular arch, shown in Fig. 2. 25,


Fig. 2. 25 Structural Arch Subject to Impulse Load From
Line-Source Charge
the distribution of fragment quantities becomes a constant.
$\frac{\text { Number of fragments }}{\text { Unit elevation angle }}=\frac{\mathrm{dx}}{\mathrm{d} \theta}=\mathrm{r}_{3}=$ constant
The distance traveled by the individual fragments is represented in Fig. 2. 26,


Figure 2. 26 Trajectory of Individual Fragment

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in which,

$$
\begin{aligned}
& v_{v}=v_{o} \sin \theta \\
& \mathrm{~h}=\mathrm{v}_{\mathrm{o}} \sin \theta \mathrm{t}-\frac{1}{2} \mathrm{gt}^{2}
\end{aligned}
$$

and

$$
h=\frac{1}{2} g t^{2}
$$

Therefore,

$$
\begin{aligned}
& \mathrm{v}_{\mathrm{o}} \sin \theta \mathrm{t}=\mathrm{gt}^{2} \\
& \mathrm{t}=\frac{\mathrm{v}_{\mathrm{o}} \sin \theta}{\mathrm{~g}}
\end{aligned}
$$

Also,

$$
\mathrm{v}_{\mathrm{h}}=\mathrm{v}_{\mathrm{o}} \cos \theta
$$

Therefore

$$
\begin{aligned}
\mathrm{d} & =\mathrm{v}_{\mathrm{o}} \cos \theta \mathrm{t} \\
& =\frac{\mathrm{v}_{0}^{2} \sin \theta \cos \theta}{\mathrm{~g}},
\end{aligned}
$$

and the trajectory distance parameter becomes

$$
\mathrm{d} \propto \sin \theta \cos \theta
$$

The fragment quantity parameters, $\frac{1}{\cos ^{2} \theta}$ and $\frac{1}{\sin \theta \cos \theta}$, are plotted
against the fragment distance parameter $\sin \theta \cos \theta$ in Fig. 2.27 for values of $\theta$ from $0^{\circ}$ to $90^{\circ}$. The relative quantity of fragments landing within any circumferential sector can be found by integrating these functions over the appropriate limits of the function $\sin \theta \cos \theta$.
This integration is carried out as follows for walls:
In Fig. 2. 27 let

$$
\begin{aligned}
& y=\frac{1}{\cos ^{2} \theta} \\
& x=\sin \theta \cos \theta
\end{aligned}
$$

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Figure 2.27 Theoretical Fragment Distribution Functions

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Then

$$
\begin{aligned}
d \mathbf{x} & =\frac{d x}{d \theta} d \theta \\
& =\left(-\sin ^{2} \theta+\cos ^{2} \theta\right) d \theta \\
& =\left(1-2 \sin ^{2} \theta\right) d \theta
\end{aligned}
$$

The relative quantity of fragments in strips centered about the line charge, (or circumferential strips centered about ground zero), then becomes:

$$
\begin{aligned}
Q_{w} & =\int_{\theta_{1}} y d x \\
& =\int_{\theta_{2}}^{\theta_{2}} \frac{\left(1-2 \sin ^{2} \theta\right)}{\cos ^{2} \theta} d \theta \\
& =[\tan \theta-2(\tan \theta-\theta)]_{\theta_{1}}^{\theta_{2}} \\
& =[20-\tan \theta]_{\theta_{1}}^{\theta_{2}} \quad \text { with limit on } \theta \leqslant \tan ^{-1} \frac{h^{\prime}}{r} .
\end{aligned}
$$

Similarly, for roofs, let

$$
\begin{aligned}
& \mathrm{y}=\frac{1}{\sin \theta \cos \theta} \\
& \mathrm{x}=\sin \theta \cos \theta \\
& \frac{d x}{\mathrm{~d} 0}=\left(1-2 \sin ^{2} \theta\right) \mathrm{d} \theta
\end{aligned}
$$

thus

$$
\begin{aligned}
\theta_{R} & =\int_{\theta_{1}} y d x \\
& =\int_{2} \frac{\left(1-2 \sin ^{2} \theta\right)}{\sin \theta \cos \theta} d \theta \\
& =\left[\log _{\epsilon} \tan \theta-2\left(-\log _{\epsilon} \cos \theta\right)\right]_{\theta_{1}}^{\theta_{2}} \\
& =\left[\log _{\epsilon} \tan \theta+2 \log _{\epsilon} \cos \theta\right]_{\theta_{1}}^{\theta_{2}} \text { in which limit } \theta \geqslant \tan ^{-1} \frac{h^{\prime}}{r} .
\end{aligned}
$$

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For arches:

$$
\begin{aligned}
y & =\operatorname{constant}=k \\
x & =\sin \theta \cos \theta \\
\frac{d x}{d \theta} & =\left(1-2 \sin ^{2} \theta\right) d \theta
\end{aligned}
$$

thus integration for arches yields a fragment quantity function as follows:

$$
\begin{aligned}
Q_{A} & =\int k\left(1-2 \sin ^{2} \theta\right) d \theta \\
& =\left[k \theta-2 k\left(\frac{\theta}{2}-\frac{\sin 2 \theta}{4}\right)\right]_{\theta_{1}}^{\theta_{2}} \\
& =\left[\frac{k \sin 2 \theta}{2}\right]_{\theta_{1}}^{\theta_{2}} .
\end{aligned}
$$

Relative fragment quantities for circumferential bands about the point of burst, (assumed here to be at zero elevation at the center of the containing structure), are obtained by equating $Q$ over the appropriate values of $\theta_{1}$ and $\theta_{2}$. For wall panels $\theta_{1}$ is $\theta^{\circ}$ and $\theta_{2}$ the angle included between surface zero and the plane through the line charge and the wall-roof intersection. For actual structures $Q$ must be corrected to account for the different material quantities in wall and roof panels, and for differentials between normal distances between the line charge and the wall and roof panels. Thus, for the structure pictured in Fig. 2. 28, the relative number of fragments falling within any band $R_{j}-R_{i}$ about the structure would be


Figure 2. 28 Typical Structure Subjected to Fragmentation by Internal Detonation

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$$
\begin{aligned}
Q_{T} & =t_{w} r Q_{w}+t_{R} h Q_{R} \\
& =t_{w} r[2 \theta-\tan \theta]_{0^{\circ}}^{\theta} A+t_{R} h\left[\log _{\epsilon} \tan \theta+2 \log _{\epsilon} \cos \theta\right]_{\theta_{A}}^{90^{\circ}}
\end{aligned}
$$

where

$$
\begin{aligned}
\mathrm{t}_{\mathrm{w}} & =\text { wall thickness } \\
\mathrm{t}_{\mathrm{r}} & =\text { roof panel thickness }
\end{aligned}
$$

## Comparison of Ideal and Actual Debris Distribution

 from Semi-Contained ExplosionsA comparison of the ideal fragment dispersion functions with three actual HE explosions is made in Fig. 2. 29. In making this comparison the quantities $t_{w}, s, t_{r}$, and $h$ were dropped because of lack of data fully describing the containing structures. Furthermore, the functions $Q_{r}$ and $Q_{w}$ in Fig. 2. 29 are plotted for all values from $0^{\circ}$ to $90^{\circ}$, since the actual value of $\theta_{A}$ was not known. Actually, this means that $Q_{r}$ and $Q_{w}$ are ploited for panels of infinite length. The theoretical curves $Q_{r}$ and $Q_{w}$ were plotted by equating ( $\sin \theta \cos \theta)_{\max }=1.0$, computing $Q_{r}$ and $Q_{w}$ for 10 percent increments of $(\sin \theta \cos \theta)_{\text {max }}$, and plotting values of $Q_{r}$ and $Q_{w}$ at the midpoints of these sectors. Curves for the actual explosions were established by equating the maximum debris distance to 1.0 , and plotting actual reported fragment counts for 10 percent increments of the maximum reported distance at the midpoints of the circumferential bands. The following explosions are included in this analysis.

Wolf Creek 1944: (Ref. 10)
Accidental Explosion.
$4,800 \mathrm{lb}$ of ammonium nitrate exploded in a building.
305 fragments were recorded within a maximum range of one mile.
Kankakee, 1943: (Ref. 10)
Accidental Explosion.
$5,300 \mathrm{lb}$ of Bi -Oil exploded in a building.
160 fragments were recorded within a maximum range of $3,750 \mathrm{ft}$.

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Figure 2.29 Comparison of Theoretical and Actual Fragment Dispersion Patterns

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> Pantex Ordnance Plant, 1960: (Ref. 11)
> Experimental explosion.
> 2, 000 lb of HE in encased warheads detonated inside U-shaped bay of reinforced concrete ordnance structure having one-foot thick walls.
> About 31,000 concrete fragments, with a total weight of about $85,000 \mathrm{lb}$, were recorded within a maximum range of about 1,450 feet.

The plotting of data from these explosions assumed that the furthest-thrown fragments were found and recorded. This is generally a reasonable assumption in reports on accidental explosions since the furthest-thrown fragments are generally of sufficient size to be observed and found, and since, in the thickly strewn close-in region, only large fragments or fragments considered significant in determining the cause of the explosion are recorded.

On the general shapes of the actual and theoretical curves are sımilar; this similarity is quite pronounced for the Pantex explosion. The average curve of fragment distribution for the Wolf Creek and Kankakee explosions is observed to fall between the theoretical and actual curves for walis and roofs -- a reasonable expectation. The steeper rise of the Pantex curve below a range of $0.8-0.9$ times maximum may result from the nature of the structure.

It would be desirable to use these curves along with the maximum debris distance curves presented earlier to develop areal density patterns or probability of hit. To do so would require that both (1) the number of fragments in the outermost circumferential sector and (2) the nature of the curve in the region where it starts rising rapidly, be defined better. These factors may be related to the size and strength of the structure and their determination would thus require extensive experimental investigation. The currently available data, considered above, could be used for this purpose although it should be recognized that they are far from well defined experimental input.

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## 2. 3. 2 Debris Dispersion Patterns of HE Explosions

Several prior studies of debris dispersion about accidental and experimental explosions have been made. One was performed by Colonel Clark S. Robinson who studied the results of six ordnance explosions, then compiled and compared their results (Ref. 12). This study is reviewed here along with additional plots of the recorded data which were prepared in an effort to find greater consistency of the plotted patterns. The second study was made at the Pantex Ordnance Plant at Amarillo, Texas, where about 31,000 fragments totaling $85,000 \mathrm{lb}$ from a reinforced concrete structure were located and weighed after a planned interior explosion of $2,000 \mathrm{lb}$ of high explosives (Ref. ll). The raw data available from this study were used in making a detailed study of the debris dispersion pattern from this explosion.

## Army-Navy Explosives Safety Board Investigation of Debris Dispersion

Colonel Robinson plotted the specific area per missile against ground range. Specifically, the following explosions were studied:

Badger Ordnance Works, 1945:
7,500 lb of nitroglycerine exploded accidentally in a barricaded storehouse building of light frame construction. Data on about 600 fragments from building and contents were recorded.

Cornhusker, Nebraska, 1945:
$10,000 \mathrm{lb}$ of explosives detonated accidentally in a loading plant manufacturing bombs. Chief interest in the investigation centered about fragments of machinery and contents and few building fragments were included.

Portage, Ohio, 1943:
Bombs amounting to $40,000 \mathrm{lb}$ of HE detonated accidentally in an arch-type, earth-covered, igloo magazine. Data on whole bombs and bomb fragments constituted the bulk of recorded debris.

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Umatilla, Oregon, 1944:
$64,000 \mathrm{lb}$ of HE in bombs exploded accidentally in an Army-type igloo. Debris described consisted chiefly of pieces of concrete and fragments of machinery stored in the igloo.

Hastings, Nebraska, 1944:
$100,000 \mathrm{lb}$ of HE ignited in a building where bombs were being loaded. Interest lay chiefly in large concrete fragments from the builidng itself, and only information on those weighing over 25 lb was included.

Arco, Idaho, 1945:
$250,000 \mathrm{lb}$ of explosive in bombs was detonated in a planned explosion in an igloo-type magazine. Data on over 13, 000 fragments were recorded,

Specific area is plotted against ground range for these six explosions in Fig. 2. 30. Several limitations of these plots, stemming from existent debris recording practices, must be recognized. First, the investigators made no attempt to record all fragments, especially in the first five explosions listed above; this tends to give higher values for specific area than actual, perhaps by several orders of magnitude. Secondly, in the thickly strewn close-in region, explosion investigators tend to record the exceptional fragments only; hence the reversal of several curves at short ground ranges in Fig. 2. 30 is unrealistic. It is stated in the ASESB report, however, that at the greater ground ranges, it was customary for investigators to record all missiles of sizes that would damage structures or kill personnel -- that is, all fragments weighing one pound or more. lt was observed that in these explosions the far-flying fragments were usually large and most were discovered and recorded.

The individual curves of specific area in Fig. 2. 30 do exhibit a consistent and characteristic shape. The relative positions and crossing over of the four intermediate explosions ( $10,000-1 \mathrm{l}, 40,000-1 \mathrm{l}, 65,000-1 \mathrm{~b}$ and 100,000-1b curves) make it difficult to derive a general expression for the functional reltaionship between specific area and ground range for contained explosions. This stems from the fact that the curves are based
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Figuse 230 Debris Dispe?sion Patterns for Selected Explosions
on incomplete fragment counts, and that the actual specific area is a function of the nature of the structure itself, as well as the explosive weight and ground range.

Several additional plots of the data from Fig. 2. 30 have been prepared in an effort to find general curves for these functions with a better $\mathrm{f}_{1}$. These additional plots, shown in Fig. 2. 31 through 2. 36, were derived by selecting five points on each of the explosion curves of Fig. 2. 30 and computing new expressions as follows:

Fig. 2. 31 First-Order Logarithmic Curves of Debris Dispersion for Selected Explosions.

Log-log curves were fitted through the five selected points by the least-squares method, producing a series of curves of the form:

$$
\log _{10} A_{D}=k\left[\log _{10} D\right]^{n}
$$

No consistent relationship between either the slopes (II) or the intercepts (k) and the equivalent yield are apparent from these plots.

Fig. 2. 32 Second-Order Logarithmic Curves of Debris Dispersion for Selected Explosions.

Log-log second-order curves were fitted through the same five points as above for each curve. Distinct separations between the curves for the various yields are absent, and the negative slopes near ground zero are questionable.

Fig. 2. 33 Modified Second-Order Logarithmic Curves of Debris Dispersion for Selected Explosions.

Log-log second-order curves were fitted, again by the least-squares method, through six points for each curve -- the five used previously and the arbitrary addition of point $(1,1)$ as a data point. The resulting curves exhibit some separation at the shorter ground ranges, but the crossing-over of curves and their relative positions cannot be explained.

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Figuee 2. 31 First-Orde:-Logainthmic Curves of Debris Dispersion for Selected Explosions


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Figure 2.33 Modified Second-Order: Logarıthmic Curves of Debris Dispersion for Selected Explosions

Fig. 2. 34 Parabolic Curves of Debris Dispersion for Selected Explosions.
These curves are parabolas of the form

$$
A_{D}=k D^{2}
$$

where, for each curve, $k$ is the average value for a series of parabolas through point $(0,0)$ and each of the five selected data points. The slope appears adequate for some curves (notably the $7,500-\mathrm{lb}$ and $250,000-\mathrm{lb}$ explosions), but the relative positions of the curves cannot be explained in terms of equivalent yield alone. Total quantity of available material for fragmentation probably is involved.

Fig. 2. 35 Logarithmic Parabolas of Debris Dispersion for Selected Explosions.
These curves are of the form

$$
\log _{10} A_{D}=k\left(\log _{10} D\right)^{2}
$$

where, for each curve, $k$ is the average value for a series of parabolas through point ( 1,1 ) and each of the five points on the curves.

Fig. 2. 36 Second-Order Semi-Logarithmic Curves of Debris Dispersion for Selected Explosions.
These curves are of the form

$$
\log _{10} A_{D}=a+b D+c D^{2}
$$

where the coefficients $a, b$, and $c$ are determined by the least-squares method.

Of the six sets of curves, the logarithmic parabolas appear to provide the best fit. Relative positions of the curves cannot be explained, $i_{\text {. }}$ e., why they should not be consecutive in the order of yield. As stated earlier this may be a function of the amount of material available for fragmentation and the degree of fragmentation, which can be expected to vary with the quantity of explosives or the impulse. Because of the lack of consistency in these results, no further analysis was made of these results.

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Figure 2.34 Parabolic Curves of Debris Dispersion for Selected Explosions

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Figure 2.35 Logarithmic Parabolas of Debrıs Dispersion for Selected Explosions

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Figure 2.36 Second-Order Semi-Lugarithmic Curves of Detiis Dispersion for Seiected Explosions

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It will be recalled that investigators of accidental explosions tend to make incomplete counts of the smaller fragments in the thickly strewn close-in region about the source. Since plotted points in the foregoing study were all taken at equal value, this would result in curves which are too shallow, and thus tend to overestimate the debris problem when exirapolations are made beyond the limits of the data.

One measure of the fit of the fragment dispersion function can be obtained visually from data on the 1945 Badger Ordnance Works explosion of $7,500 \mathrm{lb}$ of nitroglycerine from the plot of data points as shown in Fig. 2. 37. Plotted points in Fig. 2. 37 show debris dispersion expressed as specific area for each $10-\mathrm{ft}$ interval in ground range from the point of burst. The average line drawn through these points was not originally computed by the least-squares method, but was an average based on the four central curves of Fig. 2. 30 plotted at one-half the ground range. Though this original method of plotting was used merely because the Badger explosion was a barricaded structure, the line does appear to be a fajr representation of this particular set of data.

The two curves denoting the one-standard-error limits have been plotted visually to include two-thirds of all plotted data points between them. The log value of the $e_{x}$ standard error is thus measured as $\pm 0.328$, giving a standard error of $\div 2.13$ about the central line, assuming it were a true average and the distribution were normal.

## Debris Dispersion from a Reinforced Concrete Structure

The 1960 planned explosion at Pantex Ordnance Plant provided extensive debris dispersion data (Ref. 11). This explosion involved the detonation of $2,000 \mathrm{lb}$ of HE in the form of encased warheads placed inside the standard one-foot-thick reinforced concrete walls of a U-shaped bay. Two views of the structure are shown in Fig. 2. 38 and 2. 39. The structure Wis completely destroyed. Debris was dispersed over a large area, the m.tximum debris distance being about $1,500 \mathrm{ft}$ from the point of burst. Figure 2. 40 is a post-shot view of the close-in region. The area was canvassed and all fragments found were listed by terminal location. About 31,000 concrete fragments with a total weight of about $85,000 \mathrm{lb}$ were

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Figure 2. 37 Debris Dispersion and Standard Error for 1945 Badger Ordnance Works Explosion

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Figure 2.38
Pre-Shot View of Pantex Ordnance Plant Test Structure From Southeast

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Figure 2.39
Pre-Shot View of Pantex Ordnance Plant Test Structure From Northwest

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Figure 2. 40
Post-Shot View of Pantex Ordnance Plant Test Structure


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recorded along with their individual weights and terminal locations. Terminal locations were noted according to the serialized $50-\mathrm{ft}$ square zones shown in the missile map, Fig. 2.41. As the missile map indicates, debris dispersion was non-uniform with the greatest concentration being to the south -- the direction of the front wall from the point of burst. Debris data were not analyzed in the original explosion report.

Since the fragment list reported for the Pantex Ordnance Plant event constituted the most extensive compilation of data found on dispersion of fragments, these results were studied extensively to obtain data on fragment-size distributions and fragment dispersion.

Over-all dispersion of fragments from the Pantex event is tabulated in Appendix D and plotted in Fig. 2. 42 through 2.47. Two measures of dispersion are included here, specific area in sq ft per fragment, and areal dispersion in sq ft per lb of debris. The following plots have been prepared as alternate means of describing the dispersion function:

Fig. 2.42 $\begin{aligned} & \text { Log-Log Plot of Specific Area Vs Ground Range, } \\ & \text { with Second-Order Least-Squares Regression Line }\end{aligned}$
Fig. 2. 43 Semi-Log Plot of Specific Area Vs Ground Range, with First-Order Least-Squares Regression Line
Fig. 2. 44 Semi-Log Plot of Specific Area Vs Ground Range, with Second-Order Least-Squares Regression Line
Fig. 2. 45 Log-Log Plot of Areal Dispersion Vs Ground Range, with Second-Order Least-Squares Regression Line
Fig. 2.46 Semi-Log Plot of Areal Dispersion Vs Ground Range, with First-Order Least-Squares Regression Line
Fig. 2. 47 Semi-Log Plot of Areal Dispersion Vs Ground Range, with Second-Order Least-Squares Regression Line

In each case the dispersion measures, (Specific Area or Areal Dispersion), have been computed for $50-\mathrm{ft}$ and $100-\mathrm{ft}$ wide concentric circular bands about the point of detonation. Computed points were plotted at the midpoints of the circular bands. Since the original debris data from this event were collected for $50-\mathrm{ft}$ squares from coordinates near the point of burst, the approximation of the actual dispersion function is based on considering the debris within any $50-\mathrm{ft}$ square to be within the circular band containing its center. Regression lines are computed on the basis of all

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Figi. e 2. ti Debris Dispersion for Reinforced Concrete Structure
Logarithimic Plot with Second-Order Regression Line


Figure 2. 43 Debris Dispersion for Reinforced Concrete Structure Semilog Plot with First. Order Regression Line

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Figure 2. 44 Debris Dispersion for Reinforced Concrete Stiucture Semilog Plot with Second-Order Regressicn Line

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Figure 2.45 Areal Dispersion for Reinforced Concrete Structure Logarithmic Plot with Second-Order Regression Line

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points for the $50-\mathrm{ft}$ circumferential bands as listed in Table D-2.
Variation in average fragment weight with ground range is ploted in Fig. 2. 48. Again, these values were computed for concentric circular bands of $50-\mathrm{ft}$ and $100-\mathrm{ft}$ widths. Though the largest fragment recorded weighed about $4,000 \mathrm{lb}$, it is seen from this chart that the average weight of fragments at virtually all ground ranges was less than 12 lb . The bimodal nature of this curve is unexplained, and since this computation includes concrete fragments only, it does not stem from a two-phase nature of material. Relative distances from explosive sources to the various structural panels and the shielding effects of adjacent bays to the east of the bay containing the charge may have influenced this.

Fragment-size distribution for the $50-\mathrm{ft}$ and $100-\mathrm{ft}$ debris
zones is plotted in Fig. 2. 49 and 2.50. Figure 2.49 shows the total number of fragments above any indicated weight which were found at various ground ranges. This figure shows that relatively few fragments above three pounds were found at any ground range. An improved indication of variation in fragment-size distribution is obtained by plotting the cumulative percentage of fragments above indicated weights, as shown in Fig. 2.50. There is little separation between the cumulative-percentage curves over much of the range of fragment sizes in this figure. The pronounced "flattening-out" of these curves at the $3-1 \mathrm{~b}$ fragment weight suggests that an optimum design point for debris hazards may exist, providing of course that these curves are truly characteristic. At all ground ranges, less than 6 percent of total fragments was above three pounds in size. It is interesting to note that at the greatest ground ranges ( $1200-1500 \mathrm{ft}$ ) no fragments above one pound were found; but at the intermediate ground ranges (600-1200 ft) the percentage of fragments over three pounds in size was greater than at the shorter ground ranges ( $0-600 \mathrm{ft}$ ). The largest sized fragments (over 100 lb ) were found only at the shorter ground ranges. The tendency of making incomplete counts of the smallest fragments in the close-in region would actually result in the cumulative percentage curves of Fig. 2. 50 being overstated for the greater fragment sizes.

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Figure 2.46 Areal Dispersion for Reinforced Concrete Structure Semilog Plot with First-Order Regression Line




Figure 2.47 Areal Dispersion for Reinforced Concrete Structure Semilog Plot with Second-Order Regression Line

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Figure 2.48 Average Fragment Weight versus Ground Range for a Reinforced Concrete Structure

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Figure 2. 49 F::agment Quantities at Various Ground Ranges for a Reinforced Concrete Structure
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## Fragment dispersion for various weight classes is plotted in

Fig. 2. 51 and 2.52. Actual counts of fragments are presented in Fig. 2. 51. It is seen here that the heaviest fragments did not travel to the extreme ranges. Fragments of $1,000 \mathrm{lb}$ and heavier were found only within 200 ft , those of 100 lb or more were found only within 600 ft , and those of 10 lb or more were not found beyond 1200 ft . This is contrary to what would normally be expected on consideration of air drag as shown in Chapter Five. The comparison is not altogether appropriate, however, since in this actual structure, the fragments which are subject to forces of sufficient magnitude to cause large motion are also likely to be broken up to a greater extent. Proportionate distribution of fragments of various size classes, among the various ground ranges, is shown in Fig. 2. 52. Other than for the extremely heavy fragments (over 100 lb ) the proportionate distribution did not change much for the different fragment weights.

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Figure 2.51 Fragment Distances for Various Weight Classes for a Reinforced Concrete Structure

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Figure 2.52 Fragment Distance Distribution for Various Werght Classes for a Reinforced Concrete Structure

## CHAPTER THREE

FRAGMENTATION, EXPERIMENTAL OBSERVATIONS

To predict the hazards of debris from nuclear explosions, in addition to defining the maxımum range of debris and the distribution of debris within this lımit, it is well to define the nature (i.e., size or weight) of the expected missiles. Fragment weight combined with fragment velocity determines energy level of the fragment, and thus the loading upon equipment of personnel struck by the fragment. No definitive experimental investigation of structural fragmentation, relating fragment-size distribution to structural strength and loading parameters has yet been made. Recourse was made to collecting and summarizing past experimental investigations in fragmentation and to describing such general behavior and characteristic patterns as have been developed. Five such investigations are described in this chapter; research of the British Coal Utilization Research Association on coal breakage from random forces, research of the Safety in Mines Research Establishment of Great Britain on explosively detonated stone blocks, Stanford Research Institute model tests of containment structures fractured by internal explosions, an extensive study of fragments from a planned explosion of an ordnance structure at Pantex Ordnance Plant and findings of Project 4.5 of Operation JANGLE.

The fragmentation of materials produced by mine charges or ore crushers has been the subject of considerable study. Many attempts have been made to analyze the fractions produced by these processes. It has been shown experimentally that the higher the loading on the source material, the smaller the fragments produced, and that a wide range of fragment sizes are produced at any loading. An "Ideal Law of Breakage" which shows excellent fit to experimental data has been developed from coal-crushing investigations.

The extensive collection of data on the sizes of more than 30,000 concrete fragments from the Pantex Ordnance Plant event afforded the opportunity for a detailed analysis of fragment-size distribution. An interesting result of this particular study is that only about 3 percent of

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all the fragments produced weighed more than three pounds and that these fragments accounted for nearly 75 percent of total weight of all fragments recovered. Thus, as structures or equipment are built with increased resistance to fragments, the probability of being hit by a fragment sufficient to cause damage declines substantially.

The only structural fragmentation study conducted on full-scale nuclear tests was limited work performed on reinforced-concrete wall panels erected over the crater zone in Project 4.5 of Operation JANGLE (Ref. 5). Size distributions of the larger fractions were plotted as part of this project and it was noted that the JANGLE data did not preclude the possibility that the fragment-size distribution of concrete source material caused by the underground nuclear shot followed the same pattern as coal in a mine or ore in a crusher.

### 3.1 Fragmentation of Coal

A number of experimental investigations of coal fragmentation characteristics have been made. The mining industry, interested in minimizing dust formation and in producing fragments of an appropriate size for handling, has conducted numerous investigations of fragmentation from blasting and crushing operations. Assuming a brittle material with a random distribution of internal weaknesses, it has been deduced that when a single lump is fragmented by forces sufficiently violent to make breakage equally likely at any point, the broken product should have a fragment-size distribution obeying the exponential law (Ref. 13):

$$
M(x)=1-e^{-x / x}
$$

where

$$
\begin{aligned}
\mathrm{M}(\mathrm{x})= & \text { percentage of material smaller in size than } \mathrm{x} \\
\mathbf{x / \overline { x }}= & \text { dimensionless measure of fragment size, such as } \\
& \text { ratio of a characteristic length to a mean length. }
\end{aligned}
$$

This has been called the Ideal Law of Breakage. Experiments with large lumps of coal broken under conditions approximating random fracture tend to conform to the exponential law, at least for the smaller fragment sizes (under about $1 / 2$-in. equivalent spherical diameter). Fitting data to the

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equations

$$
M(x)=c\left(1-e^{-x / \bar{x}}\right)
$$

and

$$
M(x)=1-e^{-x / \bar{x}}
$$

researchers of the British Coal Utilization Board plotted experimental fragmentation data on broken coal of various types as shown in Fig. 3.1 (Ref.13).

## 3. 2 Fragmentation of Explosively Detonated Stone Blocks

In research conducted by personnel of the Safety in Mines Research Establishment, stemmed charges of $1-0 z$ and $2-0 z$ of coalmining explosive were fired in stone blocks 18 -in. in diameter by $30-1 \mathrm{n}$. long contained in a steel chamber (Ref. 14). Weights of the full-size. distribution fractions, including all dust, were measured.

Weights of the various fractions have been recomputed in terms of cumulative fragment quantities in Appendix E and are plotted in Fig. 3. 2. The cumulative curves all follow the same general trend, which is nearly linear on logarithmic coordinates. Slopes of the lines vary with the quantity of charge, and seemingly with the diameter of the shothole, both of which would be determinants of the impulse impinging on the stone blocks. Average slopes of the cumulative size-distribution curves for $2-o z$ charges were all significantly greater than those for $1-0 z$ charges and showed from only one-fifth to one-tenth the proportion of larger fractions. Thus, the larger quantities and smaller sizes of fragments were produced by detonations of the larger charges providing greater impulses -- certainly not an unexpected result. This tends to lend support to theories, that, in impulsive fragmentation, new surface area produced through fragmentation may be quantitatively related to explosive impulse or energy.

## 3. 3 Fragmentation of Concrete Shielding of Reactor Models

The Stanford Research Institute conducted experimentation on model concrete biological shielding poured directly in contact with model

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Figure 3.2 Cumulative Fragment Size Distribution for Exploded Dry Sandstone Blocks

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pressure vessels (Ref. 15). The models, shown in Fig. 3. 3, consisted of a reactor vessel, the concrete biological shield, and a stepped plug filling the access opening at the top of the vessel. Vessels were filled with varying amounts of water ( 75 and 100 percent) and subjected to the detonation of scaled energy sources centered radially and axially in the pressure vessel. The following model tests were selected for detailed plotting in this study, since individual fragment weights were available for these events. Original source data from these tests are included in Appendix $F$.

Table 3.1
MODEL TESTS SELECTED FOR PLOTTING

| Model <br> Test <br> No. | Event Simulated |  | Energy Source Material | Period, <br> (msec) | Water in Vessels, (\%) | Number of Fragments |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Pounds of TNT | Megawattseconds |  |  |  | Above 0.1 lb | $\begin{aligned} & \text { Total } \\ & \text { Recorded } \end{aligned}$ |
| 15 | 150 | 280 | Pyracore | 1 | 75 | 14 | 24 |
| 17 | 160 | 300 | Pyracore | 1 | 100 | 7 | 9 |
| 14 | 160 | 300 | Pyracore | 1 | 75 | 2 | 2 |
| 16 | 210 | 400 | MDF | 1 | 100 | 41 | 373 |
| 19 | 210 | 400 | MDF | 1 | 75 | 55 | 287 |
| 5 | 510 | 960 | Pyracore | 1 | 100 | 23 | 63 |
| 3 | 730 | 1,360 | Pyracore | 1 | 75 | 45 | 96 |

The cumulative total number of fragments above the indicated size is plotted in Fig. 3. 4 and 3.5. Figure 3.4 includes all recorded fragments, while Fig. 3. 5 includes only those fragments weighing more than 0.1 lb . Three explosion factors in the test setup which influence the loading are the percent of water in the vessel, the quantity of explosive (as indicated by the magnitudes of the simulated event), and the equivalence of the explosive source material. Tests, numbered 16 and 19 , in which MDF was used as the explosive source material produced the most

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Figure 3. 3 Model Reactor with Concrete Shielding

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Figure 3.4 Cumulative Fragment Distributions for $1 / 24$-Scaie Shielded $\underline{\text { Reactor Vesse.s (includes all recorded fragments) }}$


Figure 3.5 Cumulative Fragment Distributions for 1/24-Scale Shielded Reactor Vessels (includes all fragments over 0.1 lb)

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complete fragmentation. This certainly appears to be a function of the impulsive loading (or energy input) since the amount of gaseous reaction produce in the energy source is five times as high for MDF as for Pyracore ( $100 \%$ vs. 20\%).

Other than for Test 15, it was observed that fragmentation became more complete with increasing magnitude of the simulated event, higher water level in the vessel, and increased quantity of gaseous combustion products from the charge. These are all factors which would increase the loading on the concrete shielding. In Test No. 15, 24 fragments were produced and 85 percent of the total concrete weight was in the largest fragment.

A highly favorable result of this series of tests is the consistency in the shape of the cumulative curves of Fig. 3.4 and 3.5. Crossing-over of curves exists in these plots, but the general shape of the curves is consistent and the relative position of the successive curves appears to bear some relation to the impulsive loading.

Similar results were observed in a series of $1 / 12$-scale shielded reactor model tests (Ref. 16). Only the magnitude of the event simulated was varied in this series of tests, (data are tabulated in Appendix $G$ and plotted in Fig. 3.6). As Fig. 3. 6 shows, separation between the cumulative distribution curves for different events is decidedly pronounced, and the general shape of the curves is quite consistent. This series of tests appears to support contentions that fragmentation becomes more complete with increased loading and implies that the relationship between fragmentation and loading may be quantified, at least for simple ideal structures.

Stanford Research Institute selected three criteria to assess the fragmentation and fragment dispersion in the reactor shielding model tests. The quantity of new surface area created in the fragmentation process was regarded as an indicator of the energy absorbed in breaking the concrete shielding. It can be shown that this newly created surface area (termed fragmentation in the SRI reports) can be approximated from the individual fragment werghts:

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$$
\begin{aligned}
\text { Fragmentation } & \propto \text { New Surface Area } \\
& \propto \Sigma \mathrm{W}_{\mathrm{f}}^{2 / 3}
\end{aligned}
$$

where $\mathrm{W}_{\mathrm{f}}$ is the individual fragment weight.
Similarly, the throw, defined as the summation of the products of weight times distance thrown for individual fragments, was taken as an indicator of the total amount of input energy diverted into transport of fragments;

$$
\text { Throw } \& \Sigma W_{f} D_{f}
$$

where $D_{f}$ is distance thrown for the individual fragment.
A measure of structural integrity representative of the degree to which concrete shielding remained intact was defined as the ratio of the sum of the squares of individual fragment weights to the square of the total shielding weight;

$$
I \propto \frac{W_{f}^{2}}{(\Sigma W)^{2}}
$$

The inverse of this integrity ratio was considered to be a measure of the total damage to the shielding structure.

The measures of fragmentation, throw, and integrity were computed for the various reactor shielding models by SRI. Plots of these values are included in Fig. 3. 7 through 3.10. For the various series of models, the selected debris parameters (fragmentation, throw, and the inverse of the integrity ratio) show varying degrees of consistency in the shape of the plotted relationships. Some of the more pronounced deviations between plotted points and the curves as drawn are perhaps explanable in terms of the considerable scatter customarily experienced in explosion testing.

Figure 3.7 shows results of $1 / 12$ - and $1 / 24$-scale model tests for the debris criteria. In this figure the fragmentation and throw for the 1/l2-scale models are reduced by the reciprocal of appropriate scale factors:


Event Simulated, (Mw/sec)
Note: $20 \%$ gaseous products

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Note: All Models $100 \%$ Full of Water



Event Simulated, (Mw/sec)
Group $V$ - Reinforcing bars included in models.
Group VI - Insulated models.
Figure 3.10 Fragmentation Measures for Selected 1/24-Scale Shielded Reactor Model Tests

Fragmentation: $\quad A_{1} \sim \lambda^{2} A_{2}$, by 4
Throw: $\quad W_{f} D_{f} \sim \lambda^{3} \quad \lambda^{1 / 2} W_{f}^{\prime} D_{f}^{\prime}$, by $8 \sqrt{2}$
where $\lambda$ is the ratio of scale factors.
Comparison of debris parameters for model reactor shielding with different quantities of contained water are shown in Fig. 3. 8. Some, but not complete, separation of data points is apparent in the series of mudel tests in Fig. 3.9 and 3.10. With further tests, more conclusive definition of these functions might evolve.

No relationships can be established between the results of these reactor shielding model tests, fragmented by the impulsive loading of an internal explosion, and the fragmention behavior of other structures under nuclear blast loading. Under the longer duration nuclear loading, it is likely that peak overpressure may exert a greater significance than impulse on the fragmentation process. It is significantly demonstrated, however, that fragmentation patterns can be measured and related to loading, and that, for a uniform model structure consistent relationships are obtained. This suggests the feasibility of experimenting with suitable ideal structural elements, on a full-scale or model basis, to develop basic input data for debris hazard estimating purposes.

### 3.4 Fragmentation of a Reinforced Concrete Ordnance Structure

Fragment counts from the Pantex Ordnance Plant event described previously provided data on the individual weights of more than 31, 000 recovered concrete fragments with a total weight of more than $85,000 \mathrm{lb}$ (Ref. 11). Individual fragment weights recorded ranged from $1 / 16 \mathrm{lb}$ to $4,000 \mathrm{lb}$.

The fragment-size distribution of the recorded fragments is shown in Fig. 3. 11. Fragment classes in this figure are in a geometric progression, each class being twice as large as the preceding class. From this chart it appears that the greatest portion of fragments were in the weight range of 0.250 to 0.499 lb , with equivalent diameters of

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Figure 3.11 Fraginent-Size Distribution For a Reinforced Concrete Ordnance Structure

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1.8 to 2.2 in. The expected preponderance of relatively small fragments, under 4 lb or $4.4-\mathrm{in}$. diameter, is apparent. The relatively few fragments in the $0.060-0.249 \mathrm{lb}$ range is questionable and could result from the tendency to not record all the smallest items at close-in ranges (see Fig. 2.51).

Though the smaller fragments are produced in the greatest quantities, the amount of material involved is relatively small. The distribution of total weight of fragments, pictured in Fig. 3. 12, shows more than 65 percent of the total material fragmented into pieces weighing between 32 and 255 lb , that is, with equivalent diameters from 8.9 to 17.8 in .

Cumulative distributions of fragment sizes are approximated in Fig. 3.13. The median fragment weight is found to be about 0.22 lb , and more significantly -- more than 95 percent of all fragments weigh about 3 lb or less, i.e., have equivalent diameters of 4 in . or less. Less than 1 percent of all fragments weight 90 lb or had equivalent diameters greater than about l. 1 ft . In general, fragments above one pound in size are considered potentially lethal to personnel. Only about 20 percent of all recorded fragments from this test were above this size.

Cumulative weight distribution for fragments is presented in Fig. 3.14. Whereas the greatest number of fragments produced in this event were about 3 lb or less in size. Fig. 3.14 shows that somewhat over 50 percent of the total weight of fragments was in boulders weighing greater than 70 lb or having equivalent diameter of about 11.5 in . or more. Over 70 percent of total fragment weight was in pieces above 3 lb in size. Thus, although substantial quantities of concrete fragmented into large pieces, the number of these was small. The probability of being bit by small preces is much higher than that of being hit by large boulders.
3. 5 Fragmentation of Concrete Walls

The only fragmentation study of structural elements under nuclear loading was conducted as part of Project 4.5 of Operation JANGLE. For this event a series of six reinforced concrete wall panels was

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Equivalent Diameter, (in.)


Figure 3.14 Cumulative Distribution of Total Fragment Weight for a Reinforced Concrete Ordnance Structure

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erected over the crater zone of the underground shot as shown in Fig. 3. 15. This event involved a $1.2-\mathrm{KT}$ device emplaced at a depth of burst of 17 ft ; the apparent crater diameter was about 260 ft . The wall panels were thus placed fairly close-in relative to the total crater -- i. e., at ground ranges from 14 percent to 40 percent of apparent crater radius.


Figure 3.15 Placement of Wall Targets in JANGLE U Event
The fragment-size distribution of wall fragments reported from this test is presented in Fig. 3.16. These curves are plotted through the data points at a slope of 0.5 (the slope predicted on the basis of the Ideal Breakage Law). The JANGLE report pointed out that only the larger fragments were collected and that these would normally be expected to deviate from the ideal fragment-size distribution when a limited amount of material is available. It was concluded that the JANGLE data did not preclude the possibility that the breakup of concrete source material caused by a nuclear detonation follows the same pattern as coal in a mine or ore in a crusher.

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Figure 3.16 Size Distribution of Fragments from Wall Targets in JANGLE U Event

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It is noted that the relative positions on the curves (and their associated data points) are not arrayed in the same sequence as the initial ground ranges of the walls. The lowest curve, representing the largest fragments, corresponds to the wall closest to surface zero, where vertical components of the load may be expected to be higher. However the two topmost curves, representing the smallest fragment sizes, correspond to the middle two walls rather than the farthest-out.

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## CHAPTER FOUR

FRAGMENTATION, ANALYTICAL CONSIDERATIONS

### 4.1 State of Knowledge of Fracture Mechanics

Analytical consideration of debris resulting from the fracture of concrete or masonry structures poses two fundamental questions: 1) What is the fracture strength, i.e., the ultimate lad which the structure can sustain? and 2) What is the number and size of the resulting fragments? The first question is by far the easier to answer and it is discussed briefly from the point of view of existing fracture theory.

The theoretical strength of materials is of the order of 100 to 1000 times the observed strength. Griffith (Ref. 17) proposed that this difference can be rationalized in terms of pre-existing flaws contained in the solid. This model of a solid containing an array of flaws is the basis of the "fracture mechanics" approach to the problem. Treatment of this problem is concerned with the growth of pre-existing flaws contained in the solid. This model of a solid containing an array of flaws is the basis of the "fracture mechanics" approach to the problem. Treatment of this problem is concerned with the growth of pre-existing flaws and conditions of instability which can change the crack propagation from a slow process to that characteristic of a fast-running crack.

Griffith formulated the condition of stability, under load, of a body containing a certain type flaw upon two theorems. The theorem of minimum energy states that the equilibrium state of an elastic body, deformed by specific surface forces, is such that the potential energy of the whole system is a minimum. He obtained his new criterion of rupture by adding to the theorem of minimum energy the statement that the equilibrium position, if equilibrium is possible, must be one in which rupture of the solid has occurred, if the system can pass from the unbroken to the broken condition by a process involving a continuous decrease in potential energy. Thus, in the propagation of a crack, stored potential energy is released, but the potential energy of the system is increased by the creation of new surfaces (surface energy). Griffith's condition for

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continuing propagation of a crack is that the resultant potential energy of the system is decreasing. Likewise, the equilibrium crack size is one in which the decrease in potential energy just equals the increase in surface energy. The theory predicts strength reasonably well for bodies such as glass and ceramic which behave in a brittle fashion. Using the Griffith theory, we can reason that there will be a distribution of strengths in a given specimen in the sense that a different amount of force will be needed to fracture a specimen at one point than at another. If one assumes that the flaws are distributed at random with a certain density per unit volume, then the statistical formulation of the strength problem becomes apparent. The strength of a specimen is determined by the weakest point in the specimen or by the smallest value to be found in a sample of size $n$ where $n$ is the number of flaws. Clearly, $n$ increases as the volume increases and, therefore, the problem of finding out how the strength depends on the volume of the specimen is equivalent statistically to studying the distribution of the smallest value as a function of $n$, the sample size. This statistical problem is an important one on which much theoretical work has been done. For breaking strength the major contributions have been made by Weibull (Ref. 18).

Essentially, Weibull assumes that the probability of failure of a unit volume of material can be represented by a distribution function of the form

$$
\begin{align*}
& F(\sigma)_{.}=1-\exp \left[-\int_{V}\left(\frac{\sigma-\sigma_{y}}{\sigma_{0}}\right)^{m} d V\right]_{\sigma \geq \sigma_{y}}^{\sigma<\sigma_{u}}  \tag{4.1}\\
&=0 \\
& \text { where } \\
& F(\sigma)=\text { failure probability for stress } \sigma \\
& \sigma_{u}, \sigma_{0}, m=\text { constants of the material }
\end{align*}
$$

Once the unit probability of failure $F(\sigma)$ is known, it is a reasonably straightforward procedure to find the probability of failure of any structure under any known system of stresses. Note, however, that the failure mode is not unique since the location of the weakest link is a statistical quantity. Since the likelihood of failure is greatest in regions of high stress, it is certainly possible to anticipate the origin of failure. This predicability is relied upon when brittle test specimens are designed to break in the gage length. But even here it is quite common to get fractures outside of this region.

The applicability of the weakest link theories, and especially the Weibull theory, to the prediction of fracture strength in masonry and ceramics is a problem which is currently undergoind the most intensive investigation. Since the theories, at best, only treat the static load problem, their use in the exceedingly complex debris formation problem is simply too much to expect.

## 4. 2 Mathematical Model

Since a sophisticated treatment of fragmentation involves application of fracture mechanics to an extent that is well beyond the current state-of-the-art, our objective here is to formulate a relatively simple mathematical model which can be used to predict debris formation and which can be adapted to include advances in fracture mechanics.

The response of structures to nuclear blast loading is calculated by considering the response of an equivalent single-degree-of-freedom system (equivalent to neglecting all but the fundamental mode of vibration of the structure). The stiffness of this equivalent system is assumed to be elastic - perfectly plastic so that the structure's response into the plastic range can be considered.

Since this model has proved adequate for analyzing the response of structures up to failure, we propose a simple extension of the model to include fracture. The model is shown on Fig. 4. 1. A structure is represented as an equivalent mass $M_{e}$ and resistance $R(x)$. The

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Figure 4.l Fragmentation Model
resistance is taken to be elastic-perfectly plastic up to a displacement ( $\mathrm{X}_{\mathrm{f}}$ ) at which the structure fractures. Such a model can be used to predict the velocity and time at which the equivalent mass reaches the fracture displacement.

Limitations to this model are obvious. First, the mass distribution of fragments from the structure will not be determined by this model. Secondly, it is somewhat presumptuous to assume that all fragments are formed at a single displacement. These are questions however, which must be answered by advances (experimental and analytical) in fracture mechanics. Here, we say

$$
M_{f}=a M_{e}
$$

where,

$$
\begin{aligned}
M_{f}= & \text { mass distribution of fragments } \\
a= & \text { normalized mass distribution of fragments to be } \\
& \text { determined. }
\end{aligned}
$$

Similarly, there would be a statistical distribution about $\mathrm{x}_{\mathrm{f}}$ which would result in a corresponding distribution in the time of fracture and velocity spectra of the fragments. Thus, refinements of the basic simple model would be statistical in nature. An experimental program could be undertaken to establish the nature of these distributions for particular types of structures (e.g., brick walls, concrete structures, steel frames). Hence, it is feasible to obtain answers to these questions (at least insofar as required to modify our model) without waiting for a complete understanding of the fracture mechanics.

The equation of motion for the model of Fig. 4. 1(a) is,

$$
\begin{equation*}
M_{e} \frac{d^{2} x}{d t^{2}}+R x=F(t) \tag{4.2}
\end{equation*}
$$

For convenience, the following nondimensional parameters are defined:

$$
\begin{aligned}
& \tau=t / t_{0} \\
& \zeta=x / x_{e} \\
& \beta^{2}=\frac{R_{0} t_{0}^{2}}{M_{e} x_{0}}=\left(\omega t_{0}\right)^{2}
\end{aligned}
$$

$\omega=$ circular frequency of elastic system

$$
\begin{aligned}
\delta^{2} & =\frac{F_{o}}{R_{o}} \\
\rho & =\frac{x_{f}}{x_{e}}
\end{aligned}
$$

The equation of motion then becomes,

$$
\zeta^{\prime \prime}+\beta^{2} \zeta=\beta^{2} \delta^{2}(1-\tau)
$$

for

$$
\begin{equation*}
\zeta \leq 1 \tag{4.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\zeta^{\prime \prime}+\beta^{2}=\beta^{2} \delta^{2}(1-\tau) \tag{4.4}
\end{equation*}
$$

for

$$
1<\zeta \leq \rho
$$

where the double prime denotes differentiation with respect to $\tau$.
Equation 4.4 has a solution of the form (for zero initial conditions),

$$
\begin{equation*}
\zeta=\delta^{2}\left[-\cos \beta \tau+\frac{\sin \beta \tau}{\beta}+1-\tau\right] \tag{4.5}
\end{equation*}
$$

for
$\zeta \leq 1$,
and

$$
\begin{aligned}
\zeta= & \frac{\beta^{2}\left(\delta^{2}-1\right)}{2}\left(\tau-\tau_{e}\right)^{2}-\frac{\beta^{2} \delta^{2}}{6}\left[\tau^{2}-3 \tau\left(\tau_{e}\right)^{2}+2 \tau_{e}^{3}\right] \\
& +\zeta_{e}^{\prime}\left(\tau-\tau_{e}\right)+1
\end{aligned}
$$

for

$$
1<\zeta \leq \rho
$$

where

$$
\begin{aligned}
\tau_{e} & =\text { time at which } \zeta=1 ; \text { determined from Eq. (4.5). } \\
\zeta_{e}^{\prime} & =\text { nondimensional velocity at } \tau_{e} \text {; determined from Eq. (4.5). }
\end{aligned}
$$

Solutions for Eq. (4.5) and (4.6) were obtained for final velocity and the time of fracture. These solutions are presented graphically in Fig. 4. 2 through 4.4. Thus, the velocity of fragments and time of fracture can be found from this analysis. These would then be inputs to the analysis of debris transport which would provide a complete displacement history of the fragment.

## 4. 3 Example

As an example of the application of this analysis consider the wooden siding from a railroad car. The failure of such siding was observed during the UPSHOT-KNOTHOLE series of weapon effects tests (Ref. 19). The velocity at failure of siding located in the 10 -psi overpressure region was observed to be 60 fps .

## The ultimate elastic resistance of a simply supported rectangular

 beam is,$$
R_{0}=2 \frac{\mathrm{bd}^{2}}{L} \sigma_{y}=\frac{2 b(1.25)^{2} \times 1400}{10 \times 12}=37 b
$$



Figure 4.2 Motion of Mass At Failure $(\delta=5)$

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Figure 4.3 Motion of Mass At Failure $(\delta=10)$


Figure 4.4 Motion of Mass At Failure $(\delta=100)$

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where

| $b$ | $=$ width of member, | $L=$ span, and |  |
| ---: | :--- | ---: | :--- |
| $d$ | $=$ depth of member | $R_{o}=$ total uniform load on beam. |  |
| $\sigma_{y}$ | $=$ yield strength of material, |  |  |

The deflection of such a member is,

$$
x_{e}=\frac{R_{o} L^{3}}{6.4 \mathrm{Ebd}^{3}}=\frac{37 \times 10^{3} \times 1728}{64 \times 10^{6} \times(1.25)^{3}}=1.52 \mathrm{in} .
$$

The circular frequency of a simply supported rectangular beam
is

$$
\begin{aligned}
\omega & =\frac{\pi^{2}}{L^{2}} \sqrt{\frac{E}{12} \frac{d^{2} g}{\gamma}} \\
& =\frac{\pi^{2}}{100} \sqrt{\frac{10^{6}(1.25)(1.25) 32.2}{12 \times 40}}=31.8 \mathrm{rad} / \mathrm{sec}
\end{aligned}
$$

where

$$
\begin{aligned}
& \mathrm{E}=\text { Young's Modulus } \\
& \mathrm{g}=\text { gravitational constant } \\
& \gamma=\text { density of material }
\end{aligned}
$$

The 10 -psi overpressure region corresponds (Ref. 20) to a 2-psi dynamic pressure and a positive phase duration of 0.7 sec . The drag loading acting on the siding is then

$$
F_{o}=b L(2) c d=b \times 10 \times 12 \times 2 \times 1=240 b
$$

Therefore

$$
\delta^{2}=\frac{F_{o}}{R_{o}}=\frac{240 \mathrm{~b}}{37 \mathrm{~b}}=6.5
$$

and,

$$
\beta=w t_{o}=31.8 \times 0.7=22.3
$$

Then if we interpolate between Fig. 4. 2 and 4.3 and assume a fracture displacement of $\rho=20$,

$$
\begin{aligned}
\zeta_{f}^{\prime} & =300 \\
\tau_{f} & =0.125
\end{aligned}
$$

or, in dimensional form,

$$
\begin{aligned}
\dot{x}=\frac{x_{e}}{t_{o}} \zeta_{f}^{\prime} & =\frac{1.52 \times 300}{12 \times 0.7}=54 \mathrm{fps} \\
t_{f} & =0.125 \times 0.7=0.09 \mathrm{sec}
\end{aligned}
$$

Obviously, this fracture velocity is dependent on the value assigned to $\rho$. Therefore, the reasonably accurate prediction of this velocity does not establish the validity of our model. It does show, however, that reasonable values of $\rho$ lead to an acceptable prediction.

The subsequent motion of the debris fragments is studied in Chapter Five. The significance of debris velocity at fragmentation $x_{f}$ and time of fragmentation $t_{f}$ are also considered in subsequent chapters.

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## CHAPTER FIVE

## DEBRIS TRANSPORT BY BLAST WINDS

The motion of a particle acted on by the nuclear blast wind is considered with the fracture velocity and time of fracture forming the initial conditions of the transport problem. The model used here assumes that the force acting on the particle is proportional to the square of the relative velocity between the particle and air. The Last parameters are assumed to be constant over the range of travel of the debris, and it is further assumed that the effective lengthening of positive phase duration (i.e., the so-called time correction) which is due to the debris motion in the direction of shock propagation can be handled by a simple adjustment of positive phase duration. These assumptions reduce the problem to the solution of a one-parameter nondimensional differential equation. Thus, much can be learned about the general behavior of flying debris. Without these two assumptions, it would be necessary to treat each weapon yield and placement as a separate problem and no general observations could be made.

The general equations are derived and numerical solutions for many problems of interest are obtained. The application of these results to specific debris problems is outlined and general observations regarding the behavior of flying debris are detailed.

### 5.1 General Treatment of Problem

The motion of a piece of debris of arbitrary shape is first considered and then the problem is reduced to the point where the vertical and horizontal motions are uncoupled.

Consider the motion of the body shown in Fig. 5. 1. By equilibrium,

$$
\begin{align*}
m \ddot{x} & =f_{x} \\
m \ddot{y} & =f_{y}  \tag{5.1}\\
j \ddot{\phi} & =f_{y}\left(x_{1}-x_{0}\right)+f_{x}\left(y_{0}-y_{l}\right)
\end{align*}
$$

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## where

$m=$ mass of body
$j=$ product of inertia of body.
$\phi=$ coordinate describing orientation of body with respect to $x-y$ coordinate system


Figure 5.1 General Debris Particle

The forces ( $f_{x}, f_{y}$ ) are given by

$$
\begin{align*}
& f_{x}=k \frac{a C_{d(x)}}{2} \rho(u-\dot{x})^{2}  \tag{5.2}\\
& f_{y}=k^{\prime} \frac{a C_{d(y)}}{2} \rho \dot{y}^{2}+m g
\end{align*}
$$

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where,

$$
\begin{aligned}
a_{x, y} & =\text { projected area in } x \text { and } y \text { direction } \\
C_{d(x, y)} & =\text { drag coefficient in } x \text { and } y \text { direction } \\
\rho & =\text { mass density of air } \\
u & =\text { particle wind velocity } \\
g & =\text { gravitational constant } \\
k & =\left\{\begin{array}{l}
+1 ; u \geq \dot{x} \\
-1 ; u<\dot{x}
\end{array}\right. \\
\mathbf{k}^{\prime} & =\left\{\begin{array}{l}
+1 ; \dot{y}<0 \\
-1 ; \dot{y}>0 .
\end{array}\right.
\end{aligned}
$$

Equations (5.1) are coupled, in general, by the dependence of $\mathrm{aC}_{\mathrm{d}}$ on $\phi$. Consideration of this dependency on $\phi$ makes exact descriptions of debris particles necessary. We are, however, concerned with the gross behavior of groups of debris particles of many different sizes and shapes rather than a detailed treatment of specific shapes. Thus we restrict our attention to those particle shapes where $\mathrm{aC}_{\mathrm{d}}$ does not depend on $\phi$ and, hence, the equations of motion Eq. (5.1) are uncoupled. The trajectory dispersion which is due to $\phi$-variation of $\mathrm{aC}_{\mathrm{d}}$ can be investigated separately, thereby leading to a dispersion function which can be superposed on the results determined for the simpler problem considered here. Thus we consider only the first of Eq. (5.1). The following nondimensional parameters are defined:

$$
\begin{align*}
& T=\frac{t}{t_{0}} \\
& 0=U_{o} t_{0} \rho \frac{a C_{d}}{m} \\
& \zeta=\rho \frac{a C_{d}}{m} \times \tag{5.3}
\end{align*}
$$

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where
$U_{0}=$ characteristic air particle velocity such that

$$
\begin{aligned}
U & =U_{o} h(t) \\
t_{0} & =\begin{array}{l}
\text { characteristic time of air particle velocity-duration } \\
\text { curve }
\end{array}
\end{aligned}
$$

Substituting Eq. (5.3) into Eq. (5.1) we have

$$
\begin{equation*}
\zeta^{\prime \prime}=\frac{k}{2}\left[\theta h(\tau)-\zeta^{\prime}\right]^{2} \tag{5.4}
\end{equation*}
$$

where the prime denotes differentiation with respect to $\tau$.
Equation (5.4) is a Riccati-type nonlinear differential equation which can be linearized by the substitution,

$$
\begin{equation*}
\zeta^{\prime}=-k \frac{s^{\prime}}{s} . \tag{5.5}
\end{equation*}
$$

Equation (5.5) reduces Eq. (5.4) to

$$
\begin{equation*}
s^{\prime \prime}+k \theta h s^{\prime}+\theta^{2} \frac{h^{2} s}{2}=0 \tag{5.6}
\end{equation*}
$$

However, closed-form solutions can be obtained to Eq. (5.6) only for special classes of the forcing function $h(T)$ (e.g., $h(T)=$ constant leads to secondorder differential equations with constant coefficients; $h(\tau)=h_{o} / \tau$ leads to a Cauchy-Euler equation). The blast-induced wind is unfortunately not one of these forms. Therefore, we are forced to integrate Eq. (5.4) numerically. The case of $h=0$ is of interest and the solution for this case can be readily derived from Eq. (5.5) and (5.6) to be,

$$
\zeta=-k \log \left(\tau+C_{1}\right)+C_{2}
$$

and note that,

$$
\begin{equation*}
\zeta^{\prime}=\frac{-k}{T+C_{1}} \tag{5.7}
\end{equation*}
$$

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where $C_{1}, C_{2}$ are constants of integration to be determined from initial conditions.
Observe from Eq. (5.7) that the velocity only asymptotically goes to zero. This can be explained by recognizing that when the surrounding air is motionless the retarding force acting on the particle is proportional to the square of the particle velocity. Thus, solutions based on our model will predict infinite displacement at an infinite time. It will be necessary to place a maximum time limitation on the debris particles flight based on the time required for it to hit the ground.

Numerical solutions to Eq. (5.4) are obtained for a wide class of problems of interest. Specifically, results are obtained for a wide range of aerodynamic coefficients, various initial conditions, and a few possible negative phase representations of the blast induced winds.

We must consider first the form of the wind loading $h(\tau)$, however. The particle velocity can be related to the dynamic pressure $q$ by,

$$
\begin{equation*}
q=\frac{1}{2} \rho u^{2} . \tag{5.8}
\end{equation*}
$$

During the positive phase of the loading, the dynamic pressure is

$$
\begin{equation*}
q=q_{0} e^{-2 \tau}(1-\tau)^{2} \tag{5,9}
\end{equation*}
$$

When Eq. (5.9) is substituted into (5.8)
$u=\sqrt{\frac{2 q_{0}}{\rho}} e^{-\tau}(1-\tau)$.
Thus, in terms of Eq. (5.4),

$$
u=u_{0} h(\tau)
$$

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where, for $0<\tau<1$

$$
\begin{align*}
u_{0} & =\sqrt{\frac{2 q_{0}}{p}}  \tag{5.11}\\
h(T) & =e^{-T}(1-T)
\end{align*}
$$

Variation of $\sqrt{\frac{2 q_{o}}{\rho}}$ and $t_{0}$ with weapon yield and distance from ground zero is shown on Fig. 5.2 and 5. 3.

Relatively little data are available regarding the negative phase dynamic pressure (i.e., wind blowing toward ground zero). We have taken the negative phase dynamic pressure to be of the form,

$$
\begin{equation*}
\bar{q}=\bar{q}_{o} \sin \pi \frac{(\tau-1)}{k} \tag{5.12}
\end{equation*}
$$

so that wind ceases at $T=1+k$.
It can be argued that, to have ambient conditions at ground zero some time after detonation of the weapon, the area under the positive phase dynamic pressure curve must equal the area under the negative phase curve. This results in,

$$
\begin{equation*}
K=0.34 \frac{\mathrm{q}_{0}}{\overline{q_{0}}} \tag{5.13}
\end{equation*}
$$

The wave form descriptions apply to a fixed location with reference to ground zero. The debris particles, however, move with respect to ground zero. Note that $u_{0}$ and $t_{0}$ vary with distance from ground zero so that, to exactly represent the forcing function, the particle and pressure wave motion must be followed. As an example, consider a particle originally located at ( x ) and assume the pressure wave arrives at this location at time $t=0$. The peak particle velocity and positive phase duration are $u_{0}(x)$ and $t_{0}(x)$. At some later time $(\Delta t)$ the particle has moved to ( $x+\Delta x$ ); the pressure wave arrives at this location at some time less than $(\Delta t)$, say $\left(\Delta t^{\prime}\right)$. Thus, the particle velocity at this time and at the actual debris location is given by,

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$$
\begin{equation*}
u=u_{0} \exp \left[-\left(\Delta t^{\prime}-\Delta t\right) / t_{0}\right]\left[1-\left(\frac{\Delta t^{\prime}-\Delta t}{t_{0}}\right)\right] \tag{5.14}
\end{equation*}
$$

where

$$
u_{0}, t_{0} \text { are evaluated at }(x+\Delta x)
$$

Equations (5.14) cannot be used to represent the particle velocity unless attention is restricted to specific weapon yields. Thus, to retain generality in the study, this refinement is not considered here. Note, however, that the form of Eq. (5.14) is well suited for inclusion in a numerical integration routine such as will be used to integrate the equations of motion.

## 5. 2 Numerical Results

A computer program was written for the UNIVAC 1105 digital computer to numerically integrate Eq. (5.4). The Runge-Kutte-Gill numerical integration procedure was used. Solutions were obtained for zero initial conditions of the debris particle, and a complete range of realistic aerodynamic coefficients. The effect of negative phase wind on particle motion is also considered. It was the objective of this numerical analysis to obtain results in a form such that mass-velocity curves can be constructed for a given attach condition and location from a source of debris.

### 5.2.1 Zero Initial Conditions Excluding a Negative Phase

We consider first the piece of debris that is free to move, acted upon only by the positive phase wind loading. Solutions are obtained for values of $\theta$ in the range $0.1<\theta<10$. These are presented in Fig. 5.4 and 5.5 in the form of curves of nondimensional velocity versus nondimensional distance from original debris source. Time is included as a parameter on these curves.

As stated, attention has been limited to the horizontal motion of a particle. These solutions are invalid if the particle hits the ground. Therefore, the time parameter on the curves of Fig. 5.4 and 5.5 can be used to establish the range of validity of the results based on initial height

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$h$ of the debris and any vertical component of velocity $v_{o}$ imparted to the particle. Based on elementary kinematics the particle will hit the ground at

$$
\begin{equation*}
t=\frac{v_{0}}{g}\left[1+\sqrt{\frac{2 h g}{v_{o}^{2}}+1}\right] \tag{5.15}
\end{equation*}
$$

An interesting observation regarding the size and shape of particle that will travel the furthest can be drawn from these data. To this end, the parameter $\theta$ is plotted as a function of distance for various times on Fig. 5.6. Note that the results of this analysis predict that the size and shape (i.e., aerodynamic coefficient) of the piece of debris that travels furthest depends on the length of time required for the debris to hit the ground surface. The longer this time, the larger the piece of debris that travels furthest. This could explain the apparent randomness in highexplosive debris data.

## 5. 2. 2 Zero Initial Conditions Including a Negative Phase

As has been mentioned, little data which quantitatively describes the negative phase wind are available. A sine wave is assumed for the wave shape of the negative phase wind, and the area under the sine wave is made to equal the area under the positive phase wind. In this study, the peak negative phase wind velocity was taken to be 0.02 times the peak positive phase wind velocity.

Velocity-distance curves for this negative phase wind are given in Fig. 5. 7 and 5. 8. The sharp departure of these results from those for zero negative phase is of conside rable interest. Given sufficient time, all particles move toward ground zero. Of course, these required times of flight are sufficiently long that this reversal of velocity can probably not be realized for most practical cases. There have been instances, however, at the Nevada test site where pieces of structural debris were found closer to ground zero than the place they started. Of course, structural debris requires some time to be ripped from the structure so that a portion of the positive-phase wind impulse is not effective on the piece of debris. Thus the velocity reversal occurs sooner and can be realized with practical times of flight.


Figure 5.5 Velocity-Distance Plot for $0.2<\theta \leq 1$

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Figure 5.6 Comparison of Distance Traveled by Different Size Particles


Figure 5.7 Velocity-Distance Plot Including Negative Phase For $1<\theta<10$


Figure 5.8 Velocity-Distance Plots Including Negative Phase For $0.2<\theta<1$

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## 5. 2. 3 Effect of Initial Conditions

The results described in the previous two sections must be modified to include the initial condition of failure time and velocity at failure. One could accomplish this modification by constructing velocity-distance curves (similar to Fig. 5.4 and 5.5) for all combinations of initial velocity and failure times of practical interest. Such an approach is impractical because of the number of combinations that would need to be considered. Rather, an approximate method of modifying the existing data (Fig. 5.4 and 5.5 ) to include the initial conditions must be found.

Two approximations are made in this regard, and the errors induced in practical problems are evaluated. First, the failure time conditions is approximated by applying to the debris particle a negative impulse equal to the area under the dynamic pressure time curve from time $t=0$ to time $t=t_{f}$. Mathematically then, the apparent initial impulse is given in nondimensional form as,

$$
\zeta_{a}=\zeta^{\prime}(0)-\frac{\rho a^{2} c_{d}{ }^{2} t_{o}{ }^{2} q_{o}}{m^{2}} \int_{0}^{T} e^{-2 \tau}(1-\tau)^{2} d \tau
$$

where

$$
\begin{aligned}
\zeta_{a}^{\prime} & =\text { apparent initial velocity } \\
\zeta 1(0) & =\text { actual initial velocity } .
\end{aligned}
$$

when this is carried out the apparent initial velocity becomes,

$$
\zeta_{a}^{\prime}=\zeta(0)-\frac{\theta^{2}}{4}\left[e^{-2 \tau_{f}}\left(\tau_{f}^{2}-\tau_{f}+\frac{1}{2}\right)-\frac{1}{2}\right] .
$$

The failure time $\tau_{f}$ is always small ( $\tau_{f}<l$ ) so that if $e^{-2 \tau_{f}}$ is expanded in a power series and $\tau_{f}^{2}$ is neglected in comparison to $\tau_{f}$. Then,

$$
\begin{equation*}
\zeta_{a}^{\prime}=\zeta^{\prime}(0)-\frac{\theta^{2}}{2} \tau_{f} \tag{5,16}
\end{equation*}
$$

results.

The second approximation made was to assume that the particle motion resulting from the apparent initial velocity acting alone can be superposed on the particle motion which is due to the blast wind, to give the resulting motion of the particle which is due to the combined effect of initial velocity and blast wind. By virtue of the general dependence of air friction on the square of the relative velocity, the two results are superposed by the square root of the sum of the squares. The motion which is due to the initial velocity can be determined from Eq. (5.7) subject to the initial conditions

$$
\begin{aligned}
\zeta(0) & =0 \\
\zeta^{\prime}(0) & =\zeta_{a}^{\prime}
\end{aligned}
$$

The solution of Eq. (5.7) is then,

$$
\begin{align*}
& \zeta_{1}^{\prime}=\frac{-k \zeta_{a}^{\prime}}{\zeta_{a}^{\prime} \tau-k} \\
& \zeta_{1}=-k \log \left(\frac{\tau \zeta_{a}^{\prime}-k}{-k}\right) . \tag{5.17}
\end{align*}
$$

For the case of interest,

$$
k=\left\{\begin{array}{l}
-1 ; \zeta_{a}^{\prime} \geq 0 \\
+1 ; \zeta_{a}^{\prime}<0
\end{array}\right.
$$

In summary, the procedure for applying these results to a real problem is as follows:
(1) Compute $\zeta_{a}^{\prime}$ from Eq. (5. 16)
(2) Compute $\zeta_{1}^{\prime}(\tau)$ and $\zeta_{1}(\tau)$ from Eq. (5.17) for the range of $\tau^{\prime}$ s of interest.
(3) Use Fig. 5.4 or 5.5 to determine $\zeta_{2}^{\prime}(\tau)$ and $\zeta_{2}(\tau)$ for zero initial conditions for the same values of $T$ as above.

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Figure 5.9 Comparison of Exact and Approximate Method of Handling Initial Conditions

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(4) At each given ( $\tau$ ) compute the final results,

$$
\begin{aligned}
& \zeta=\sqrt{-k\left(\zeta_{1}\right)^{2}+\left(\zeta_{2}\right)^{2}} \\
& \zeta^{\prime}=\sqrt{\left.-k\left(\zeta_{1}^{\prime}\right)^{2}+\zeta_{2}^{\prime}\right)^{2}}
\end{aligned}
$$

The validity of these results was tested by comparing the approximate results obtained by the above procedure to "exact" solutions obtained by numerically integrating Eq. (5.4) subject to the initial conditions. These comparisons are presented on Fig. 5. 9. It can be noted that the results are reasonably accurate in view of the over-all accuracy requirements of the debris prediction problem. For very high initial velocities (e.g., velocity displacement curves convex downward with a high initial peak) the approximate analysis is not very good. However, this problem is probably not very important when considering the nuclear environment.

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## CHAPTER SIX

## VULNERABILITY OF FIELD TROOPS TO TREE DEBRIS

A supplementary study involving vulnerability of field troops to casualties from tree debris caused by a nuclear explosion in the proximity of a forest, instituted upon recognition of needs by the Office of the Chief of Engineers for estimating the hazards to engineer and field troops, was undertaken. Making certain simplifying assumptions (i. e., zero-strength tree limbs, plane blast wave loading, unobstructed trajectories, and that a hit upon personnel by a tree limb is a casualty), tree limbs are followed in their trajectories from the time of shock loading to their impact with the ground. The safe distance may fall either inside or outside the forest and both cases are treated.

Results of this analysis show that for the lower yields (l KT, for example) a uniform horizontal translation of all branches is obtained since trajectories become vertical with an attendent small vertical drop, and the appearance of the area in front of the forest up to the "safe distance" would be similar to that of a forest floor after all of the branchwood were allowed to drop vertically. Basically similar behavior is observed for the higher yields where trajectories terminate before they become vertical ( 20 MT , for example), with the exception that the highest branches of the first few rows of trees pile up in a lower density than those following the closer-in trajectories.

## 6. 1 Previous Studies

The results of previous studies on tree vulnerability which have influenced our assumptions concerning debris (Ref. 21, 22, and 23) are as follows:
(1) Results of OPERATION UPSHOT-KNOTHOLE indicated that stands of 145 Ponderosa pines of heights 50 to 75 ft offered no attenuation of peak overpressure or dynamic pressure.
(2) That low burst heights were found to cause more damage to trees than large burst heights when the peak dynamic pressure was the same.
(3) OPERATION CASTLE indicated no pressure attenuation from the trees in natural tree stands. Damage predictions for two weapon yields compared favorably with the observed damage.
(4) Damage to broadleaf stands is principally limb breakage and defoliation with occasional breakage of the main stem or uprooting.
(5) The deflection and breakage of trees in the stand on UPSHOT-KNOTHOLE Shot 9 were approximately twice the values predicted on the basis of calculations of the first maximum deflection. By including the probability of breakage during the second maximum deflection under the negative phase, the predictions were brought into agreement with experimental values.
(6) Trees are drag structures. The best parameter with which drag can be correlated is the dry weight of the crown.
(7) In OPERATION SNAPPER, it was observed that when stem breakage occurred the stems broke at the tree base.
(8) A considerable amount of data describing the mechanical and aerodynamic properties of tree stems and crowns is available. Almost none exist for the isolated branches.

### 6.2 Problem Approach

Consider first the case shown in Fig. 6. 1 where the troops are dispersed outside of the forest.


Figure 6.1 Relative Positions of Troops, Forest, and Ground Zero

Our objective is to determine a conservative or upper bound "safe distance" rather than tackle the enormous task of finding the "smallest safe distance". To accomplish this goal, we made the following assumptions:
(1) Ground burst is assumed for all attack conditions since it produces the most severe tree damage.
(2) Trees are assumed not to attenuate static or dynamic pressures.
(3) The trajectories of the flying branches are considered to be unobstructed by the other branches and trees.
(4) If any debris strikes an individual, he is assumed to be incapacitated.
(5) The total "flight time" of a branch is considered to be equal to its free-fall time in a vacuum. This assumption turns out to be unimportant for yields no greater than 1 MT .

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(6) Positive wind phase is taken equal to the positive shock wave phase. They are approximately equal.
(7) The branches travel the same horizontal distarice as an air particle during the positive wind phase.

Assumption 7 is the most far-reaching and important on the list. Physically it corresponds to a branch which has zero strength, is completely diffraction insensitive, and which has an infinite acceleration coefficient. The acceleration coefficient is the product of projected area and drag coefficient divided by mass. Clearly, for tree branches this is a relatively large factor. Assuming the branches to be only drag sensitive is fairly good; assuming them to have zero strength is conservative but poor. The items tend to counteract one another.

If only drag forces were operative, it is quite clear that a particle of air would be transported further than a solid object. The questions arise when the solid object has an initial velocity; for example, particles originating as crater throwout. Such particles are not of concern here; however, it is conceivable that sufficient impulse is delivered to a branch to both sever it from the tree stem and give it an initial velocity greater than the peak particle velocity. In such a case the branch or any other particle would experience a deceleration resulting from a drag force opposing its motion. The high acceleration coefficient of a branch would quickly bring its velocity into coincidence with the particle or wind velocity behind the shock front.

During the negative phase, the reversed winds decelerate any airborne objects. Particles haveing sufficiently high acceleration coefficients have their forward motion reversed. In the spirit of conservatism, the negative phase is neglected.

At low pressure levels ( 2.4 psi ), the stems of trees remain standing and offer considerable interference to the flight of branches. However, at pressures of interest (between 5 and 30 psi ) the stems are all broken and probably create no interference to flying debris. Because the shock wave moves at a higher velocity than the winds behind it, there is a strong possibility that debris from remote trees travels ahead of debris

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from trees closer in to ground zero.
The trajectories of the wind particles have been computed as described in Appendix $H$ and the results are presented in Fig. 6. 2. The following remarks are based on, or are concerned with, this figure.
(1) Pressures over 30 psi are not included in trajectories since these pressures represent a greater hazard than the debris.
(2) As the wind velocity drops to zero the traje tories become vertical. The curves for the $20-\mathrm{MT}$ devices were terminated at about 255 ft of vertical drop since trees of greater height are not of interest.
(3) The horizontal distance associated with the vertical tangent to the trajectories can be scaled approximately according to cube-root scaling; e.g., $D / D_{0}=\left(W / W_{0}\right)^{1 / 3}$ where $D 1 s$ the distance and $W$ the yield.
(4) In the cases where the trajectories become vertical with an attendant small vertical drop (e.g., the case of the l-KT device), a uniform horizontal translation of all the branches was obtained. This situation is illustrated in the sketch shown in Fig. 6. 3. The appearance of the area in front of forest labeled "safe distance from forest' would be similar to that of a forest floor after all of the tree branchwood were allowed to drop vertically. This situation is illustrated in Fig. 6. 4 and 6.5 for forests exposed to low pressure blasts.
(5) In cases where the trajectories terminate before they become vertical (e.g., the case of the 20-MT device), basically the same behavior as in the previous case was found, with the exception of the highest branches on the first few rows of trees. These branches pile up with a lower density than those following the closerin trajectories. This situation is depicted in Fig. 6.6 where the final horizontal locations of branch centers

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Figure 6.3 Trace of Branchwood for Low Yield Weapons
for various branch elevations in a typical forest are indicated. Each cross represents the branchwood in a 5 -ft vertical distance along the tree stem. The most cursory examination of typical trees indicates that the limbs in the top 15 ft would completely cover the area surrounding the tree. Referring to Fig. 6.6, we find that "complete kill" is experienced for distances up to 1800 ft . Going from 1800 ft to 1880 ft , the density of limbs diminishes and 100 percent kill is not expected in this region. The precise determination of the kill probability in this $80-\mathrm{ft}$ area is not warranted and in such case the safe distance from forest is set equal to the maximum trajectory (e.g., 1880 ft ).


Figure 6.4
Forest Stand After a Nuclear Explosion
(2. 4 psi overpressure)


Figure 6.5
Forest Stand After a Nuclear Explosion
(3.8 psi overpressure)

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Figure 6. 6 Horizontal Displacement of Tree Limbs

We now turn our attention to the very straightforward task of determining the safe distance from ground zero in a forest of infinite extent. In Table 6.1 we have reproduced a schedule of damage criteria for forests from the ART 6. 24 of Reference 20. The damage level described in class $D$ of this table represents the borderline condition for troop safety.

Table 6.1
DAMAGE CRITERIA FOR FORESTS

| Damage <br> Class | Nature of Damage | Equivalent Hurricane <br> Wind Velocity, <br> (mph) |
| :--- | :---: | :---: |
| A and B | Up to 90 percent of trees blown down; <br> remainder denuded of branches and <br> leaves. (Area impassable to vehicles <br> and very difficult on foot.) | $130-140$ |
| C | About 30 percent of trees blown down; <br> remainder have some branches and <br> leaves blown off. (Area passable to <br> vehicles only after extensive clearing.) | $90-100$ |
| DVery few trees blown down; some leaves <br> and branches blown off. (Area passable <br> to vehicles.) | $60-80$ |  |

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Table 6.2
SAFE DISTANCES TO PREVENT CASUALTIES FROM TREE DEBRIS

| $\begin{gathered} \text { Yield } \\ \text { of } \\ \text { Weapon } \end{gathered}$ | Height of Tree, | Safe Distance Fally lnvide of Forest | Safe Distance Fall <br> Safe Distance from Ground Zero for Various Overpreseures at the Front of the Forest, (yd) |  |  |  |  | Safe Distance from Forent for Various Overproserrea at the Front of the Forest,$(y d)$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | feel | 1.7 - 2.2 psi | 5 psi | 10 psi | 15 psi | 20 psi | 30 psi | 5 psi | 10 psi | 15 psi | 20 pbi | 30 psi |
| 0.05 KT | 20 |  | 193 | 130 | 109 | 98 | 87 | 7 | 8 | 10 | 12 | 16 |
|  | 40 |  | 200 | 136 | 113 | 100 | 87 | 14 | 14 | 14 | 14 | 1\% |
|  | 60 |  | 206 | 142 | 119 | 106 | 91 | 20 | 20 | 20 | 20 | 20 |
|  | 80 |  | 213 | 149 | 126 | 113 | 98 | 27 | 27 | 27 | 27 | 27 |
|  | 100 | 311-376 yards | 220 | 156 | 133 | 120 | 105 | 34 | 34 | 34 | 34 | 34 |
|  | 120 |  | 226 | 162 | 139 | 126 | 111 | 40 | 40 | 40 | 40 | 40 |
|  | 160 |  | 239 | 175 | 152 | 139 | 124 | 53 | 53 | 53 | 53 | 53 |
|  | 200 |  | 253 | 189 | 166 | 153 | 138 | 67 | 67 | 67 | 67 | 67 |
| 0.1 KT | 20 |  | 241 | 164 | 137 | 123 | 108 | 7 | 10 | 12 | 15 | 19 |
|  | 40 |  | 248 | 168 | 139 | 123 | 108 | 14 | 14 | 14 | 15 | 19 |
|  | 60 |  | 254 | 174 | 145 | 128 | 109 | 20 | 20 | 20 | 20 | 20 |
|  | 80 |  | 261 | 181 | 152 | 135 | 116 | 27 | 27 | 27 | 27 | 27 |
|  | 100 | 392-473 yards | 268 | 188 | 159 | 142 | 123 | 34 | 34 | 34 | 34 | 34 |
|  | 120 |  | 274 | 194 | 165 | 148 | 129 | 40 | 40 | 40 | 40 | 40 |
|  | 160 |  | 287 | 207 | 178 | 161 | 142 | 53 | 53 | 53 | 53 | 53 |
|  | 200 |  | 301 | 221 | 192 | 175 | 156 | 67 | 67 | 67 | 67 | 67 |
| 0.5 KT | 20 |  | 410 | 279 | 235 | 210 | 186 | 10 | 16 | 21 | 26 | 33 |
|  | 40 |  | 414 | 279 | 235 | 210 | 186 | 14 | 16 | 21 | 26 | 33 |
|  | 60 |  | 420 | 283 | 235 | 210 | 186 | 20 | 20 | 21 | 26 | 33 |
|  | 80 |  | 427 | 290 | 241 | 311 | 186 | 27 | 27 | 27 | 27 | 33 |
|  | 100 | 670-809 yards | 434 | 297 | 248 | 218 | 187 | 3.4 | 34 | 34 | 34 | 34 |
|  | 120 |  | 440 | 303 | 254 | 224 | 193 | 40 | 40 | 40 | 40 | 40 |
|  | 160 |  | 453 | 316 | 267 | 237 | 206 | 53 | 53 | 53 | 53 | 53 |
|  | 200 |  | 467 | 330 | 281 | 251 | 220 | 67 | 67 | 67 | 67 | 67 |
| 1 KT | 20 |  | 517 | 351 | 295 | 26.4 | 233 | 12 | 20 | 26 | 32 | 41 |
|  | 40 |  | 519 | 351 | 295 | 264 | 233 | 14 | 20 | 26 | 32 | 41 |
|  | 60 |  | 525 | 351 | 295 | 264 | 233 | 20 | 20 | 26 | 32 | 41 |
|  | 80 |  | 532 | 358 | 296 | 264 | 233 | 27 | 27 | 27 | 32 | 41 |
|  | 100 | 845-1021 yards | 539 | 365 | 303 | 266 | 233 | 34 | 34 | 34 | 3.4 | 41 |
|  | 120 |  | 545 | 371 | 309 | 272 | 233 | 40 | 40 | 40 | 40 | 41 |
|  | 160 |  | 559 | 385 | 323 | 286 | 246 | 54 | 54 | 54 | 54 | 54 |
|  | 200 |  | 572 | 398 | 336 | 299 | 259 | 67 | 67 | 67 | 67 | 67 |
| 1 MT |  |  |  |  |  |  |  |  |  |  |  |  |
|  | 40 |  | 5,137 | 3,460 | 2,895 | 2,574 | 2, 243 | 86 | 151 | 204 | 251 | 326 |
|  | 60 |  | 5,147 | 3,475 | 2,916 | 2,599 | 2, 273 | 96 | 166 | 225 | 276 | 356 |
|  | 80 |  | 5,153 | 3,485 | 2,927 | 2.612 | 2. 290 | 102 | 176 | 236 | 289 | 373 |
|  | 100 | 8,448-10,208 yards | 5,157 | 3,490 | 2.934 | 2,620 | 2. 299 | 106 | 181 | 243 | 297 | 382 |
|  | 120 |  | 5,160 | 3,493 | 2,937 | 2,623 | 2, 302 | 109 | 184 | 246 | 300 | 385 |
|  | 160 |  | 5,163 | 3,495 | 2,938 | 2,623 | 2, 302 | 112 | 186 | 247 | 300 | 385 |
|  | 200 |  | 5,163 | 3,495 | 2,938 | 2,623 | 2, 302 | 11: | 186 | 247 | 300 | 385 |
| $\therefore \mathrm{MT}$ | $\therefore 0$ |  | 13,792 | 9, 128 | 7, 505 | 6,553 | 5,528 | 81 | 146 | 200 | 247 | 525 |
|  | 40 |  | 13,821 | 9, 179 | 7, 575 | 6,639 | 5,640 | 110 | 197 | 270 | 333 | 437 |
|  | 60 |  | 13,842 | 9,215 | 7,624 | 6,699 | 5, 718 | 131 | 233 | 319 | 393 | 515 |
|  | 80 |  | 13,858 | 9,243 | 7,662 | 6,746 | 5,779 | 147 | 26) | 357 | 440 | 576 |
|  | 100 | 22,932-27.709 yarde | 13,871 | 9,267 | 7. 693 | 6,786 | 5,829 | 160 | 285 | 388 | 480 | 626 |
|  | 120 |  | 13,882 | 9,287 | 7,720 | 6, 818 | 5,872 | 171 | 305 | 415 | 512 | 667 |
|  | 160 100 |  | 13,901 13,916 | 9,304 9.345 | 7,764 | 6,872 6,914 | 5.940 5.994 | 190 205 | 322 363 | 459 49. | 56 506 608 | 737 |
|  | 200 |  | 13,916 | 9,345 | 7,799 | 6.914 | 5.994 | 205 | 363 | 49.4 | 608 | 791 |

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The wind velocities $60-80 \mathrm{mph}$ are associated with overpressures of 1.7 2. 2 psi. The distance from ground zero for a surface burst at which these pressures are realized can be scaled from Fig. 3. 94a in Reference 20.

### 6.3 Results

Based on the approaches described in this chapter, we have computed the safe distances to prevent casualties from tree debris for various conditions and presented them in Table 6.2. Whenever the safe distance from the forest, which was computed from the maximum trajectory range, fell below the associated tree height, the tree height was used as the safe distance from the forest. At the pressure levels considered the trees will be blown over; hence, the stems and not the branchwood represents the more severe hazard under these conditions.

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## CHAPTER SEVEN <br> VULNERABILITY OF FIELD TROOPS TO THROWOUT DEBRIS FROM CRATERING AND STREAM-BED CHARGES

This study was undertaken to define a "safe line" for positioning engineer or field troops in the proximity of very-low yield nuclear cratering and stream-bed charges, based on debris criteria. Since no reliable analytical method of predicting crater throw out debris was avalable, the problem involved locating and utilizing experimental data on debris distribution under variations in the major controlling parameters - weapon yield, depth of burst, and soil characteristics. A number of sources were found to include data concerning crater throw out debris, some including variations in parameters (Ref. 2, 24, 25, 26, 27, 28, 29). Of these, two modes of measurement are used: the U.S. Geological Survey (Ref. 24) expresses debris distribution in terms of fragment sizes; the Suffield Experiment Station (Ref. 27) presents debris data in terms of a real density; and the Boeing Airplane Company (Ref, 25 and 26) presents data in both ways.

The procedure described here for estimating the throwout environment about cratering charges is based primarily on the U.S. Geological Survey reports of cratering tests in basalt in Area 18 at the Nevada test site (Ref. 24). Contours were plotted by USGS defining ground ranges for several average particle sizes. Average ranges for various fragment sizes were derived from these plotted contours and related to the yield and depth of burst. Nomographic methods are included whereby the average debris distance can be obtained for a wide range of weapon yields, depths of burst, and fragment size in basalt. A multiplication factor is introduced to convert the average ground range to maximum ground range, based on observations of the ray-like patterns noted in the USGS report on basalt craters. Data from the earlier Panama Canal series of tests were used to develop correction factors in converting the estimates for basalt to estimates for other soil media (Ref. 30). Likewise, debris measurements for cratering tests in marine muck, conducted as part of the Panama Canal series of experiments, were used to provide estimates of debris distance for streambed charges.

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## 7. 1 Method of Solution

In general, the debris beyond the crater lip has been observed to conform to the following two expressions;

$$
\begin{equation*}
\rho=\frac{C_{1}}{x^{n_{1}}} \tag{7.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathrm{D}=\frac{\mathrm{C}_{2}}{\mathrm{x}^{\mathrm{n}_{2}}} \tag{7.2}
\end{equation*}
$$

where

$$
\begin{aligned}
\rho & =\text { areal density in weight of debris per unit area } \\
D & =\text { fragment size } \\
\mathbf{x} & =\text { distance from ground zero } \\
C_{1}, C_{2} & =\text { constants } \\
n_{1}, n_{2} & =\text { exponents }
\end{aligned}
$$

Data which verify Eq. 7.1 are available (Ref. 2, 25, 26, and 27). Data which verify Eq. 7. 2 are also available (Ref. 2 and 26).

Equations 7.1 and 7.2 comprise the basis for two different sets of relationships which can be used for determining safe line distances.

Certain question regarding the use and validity of Eq. 7.2 must remain. First, there is the greater tendency for fragments which travel greater distances to break into smaller fragments when they hit the ground. Therefore, even if all fragments are of equal size at takeoff, a decrease in particle size for increased distance would still be witnessed. Secondly, when air resistance is considered, it can be reasoned that, in general, large particles will travel further than small particles. Of course, there are other considerations such as the origin of different-sized particles relative to the point of burst, which lend reasonability to the observed distribution. It is not known which consideration is most important. In future tests, it would be desirable to consider the occurrence of impact breakage, since the criteria for safety involves the fragment characteristics before impact rather than after impact.

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Let us now consider the criteria for the determination of a safe line ground range. The Armed Services Explosives Safety Board has used a $1-1 b$ fragment as being capable of producing a fatal injury. If the safety criteria were established considering any injury less than fatal as safe, then if material density is known it is possible to use Eq. 7.2 to find the safe line ground range.

To apply the safety criteria to Eq. (7.1), consider the area covered by an average soldier laying flat on the ground. For a given areal density, the most severe situation is that in which all of the material landing in the area covered by the soldier were lumped into a single fragment. The distance at which the areal density is such that would require this fragment to be equal to the critical weight (e.g., l lb) is the safe line ground range.

The one-pound fragment considered here is only an example. The actual size of fragment selected is a function of the probability of injury which is considered desirable. A much smaller fragment size might be selected if velocities were high.

When a criterion for determining the equipment damage capabilities of fragments is developed, it will also be possible to apply Eq. (7.1) and (7.2) to equipment.

Before any problems can be solved, we must find the functional relationship between the constants $C$ and $n$, and the independent blast parameters. These parameters consist of yield, depth of burial, and soil characteristics. A fourth parameter, geological conditions, is also important (Ref. 24). It is very difficult, however, to evaluate geological factors, especially in military situations, and these will therefore not be considered as a parameter. Let us first attempt to develop relationships between the three blast parameters and the $C$ and $n$ constants for the distribution in fragment size.

## 7. 2 Fragment Size Distribution Method

Ross B. Johnson (Ref. 24) has used maps to present data concerning throw out from a series of high explosion craters in basalt. Contour lines of equal average particle size for $1.0-\mathrm{ft}$ and $0.5-\mathrm{ft}$ diameter particles

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are presented. (See Fig. 7.1, for example.) This information is used to obtain a first approximation of the relation between $C, n$, and the blast parameters. It must be emphasized at this point that the relations developed here are only first approximations.

Data from eight charges of 1000 lb TNT and three charges of $40,000 \mathrm{lb}$ TNT exploded at various depths are tabulated in Table 7. 11. The areas within the contour lines of each map were measured and the radii of circles of equivalent areas calculated. Each radius so calculated is termed the "area mean distance for particles of given size". Accepting the size distribution law,

$$
\begin{equation*}
x=\frac{C}{D^{n}} \tag{7.3}
\end{equation*}
$$

where Eq. (7.3) is simply a modification of Eq. (7.2), the area mean distance was plotted as a straight line on a log-log plot (Fig. 7. 2). The slopes of the curves $n$ and the coefficients $C$ were found and tabulated in Table 7.1. Table 7. 1 expresses $C$ in scaled form. The resulting data were then used in plotting Fig. 7. 3 and 7.4. Figure 7.3 is a plot of the Distribution Law Exponent, $n$, vs scaled depth of burial, $d / W^{1 / 3.5}$, and Fig. 7.4 is a plot of scaled distribution law coefficient, $c / W^{1 / 3.5}$, vs scaled depth of burst, d/w ${ }^{1 / 3.5}$.

A scaling factor of 3.5 was found to produce the best fit. Its use agrees with data presented by R.B. Vaile of Stanford Research Institute (Ref. 31). Figure 7.5 shows the Vaile curves of crater radius vs yield for different materials. The curve for sandstone yields a scaling exponent of 3.6. Since Johnson's data are for basalt it may be expected that they also have ascaling coefficient in the vicinity of 3.6 , hence, 3 . 5 .

Note that both curves of Fig. 7.3 and 7.4 have the general appearance of an inverted parabola, similar to the curves of scaled crater radius and scaled crater volume versus scaled depth of burial found in many references on cratering. Notice on Fig. 7. 4 that not all points fit the curve closely. At first glance one might expect an elipse. However, the results of the $1000-1 \mathrm{~b}$ TNT shots are encouraging. Each data point for the $1000-\mathrm{lb}$ shots represents two shots whereas each point for the $40,000-1 \mathrm{~b}$ shots


Figure 7.1
High Explosion Crater 13 (40,000 pounds), Buckboard Mesa, Nevada Test Site, Nye County, Nevada


Outer limit of area one-half covered by fragments of all sizes
$\qquad$
Outer limit of fragments of one-foot or greater maximum dimension
0.5
$1.0=$
Contour lines of equal average fragment size
Only contours for average size of 0.5 and 1.0 foot fragments are shown

5230 ~
Post-explosion topographic contours (by Am. Arial Surveys) Interval 10 feet. Datum is sea level
Table 7.1
DEBRIS DISTANCES FOR CRATERING TESTS IN BASALT

| Crater Number | Charge <br> Weight, W, (lb TNT) | Depth of Charge (ft) | Area Mean Distance of 0.5-ft Diam <br> Fragments, (ft) | Area Mean Distance of $1.0-\mathrm{ft}$ Diam <br> Fragments, (ft) | Average of Mean for 0.5-ft Diam Fragments, (ft) | Average of Mean for 1.0-ft Diam Fragments, (ft) | Scaled Depth of Charge d/3. $5 \sqrt{W}$ | n | Scaled <br> Coeffi- <br> cient <br> c/3.5 $\sqrt{\mathrm{W}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2 | 1,000 | 20 | 32.1 | 18.2 ) | 34:8 | 17:2 | 2:0 | 1:01 | 2:39 |
| 7 | 1,000 | 20 | 37.4 | $17.5\}$ |  |  |  |  |  |
| 3 | 1,000 | 15 | 33.9 | 16.8 \} | 46.4 | 20.6 | 1. 5 | 1.17 | 2. 88 |
| 8 | 1,000 | 15 | 58.8 | 24.4 |  |  |  |  |  |
| 4 | 1,000 | 10 | 45.6 | 18.2 \} | 41.6 | 16.0 | 1.0 | 1.43 | 2. 22 |
| 9 | 1,000 | 10 | 37.5 | 13.8 ) |  |  |  |  |  |
| 5 | 1,000 | 5 | 33.5 | $11.9\}$ | 32.6 | 10.7 | 0.5 | 1.06 | 1.49 |
| 10 | 1,000 | 5 | 31.7 | 9.5 |  |  |  |  |  |
| 11 | 40,000 | 25.6 | 33.9 | 15.1 | 33.9 | 15.1 | . 75 | 1.17 | 0.73 |
| 12 | 40,000 | 42.8 | 34.8 | 15.1 | 34.8 | 15.1 | 1.25 | 1.21 | 0.73 |
| 13 | 40,000 | 59.8 | 57.0 | 28.8 | 57.0 | 28.8 | 1.75 | 0.99 | 1.39 |
| Note: | $\begin{aligned} x_{\operatorname{avg}} & = \\ W & = \\ d & = \end{aligned}$ | verage last yie depth of | mean distanc in lb TNT urial |  |  |  |  |  |  |



Figure 7.2 Area Mean Distance versus Fragment Size

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: 'unaodxt met wo!nq!insia
Figure 7.3 Distribution Law Exponent versus Scaled Depth of Burial

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Figure 7.5 Crater Radius versus Yield, Surface Shots, Various Soils

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represent only one shot. Data for the $1000-1 \mathrm{~b}$ shots show in gene ral, what would be expected, i.e., a curve similar to an inverted parabola. Especially encouraging is the fact that the vertex of the parabola lies very close to the optimum scaled depth of burial for basalt. The data for the 40,000-1b shots show a great deal of scatter.

Figure 7.6 is used to determine the equations for the curves of Fig. 7. 3 and 7.4. Offsets from axes placed through the vertical were plotted in Fig. 7.6 to obtain coefficients and exponents. The resulting equations are:

$$
\begin{align*}
& \mathrm{n}=1.28-0.25\left(\lambda_{\mathrm{c}}-1.62\right)^{1.8}  \tag{7.4}\\
& \mathrm{C}=\mathrm{w}^{3.5}\left[2.34-0.52\left(\lambda_{\mathrm{c}}-2.25\right)^{1.73}\right] \tag{7.5}
\end{align*}
$$

where
$\mathrm{n}=$ distribution law exponent
$C=$ distribution law coefficient
$W=$ charge weight in lb of TNT
$\lambda_{c}=\mathrm{d} / \mathrm{W}^{1 / 3.5}$
$\mathrm{d}=$ depth of turial

Equations (7.3), (7.4), and (7.5) are the basis of the charts of Fig. 7.7 and 7. 8. These charts can be used to predict safe line ground range. Let us reiterate that this represents only a first approximation. Instructions for the use of the charts follows. Two examples are presented on the charts.
7.2.1 Nomographic Calculation of Safe Distances

Based on Size Distribution

Instructions for Use of Charts
Find n, and C as follows:
(1) Select yield of device in Fig. 7.7 on horizontal axis.
(2) Draw a vertical straight line upward to the diagonal line representing the desired depth of burst, $d$, (ft).

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Figure 76 Distribution Law Exponent and Constant versus Scaled Depth of Burial

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Figure 7.7 Nomographic Chart No. 1 for Computing Values of $C$ and $n$

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Fig're 3.8 Nemigrapher Chat No ¿for Compting Debris Distance
(3) Draw a horizontal line from the resulting intersection with the diagonal to both the $n$ and $C$ curves.
(4) Draw a line vertically upward from intersection on $n$ curve to the top scale and find the value of the distribution law exponent $n$.
(5) Draw a line vertically from the intersection with the $C$ curve to the line $A-A$.
(6) From the intersection with A-A, draw a horizontal $l_{\text {ne }}$ to intersection with the diagonal (parallel to A-A) line of constant device yield.
(7) Draw a vertical line from the resulting intersection upward to the $C$ scale and find the value of the distribution law coefficient, $C$. Use values of $n$ and $C$ found above as input for Fig. 7.8 and find maximum average distance of encounter with particle of diameter $D$.
(1) Select diameter of particle on the lower horizontal axis.
(2) Draw a vertical line upward to the diagonal representing the value of $n$ found previously.
(3) Draw a horizontal line from the resulting intersection to the diagonal line of constant value of the $C$ found previously.
(4) Draw a line vertically to the upper horizontal scale to find the maximum radial distance from ground zero at which the particle of diameter $D$ may be expected.

Note now that distances found by this chart are average ground ranges at which a given average particle size will be found. In reference to Fig. 7. 1 it will be recognized that a large degree of variability can be attributed to the geological characteristics of the soil, such as discontinuities, nonhomogeneity and anisotropy.

To estimate a safe line, it is necessary to take these factors into consideration. This can be done by considering the maximum variation about the average. Table 7. 2 tabulates the variations by showing the calculation of the ratio of maximum contour line distance and average contour

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Table 7.2
DEBRIS CONTOURS FOR CRATERING EXPLOSIONS IN BASALT


[^1]
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line distance. The resulting ratios were plotted on semilog paper versus yield in lb of TNT in Fig. 7.9.

Note the general decrease in $x_{\text {max }} / x_{\text {avg }}$ ratio for increasing yield. This is exactly what would be expected for the following reason. The soil will appear to be very nonhomogeneous, discontınuous, and anisotropic to a small blast. As the blast yield increases the appearance of the soil will become homogeneous, continuous, and isotropic. The radomeness in throw out will therefore become less as blast yield is increased and consequently we can expect the asymmetry ratio-vs-yield curve will be asymptotic to the line
$\frac{x_{\text {max }}}{x_{\text {avg }}}=1$.
Data from DANNY BOY have been included in Fig. 7. 9. However, compatibility with the other data is questionable since these values of asymmetry are based, at least in part, on elevation contours instead of particle size contours. The term asymmetry is used since it describes the throw out pattern. Figure 7.10 shows the DANNY BOY contours.

Figure 7.9 can be used to find an asymmetry factor. This factor should be multiplied by the average contour line distance found from Fig. 7. 3. Since the results are still only a first approximation it is advisable to apply a safety factor of at least 1.5 to the results. Additional charts can be developed to take the asymmetry and safety factors into account automatically. The final result will be a safe line ground range for blast set off in basalt.

If it is desired to scale from basalt ot the other materials, Fig. T. 11 will provide factors which can be applied to the final result obtained, However, in some cases such as Cucaracha and Culebra, and marine muck, this scaling may be questionable since it is so difficult to consider fragments when speaking of these sorls. The scaling may possible be good only for the materials capable of fragmentation. The curves of Fig. 7. 11 were developed from the debris curves of Fig. 7.12 taken from the Panama Canal Studies (Ref. 29).

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Figure 7.9 Asymmetry versus Yield


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Figure 7.11 Debris Diameter Ratio versus Scaled Depth of Burial

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Scaled Depth of Charge, (ft/lbs ${ }^{1 / 3}$ )

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### 7.2.2 Explosive Equivalence

The data presented in Fig. 7. 7 and 7.8 are for TNT charges. If nuclear charges are to be used, it is necessary to find an equivalence factor which would provide a means of finding the TNT equivalent of the nuclear device. The state-of-the-art concerning equivalence is uncertain at present. One can only hypothesize on the basis of the limited data that exist.

The first problem concerning equivalence is one of semantics. Depending upon the terminology used, almost any factor ranging from 2 percent to 120 percent is acceptable. First, let us here define equivalence. For this study, a TNT and a nuclear device would be equivalent when each produces the same amount of throw out debris distributed at the same distance.

One approach to equivalence is to hypothesize on the basis of the energy available in a nuclear blast. Effects of Nuclear Weapons (Ref. 20) states that energy of a nuclear blast is divided as follows*:

| Blast and Shock | -50 percent |
| :--- | :--- |
| Thermal Radiation | -35 percent |
| Residual Nuclear Radiation | -10 percent |
| Initial Nuclear Radiation | -5 percent |

On this basis it can be expected that at the ground surface a nuclear blast will be at least 50 percent efficient in terms of debris-producing capabilities. If the blast is set off underground a portion of the 35 percent thermal radiation will be transformed into mechanical effects as a result of the vaporization of material in the immediate vicinity of the device. Therefore, for buried bursts efficiencies between 50 percent and 85 percent are expected. The above does not take into account variation due to partition of energy between the air and soil - this could, in fact, result in a substantially lower efficiency. For our purposes the more conservative estimate has been selected.

[^2]For bursts near the surface, venting will occur and therefore a portion of the thermal radiation will still escape. On this basis one would expect a curve of efficiency versus scaled depth of burst tc be similar to that of Fig. 7.13


Figure 7.13 Percent of Efficiency vs Scaled Depth of Burst

## -. 2. 3 Comparison with DANNY BOY Results

Maximum distances for fragments of various sizes have been computed by the methods of this chapter, using DANNY EOY explosion parameters as inputs. This was an $0.430-\mathrm{KT}$ device buried at 110 ft in basalt. Predicted debris distances are plotted in Fig. 7. 14 for various equivalence factors according to the methods described here.

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Figure 7.14 Estimated Maximum Fragment Distances for DANNY BOY Event

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## 7. 3 Throwout Debris from Stream-Bed Charges

Since energy partitioning for an explosion at the interface between two media is inversely proportional to the ratio of the densities of the media, the bottom burst of a stream-bed charge would be expected to impart more blast energy to the soil than a comparable surface burst at an air-soil interface. Thus, the stream-bed charge would dislodge more material for contribution to the debris, which we shall here regard only as the above-water ejecta. Initial velocities of the ejecta from stream-bed charges are unknown, but on the basis of energy partitioning we can concede, conservatively, that they may exceed initial velocities of ejecta from the comparable surface burst on land. Ejecta from the bottom burst may follow three paths, assuming the initial gas bubble does not break the surface.
(1) Upward (nearly vertical) within the rising gas bubble into the atmosphere at relatively high velocities.
(2) Upward (nearly vertical) through water, with considerable retardation.
(3) Through the water at various angles to the water surface, with considerable retardation.

Bottom charges in deep stream beds would not be expected to throw debris to greater distances than comparable surface bursts on land, for only the debris rising through the gas bubble would be expect.ed to reach the atmosphere at velocities comparable to those of the ejecta from surface bursts on land.

Although little consideration has been given crater throw out debris from underwater explosions or bottom bursts, either nuclear or high explosive, the following observation supports the above conclusions.
(1) The displaced bottom material did not produce airborne debris considered of consequence in the "Baker" test of OPERATION CROSSROADS. This event involved an explosion of a $20-\mathrm{KT}$ device at $180-\mathrm{ft}$ depth in the Bikini Lagoon, which is considred relatively shallow for that yield (Ref. 30 ).

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[^3]

$\begin{array}{lll}W=300 \mathrm{fbs} & D_{\mathrm{b}}=2 \mathrm{ft} & \lambda_{\mathrm{c}}=0.30 \mathrm{ft} / \mathrm{bs} \mathrm{s}^{1 / 3} \\ \mathrm{D}_{\mathrm{f}}=6 \mathrm{ft} & \mathrm{t}_{\mathrm{P}}=0.25 \mathrm{ecc} & \mathrm{h}_{\mathrm{p}}=120 \mathrm{ft}\end{array}$


(charge on bottom) $W=300 \mathrm{lb}$
$\mathrm{D}_{\mathrm{f}}=4 \mathrm{ft}$
$\mathrm{R}_{\mathrm{p}}=23 \mathrm{ft}$

Figure 7.15
Plumes From Low-Yield Underwater Explosions

$D_{b}=$ depth of burat below water surface
$D_{f}=$ depth of bottom
$h_{p}=$ reported plume height from photograph

( 1 hsin linear acale Bkinl Tent Baker
Note-Note-
W=yield
$\lambda_{c}=$ ocaled depth of charge
epetime of plume photograph
$\mathbf{R}_{\mathrm{p}}=$ plume radius


Figure 7.15 (Cont'd)
Plumes From Low-Yield Underwater Explosions

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Figure 7.16
Plume From Shallow Underwater Nuclear Explosions

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Among all the experimental work reviewed on cratering and debris, the Panama Canal tests of cratering charges in marine muck most closely resemble the configuration of the shallow stream-bed charge -- at least they assuredly provide a situation where the ejecta is not retarded in its trajectory by passing through water. Energy partitioning may not be exactly comparable to the stream-bed charge, for the immediate escape of the gas bubble would also cause more energy partitioning away from the bottom material.

In the absence of analytical methods for estimating debris behavior from bottom bursts, recourse is made to the available experimental data on crater tests in marine muck. Debris limits for stream-bed charges can be computed by estimating debris distances for basalt as prescribed earlier for cratering charges, and then applying a correction to account for the ratio of debris limits for basalt and bottom material. This correction factor,

## Debris diameter for bottom material Debris diameter for basalt

can be taken as the debris-diameter-ratio for marine muck in Fig. 7.11. In using Fig. 7.11 in this manner the scaled depth of burial of the charge should be taken relative to the stream bed. Thus for bottom bursts, scaled depth of burial will be zero.

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? $\underset{\sim}{a}=$






| Sunflower Ordnance Wks Gibbstown, N.J. | 1903 | Black Powder Dynamite | Mixer in Building Barricaded Building |
| :---: | :---: | :---: | :---: |
| Upton Towans, England | 1899 | NG | Barricaded Building |
| Guinan, Samar, P.I. | 1945 | Frag. Bombs | Barge |
| Upton Towans, England | 1899 | Gelatin | Barricaded Building |
| Landing, N.J. Lower Hope Point, England | $\begin{array}{r} 1909 \\ 1902 \end{array}$ | Dynamite NG | Barricaded Building Barricaded Building |
| Uplee's Marshes. Faversharn, England | 1903 | NG | Bar:onatj Buidity |
| mahburn, Mo. | 1906 | NG | Trar |
| Savanna Urdnance Depot | 1949 | Fuzes, Minejetc. | Eartratoveiedighon |
| fangnung. Kurea | 1947 | Projectiles | Truc |
| Lewisburg, Ala. | 1907 | Dymarite | Guilding neat Hint |
| Ashburn, Mo. | 1906 | Gclatin | Barr.ataded Yuildng |
| Penson, Ariz. | 1949 | mon. Nit. Dyn. | Barricadedictib, 8'al3 |
| Lgyptian Powder Co. |  | Black Powaer | Unbarricaded Shed |
| Schaghticoke, N.Y. | 1912 | Black Powder | Barricaded Building |
| Chattanooga, Tenn. | 1903 | Dynamite | Barricaded Building |
| Moosic, Pa. . |  | Black Powder | Unbarricaded Building |
| Baraboo, Wis. | 1943 | E. C. Blank Fire Powd | , Barricaded Building |
| Ashburn, Mo. | 1906 | Dynamite | Barricaded Building |
| Pembrey |  | TNT | Barricaded Building |
| Chilworth, England | 1879 | Black Powder | Barricaded Building |
| Reynold, Pa. | 1951 | NG | Barricaded Building |
| Carney'a Point |  | Shotgun Powder | Continuous Dryer |
| Umbogintwini, Natal | 1909 | NG | Barricaded Building |
| Cliffe, England | 1904 | NG | Barricaded Building |
| Hazelton, Pa. | 1905 | Dynamite | Unbarricaded Building |
| Liberty Powder Co. |  | Eynamite | Barricaded 3 Sides |
| Airfield, England | 1944 | G. P. Bombs | Aircraft Loading |
| Mt. Braddock, Pa | 1942 | Dynamite | Barricaded Building |
| Eldred, Pa | 1950 | 60\% $\mathrm{Amm} . \mathrm{Gel}$. Dym. | Barricaded Building |
| Waltham Abbey, England | 1894 |  | Barricaded Building |
| Irwine |  | TNT | Barricaded Building |
| Gibbstown, N. J. | 1913 | Gelatin | Barricaded Building |
| Boulder, Colo. | 1907 | Dynamite | Freight Car |
| Upton Towans, England | 1904 | NG | Unbarricaded Building |
| Kuba, Okinawa | 1945 | Dynamite | 20' x 20' Tomb |
| Winsted, Conn. | 1892 | Black Powder | Freight Car |
| Nanaimo, B.C. | 1911 | Dynamite | Barricaded Building |
| Carney's Point, N.J. Birmingham, Ala. | $\begin{aligned} & 1926 \\ & 1951 \end{aligned}$ | Cannon Powder Dyn. 30\% Amm. | Unbarricaded Building Barricaded Building |
| Birmingharn, Ala. |  | $\begin{aligned} & \text { Dyn. So } \\ & \text { ell. NG } \end{aligned}$ | Barricaded Building |
| Moosic, | 1944 | Black Powder | Building |
| Gibbstown, N. J. | 1929 | NG | 24' Frame Shed |
| Radford Arsenal | 1954 | Mid Cannon Powder | Barricaded Building |
| Hercules, Calif | 1953 | Mixed Dope | Unbarricaded Building |
| Pinar Del Rio, Cuba | 1910 | Dynamite | Unbarricaded Building |
| Martin, Pa | 1913 | Dynamite | Unbarricaded Building |
| Great Ashfield, England | 1943 | Bombs | Aircraft |
| Deenethorpe |  | 50/50Amatal Bombs | Aircraft |
| Hercules, Calif. | 1953 | Black Powder | Unbarricaded Building |
| Blackbeck, England | 1898 | Black Powder | Barricaded Building |
| Gibbetown, N.J. | 1908 | Nitro-Starch | Unbarricaded Building |
| Ardeer, Scotland | 1902 | NG | Barricaded Building |
| Hercules, Calif. | 1948 | NG | Barricaded Building |
| Barge Pier, Korea | 1945 | C-2 | Landing Craft |
| Ardeer, Scotland | 1902 | NG | Barricaded Building |
| lehpeming, Mich. | 1905 | NC | Unbarricaded Building |
| Ardeer, Scotland | 1914 | NG | Barricaded Building |
| Seneca, Ll . | 1951 | 40\% Amm. Gel. | Barricaded Building |
| Barksdale, Wis. | 1907 | NG | Barricaded Building |
| Louviera, Colo. | 1908 | NG | Barricaded Building |
| A | 1943 | TNT | Unbarricaded Building |
| Haskell, Okla. | 1912 | NG | Unbarricaded Building |
| Hoyle Factory | 1916 | Nitro-Cotton | Barricaded Building |
| Wolf Creek | 1944 | Ammonium Nitrate | Building |


| \# | $\stackrel{\text { N }}{\sim}$ |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Indiana Ordnance Wks |  | Smokeless Powder | Building | 81.59 | 0.58 |
| Beira, Portugal | 1880 | Black Powder | Mag. | 82.50 | 0.58 |
| Highland Station, Calif. | 1892 | Dynamite | Unbarricaded Building | 103.6 | 0.79 |
| Haskell, N.J. | 1917 | Smokeless Powder | Unbarricaded Building | 117.0 | 0.58 |
| Rio de Janeiro, Brazil | 1925 | Dynamic | Unbarricaded Building | 124.0 | 0.79 |
| Arco, Idaho | 1945 | 50/50 Amatol | Barricaded Igloo | 125.0 | 0.87 |
| Savanna Ordnance Depot | 1948 | Tetrytol | Earth Covered Igloo | 147.0 | 1.20 |
| Kobe, Japan | 1910 | Dynamite | Barge | 150.8 | 0.79 |
| Manila, P.I. | 1924 | Dynamite, etc. | Unbarricaded Mag. $50 \times 150$ | 173.1 | 0.79 |
| Indiana |  | Gun Powder Rifle Powder | Frame Bldg, and RR Cars | 177.6 | 0.58 |
| Charleston |  | Rifle Powder | Magazine |  | 0.58 |
| Tessenderloo, Belgium |  | Ammon. Nitrate | Unbarricaded Building | 193.0 | 0.42 |
| Black Tom Island, N. Y. Harbor | 1916 | TNT | Freight Cars | 200.0 | 1.00 |
| Mindi Magazine, Canal Zone | 1914 | Dynamite | Unbarricaded Mag. | 225.5 | 0.79 |
| Sonemachi, Japan | 1946 | HE | Unbarricaded Dump | 270.0 | 1.00 |
| Baltimore, Md. | 1913 | Dynamite | Steamer | 300.0 | 0.79 |
| Acisate (Varese), Italy | 1948 | Ammo. | Underground Bunkers | 350.0 | 1.00 |
| Guadalcanal |  | Torpex | Steamer | 400.0 | 1.25 |
| Bari, Italy |  |  | Steamer | 544.0 | 1.10 |
| Hastings, Neb. |  | Torpex, TNT and DB Powder | Barricaded Bldg. $500 \times 25$ | 550.0 | 1.17 |
| Pleasant Prairie, Wis. |  | Black Powder | Unbarricaded Mag. | 587.5 | 0.58 |
| Bombay, India | 1944 | HE . | Steamer | 400.0 | 1.00 |
| Lake Denmark, N.J. | 1926 | TNT | Unbarricaded Mag. | 800.0 | 1.00 |
| Bucharist, Rumania | 1924 | HE | Building | 1,000.0 | 1.00 |
| Mt. Hood, Pacific Theatre | 1944 | HE | Steamer | 1,000.0 | 1.00 |
| Brest, France | 1947 | Ammonium Nitrate | Steamer | 730.0 | 0.42 |
| Texas City, Texas | 1947 | Ammonium Nitrate | Steamer | 2,280.0 | 0.42 |
| Port Chicago, Calif. | 1944 | HE | Steamer | 2,136.0 | 1.00 |
| Halifax, Nova Scotia | 1917 | HE | Steamer | 2,600.0 | 1.00 |
| Burton-on-Trent (Fauld) |  | Misc. Bombs | Barricade | 2,670.0 | 1.00 |
| Oppau, Germany | 1921 | Ammonium Nitrate | Open Pile | 4,500.0 | 0.42 |

## MUM DEBRIS DISTANCE AND EXPLOSION PARAMETERS FOR SELECTED EXPLOSIONS

|  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Building | 81.59 | 0.58 | 43.75 | 3.62 | 900 | 249 | 15 | 11,268 | 3,119 |
| Mag. | 82.50 | 0.58 | 47.85 | 3.62 | 2,650 | 733 | 8 | 14, 385 | 3,416 |
| Unbarricaded Building | 103.6 | 0.79 | 81.80 | 4.34 | 7,920 | 1,828 | 8 | 15,052 | 3,472 |
| Unbarricaded Building | 117.0 | 0.58 | 67.90 | 4.08 | 2, 250 | 552 | 15 | 12,991 | 3, 190 |
| Unbarricaded Building | 124.0 | 0.79 | 98.00 | 4.60 | 3,000 | 652 | 15 | 14,976 | 3,254 |
| Barricaded Igloo | 125.0 | 0.87 | 108.8 | 4.77 | 3,950 | 828 | 13 | 15,893 | 3,335 |
| Earth Covered Igloo | 147.0 | 1.20 | 176.3 | 5.60 | 6,000 | 1,071 | 13 | 19,023 | 3,398 |
| Barge | 150.8 | 0.79 | 119.0 | 4.91 | 17,920 | 3,650 | 20 | 15,383 | 3,131 |
| Unbarricaded Mag. $50 \times 150$ | 173.1 | 0.79 | 136.9 | 5.15 | 1,800 | 350 | 25 | 15,510 | 3, 015 |
| Frame Bldg. and RR Cars | 177.6 | 0.58 | 103.0 | 4.69 | 7.920 | 1,691 | 15 | 15,268 | 3,262 |
| Magazine | 179.2 | 0.58 | 104.0 | 4.70 | 2, 400 | 511 | 8 | 16,241 | 3,492 |
| Unbarricaded Building | 193.0 | 0.42 | 81.05 | 4.32 | 5,280 | 1,220 | 15 | 13,921 | 3,221 |
| Freight Cars | 200.0 | 1.00 | 200.0 | 5.85 | 5,280 | 903 | 5 | 21,193 | 3,630 |
| Unbarricaded Mag. | 225.5 | 0.79 | 178.2 | 5.62 | 7,920 | 1,410 | 8 | 19,879 | 3,539 |
| Unbarricaded Dump | 270.0 | 1.00 | 270.0 | 6.46 | 2,400 | 372 | 30 | 19,650 | 3,046 |
| Steamer | 300.0 | 0.79 | 237.0 | 6.19 | 15,840 | 2,560 | 20 | 20, 118 | 3,257 |
| Underground Bunkers | 350.0 | 1.00 | 75.0 | 4.21 | 1,500 | 356 | 24 | 12,189 | 2,894 |
| Steamer | 400.0 | 1.25 | 500.0 | 7.93 | 9,000 | 1,135 | 20 | 26,663 | 3,366 |
| Steamer | 544.0 | 1.10 | 599.0 | 8.42 | 4,500 | 535 | 20 | 28,498 | 3,389 |
| Ba=ricaded Bldg. $500 \times 25$ | 550.0 | 1.17 | 643.5 | 8.63 | 7,300 | 846 | 12 | 30,533 | 3, 544 |
| Unbarricaded Mag. | 587.5 | 0.58 | 341.0 | 6.99 | 13,200 | 1,890 | 15 | 23,872 | 3,424 |
| Steamer | 400.0 | 1.00 | 400.0 | 7.36 | 3,900 | 530 | 20 | 24, 532 | 3, 336 |
| Unbarricaded Mag. | 800.0 | 1.00 | 800.0 | 9.28 | 5,280 | 569 | 15 | 32,495 | 3,508 |
| Building | 1,000.0 | 1.00 | 1,000.0 | 10.00 | 5, 280 | 528 | 15 | 35, 185 | 3,527 |
| Steamer | 1,000.0 | 1.00 | 1,000.0 | 10.00 | 6,600 | 660 | 20 | 34, 395 | 3,447 |
| Steamer | 730.0 | 0.42 | 306.5 | 6.74 | 5, 280 | 784 | 20 | 22, 192 | 3, 297 |
| Steamer | 2,280.0 | 0.42 | 958.0 | 9.85 | 11,500 | 1,168 | 20 | 33, 857 | 3,443 |
| Steamer | 2, 136.0 | 1.00 | 2,136.0 | 12.88 | 13,000 | 1,010 | 20 | 45,193 | 3,518 |
| Steamer | 2,600.0 | 1.00 | 2,600.0 |  | 18,480 |  | 20 | 50,528 |  |
| Barricade | 2,670.0 | 1.00 | 2,670.0 | 13.86 | 4, 346 | 314 | 20 | 48,924 | 3,536 |
| Open Pile | 4,500.0 | 0.42 | 1,890.0 | 12.34 | 4,920 | 399 | 70 | 35,737 | 2,898 |

\begin{tabular}{|c|c|c|c|c|c|c|c|c|c|c|c|c|}
\hline  \& \&  \& - \& \& \(\longrightarrow\) \&  \& 2n \& m- \& \%emb \& \(\cdots\) \& \& \\
\hline  \& \& MAXIMU \& DEBRIS DISTANCE
FOR SELEC \& \[
\begin{aligned}
\& E A N D \\
\& {[E D E}
\end{aligned}
\] \& \[
\begin{aligned}
\& \text { EXPL } \\
\& \text { XPLOS }
\end{aligned}
\] \& OSION SIONS \& PARA \& ETERS \& \& \& \& \\
\hline 菏 \& \[
\stackrel{\text { gh }}{\text { L }}
\] \&  \&  \&  \&  \&  \&  \&  \&  \&  \&  \&  \\
\hline \begin{tabular}{l}
King Powder Co. \\
King'a Mill, Ohio \\
Furnace, Lochlyne, Scotland Essex, Ontario
\end{tabular} \& 1954
1942
1883
1907 \& Black Powder
Black Powder
Black Powder
Dynamite \& Barricaded Building
Barricaded Building
Unbarricaded Building
Freight Car \& 2.500
2.500
2.500
2.500 \& 0.58
0.58
0.58
0.79 \& 1.450
1.450
1.450
1.975 \& 1.130
1.130
1.130
1.252 \& 850
400
1,350
500 \& 752
354
1.195
399 \& 15
15
15
5 \& 2,187
2,187
2,187
3,944 \& 1,933
1,933
1,933
3,144 \\
\hline Offley \& \& \& Truck \& 2.600 \& 1.10 \& 2.860 \& 1. 419 \& 1,845 \& 1.301 \& 4 \& 4,712 \& 3,321 \\
\hline Schaghticoke, N.Y. \& 1904 \& Black Powder \& Unbarricaded Building \& 2.900 \& 0.58 \& 1.682
1.740 \& 1.188
1.202 \& 300
500 \& 253
415 \& 15 \& 2, 384 \& 2, 205
2,022 \\
\hline Schaghticoke, N.Y. \& 1910 \& Black Powder
Black Powder \& Unbarricaded Building
2 Sides Barricaded Building \& 3.000
3.000 \& 0.58
0.58 \& 1.740
1.740 \& 1.202
1.202 \& 500
500 \& 415
415 \& 15 \& 2,431
2,431 \& 2,022
2,022 \\
\hline Fairchance, Pa. \& 1942 \& Black Powder
Bi-Oil \& 2 Sides Barricaded Building
Building \& 3.000
2.650 \& 0.58
1.35 \& 1.740
\(\mathbf{3 . 7 5}\) \& 1.202
1.528 \& 500
3,750 \& 2,455 \& 15 \& 3,431 \& 2,022
2,343 \\
\hline \& 1943 \& Torpex \& Bomb Trailer \& 3.025 \& 1.25 \& 3.781 \& 1.555 \& 1,000 \& 644 \& 5 \& 5,231 \& 3, 360 \\
\hline Hazardville, Conn. \& 1913 \& Black Powder \& Unbarricaded Building \& 3. 350 \& 0.58 \& 1.939 \& 1.246 \& 800 \& 642 \& 15 \& 2,587 \& 2, 074 \\
\hline Wilmington, Del. \& 1909 \& Black Powder \& Barricaded Building \& 3.350
3.510
3.7 \& 0.58
0.58
0. \& 1.939
2.035 \& 1.246
1.265 \& 900
1,500 \& 723
1,186 \& 15
15 \& 2,587
2,655 \& 2,074
2,096 \\
\hline |l \({ }^{\text {Faversham, England }}\) Sunflower Ordnance Wks. \& 1879 \& Black Powder
NG \& Unbarricaded Building
Barricaded Building \& 3.510
3.700 \& 0.58
1.35 \& 2.035
4.985 \& 1.265
1.707 \& 1,500
2,425 \& 1,186
1,422 \& 15
15 \& 2,655
4,228 \& 2,096
2,475 \\
\hline Kenvil, N.J. \& 1943 \& NG \& Barricaded Building \& 3.750 \& 1.35 \& 5.06 \& 1.715 \& 2,640 \& 1.540
4.080 \& 15 \& 4, 256
2,754
4,256 \& 2,480
2,126 \\
\hline Hercules, Calif. \& 1953 \& Black Powder \& Unbarricaded Building \& 3.750
3.750 \& 0.58 \& 2. 175 \& 1.295
1.715 \& 5,280
1.400 \& 1,080
816 \& 15 \& 2, 754
4,256 \& 2,126
2,480 \\
\hline Badger Ordnance Wks., Wis. \& 1945 \& \(\underset{\text { Picric Acid }}{ }\) \& Barricaded Building
Scrap Pile \& 3.750
3.750 \& 1.35 \& 5.060
4.160 \& 1.715
1.607 \& 1,400
100 \& 8187
187 \& 15
5
15 \& 4, 256 \& 2,480 \\
\hline Marugame. Shikoku, Japan Barksdale, Wis. \& 1945 \& Picric Acid
NG \& Scrap Pile \(\begin{aligned} \& \text { Barricaded Building }\end{aligned}\) \& 3.750
3.750 \& 1.35 \& 5.60 \& 1.715 \& 1,500 \& 875 \& 15 \& 4, 256 \& 2,480 \\
\hline Hounslow, England \& 1888 \& Black Powder
Black Powder \& Unbarricaded Building \& 3.750
4.000 \& 0.58
0.58 \& 2.175
2. 320 \& 1.295
1.322 \& 930
1.125 \& 319
850 \& 15 \& 2,754
2,853 \& 2,126
2,156 \\
\hline Mi. Carmel, Calif. Selma, N. C. \& 1907 \& Black Powder
Tetryl and TNT \& Unbarricaded Building \& 4.000
4.000
4.680 \& 0.58
1.14 \& 2.320
4.560
3.698 \& 1.322
1.657
1.544 \& \(\begin{array}{r}1.125 \\ \hline 825 \\ \hline 100\end{array}\) \& \(\begin{array}{r}850 \\ 498 \\ \hline 160\end{array}\) \& 15
4
10 \& 2,853
5,608
4,306 \& 2, 126
3, 384
2.786 \\
\hline Mahanoy \({ }^{\text {City, }}\) C. Pa. \& \& Dynamite \& Tractor-Trailer
Barricaded Building \& 4.680
4.835 \& 1.79
1.35 \& 3.698
6.525 \& 1.544
1.866 \& 2. 100 \& 1,360
268 \& 10 \& 4,306
4,806 \& 2,786
2,573 \\
\hline Barksdale, Wis. \& 1906 \& NG \& Barricaded Building \& 4.835 \& 1.35 \& 6.525 \& 1.866 \& \& 268 \& 15 \& 4,806 \& \\
\hline Pittsburgh, Pa.
Fairchance, Pa. \& 11894 \& Dynamite
Black Powder \& Unbarricaded Mag.
Barricaded Building \& 5.000
5.000 \& 0.79
0.58 \& 3.950
2.900 \& 1.580
1.425
1.910 \& 1,000
1,050 \& 633
737 \& 8
15 \& 4,718
3,214 \& 2,986
2,254
2,597 \\
\hline Fairchance, Pa.
Seneca, 11. \& 1944 1949 \& Brack Powder
NG \& Barricaded Building \& 5.174 \& 1.35 \& 6.990 \& 1.910 \& 5,280 \& 2, 765 \& 15 \& 4,961 \& 2,597 \\
\hline Antwerp, Beigium \& 1889 \& Black Powder \& Unbarricaded Building \& 6.000 \& 0.58 \& 3. 480 \& 1.512 \& 2,640
3,000 \& 1,745
1,473 \& 15 \& 3, 532
5,425 \& 2,331
2,663 \\
\hline Kenvil, N.J. \& 1948 \& NG \& Barricaded Building \& 6.275 \& 1.35 \& 8.470 \& 2,037 \& 3,000 \& 1,473 \& 15 \& 5,425 \& 2,663 \\
\hline Wayne, N.J. \& \& Black Powder
Di-Nitrol-Phenol \& Unbarricaded Building
Barricaded Building \& 6.438
6.700 \& 0.58
0.96 \& 3.735
6.435 \& 1.550
1.860 \& 2,000
1,525 \& 1,290
820 \& 15
15 \& 3,660
4,773 \& 2,360
L, 568 \\
\hline Rainbow Factory, Essex \& 1916 \& Di-Nitrol-Phenol \& Barricaded Building
Freight Car \& 6.700
6.800 \& 0.58
1.90 \& 6. 435
6.800
7.850 \& 1.850
1.893 \& 1,525
2,000 \& 1.820
1.058 \& 15
5 \& 3,773
6,340
6,649 \& 2,568
3,348
3,438 \\
\hline La Jolla, Calif. \& 1945 \& Ammo. \& Ammo. Van \& 7.250 \& 1.00 \& 7.250 \& 1.935
2.003 \& 1,675 \&  \& 4 \& 6,649
6,918 \& 3.438
3,449 \\
\hline Nebraska \& 1953 \& TNT \& Trailer \& 8.085 \& 1.00 \& 8.085 \& 2.003 \& 1,000 \& 499 \& 4 \& 6,918 \& 3,449 \\
\hline \begin{tabular}{l}
McAleater, Okla. \\
Allendorf, Germany
\end{tabular} \& 1944 \& Torpex
Torpex \& Trailer in Front of lgloo
Unbarricaded Building \& 8.100
9.000 \& 1.25
1.25 \& 10.12
11.25 \& 2.162
2.240 \& 2,500

900 \& 1,158
402 \& 4
15 \& 7.505
6.166 \& 3,472
2.754

2 <br>
\hline Allendorf, Germany City Point, Va. \& 1944 \& Torpex ${ }^{\text {Black Powder }}$ \& Unbarricaded Building \& 8.000 \& 0.58 \& 1.250
4.640 \& 2.667 \& 1,500 \& 900 \& 20 \& 3,462 \& $\begin{array}{r}2 \\ \hdashline . \\ \hline\end{array}$ <br>
\hline Rradford, England Kelloge, 11 . \& 1917 \& Picric Acid
Black Powder \& Unbarricaded Building
Barricaded Building \& 9.500
10.00 \& 1.11
0.58 \& 10.54
5.800 \& 2.196
1.796 \& 750
3,300 \& 352
1,840 \& 15
15 \& 5,990
4,544 \& $\therefore .734$
-.531 <br>
\hline Kellogg, $\mathbf{~ I l}$. \& \& Black Powder
Black Powder \& Barricaded Building
Barricaded Building \& 10.00 \& 0.58
0.58 \& 5.800
5.800 \& 1.796
1.796 \& 3,300
1,000 \& 1,840
557 \& 15 \& 4,544
4.544 \& 6.531
$\therefore .531$ <br>
\hline Holmes Park, Mo. Oakdale, Calif. \& 1908 \& Black Powder
Powder \& Barricaded Building
Unbarricaded Mag. \& 10.00
10.00 \& 0.58 \& 5. 800 \& 1.796 \& 1750 \& 418 \& 8 \& - 5.521 \& 3, 075 <br>
\hline Ce Elum, Wash. \& 1908 \& Black Powder \& Unbarricaded Mag.
3 Sides Barricaded Building \& 10.02
10.50 \& 0.58
0.58 \& 5.815
6.095 \& 1.798
1.827 \& 1,400
6,600 \& 779
3,614 \& 8
15 \& ¢
+
4.626
4.651 \& 3,075
2,548 <br>

\hline | Wilpen, Minn. |
| :--- |
| Wilmington,- Del. | \& 1909 \& Black Powder

Black Powder \& 3 Sides Barricaded Building Unbarricaded Building \& 10.50
10.59 \& 0.58
0.58 \& 6.095
6.140 \& 1.827
1.830 \& 6,600
75 \& 3,614
41 \& 15

15 \& | 4,651 |
| :--- |
| 4,659 | \& 2, 548

2,551 <br>
\hline Oakdale, Pa. \& \& TNT \& Building
Loading Platform \& 12.00
12.00 \& 1.00
1.25 \& 12.00
15.00 \& 2.286
2.465 \& 2,640
2,000 \& 1,155 \& \& \&  <br>
\hline Oahu, Hawaii
Fontanet, Ind. \& 11944 \& Torpex ${ }^{\text {Black Powder }}$ \& Loading Platform \& 12.00
12.50 \& 1.25
0.58 \& 15.00
7.250

7.250 \& | 2.465 |
| :---: |
| 1.935 | \& 2.000

375 \& 194 \& 15 \& * \& $\cdots$ <br>
\hline Fontanet,
Charleston \& . \& Black Powder \& Barricaded Mags. \& 12.50 \& 0.58 \& 7. 250 \& 1.935
2.180 \& 1,200
1.500 \& 621 \& 8
15 \&  \& - <br>
\hline Pindle, Calif. \& 1908 \& Dynamite \& Barricaded Building \& 13.18 \& 0.79 \& 10.40 \& 2. 180 \& 1,500 \& 688 \& 15 \& \& <br>

\hline Newburgh Heights, Cleveland, Ohio \& 1912 \& Dynamite Black Powder \& Unbarricaded Mags. \& $$
\left\{\begin{array}{l}
13.38 \\
13.50
\end{array}\right.
$$ \& 0.69

1.00 \& 9.235
13.50 \& 2.097
2.380 \& 1,800
447 \& 859
188 \& \& \% \& \% <br>

\hline | Gascoigne Woods |
| :--- |
| Reddick, ml . |
| Navajo Ordnance Depot | \& 1907 \& | TNT |
| :--- |
| Dynamite |
| Pentolite | \& | Unbarricaded |
| :--- |
| Freight Car Igloo | \& ( 13.50 \& 1.00

0.79
1.07 \& 13.50 \& 2.380
2. 280
2. 540 \& 747
7.920
3,600 \& 188
3,474
1,418

7 \& $$
\begin{array}{r}
15 \\
5 \\
13
\end{array}
$$ \& \[

$$
\begin{aligned}
& 780 \\
& 7.56
\end{aligned}
$$

\] \& \[

2,978
\] <br>

\hline
\end{tabular}

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|  |  |  |  |  |  |  |  |  | $\begin{aligned} & 088 \infty 8 \\ & \text { inminim } \\ & \text { nin } \end{aligned}$ |  |  |  |
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|  |  |  |  | 䂞号 |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  |  |  |  |  |  |



## MUM DEBRIS DISTANCE AND EXPLOSION PARAMETERS FOR SEIECTED EXPLOSIONS

|  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Building | 81.59 | 0.58 | 43.75 | 3.62 | 900 | 249 | 15 | 11,268 | 3,119 |
| Mag. | 82.50 | 0.58 | 47.85 | 3.62 | 2,650 | 733 | 8 | 14,385 | 3,416 |
| Unbarricaded Building | 103.6 | 0.79 | 81.80 | 4.34 | 7,920 | 1,828 | 8 | 15, 052 | 3,472 |
| Unbarricaded Building | 117.0 | 0.58 | 67.90 | 4.08 | 2, 250 | 552 | 15 | 12,991 | 3, 190 |
| Unbarricaded Building | 124.0 | 0.79 | 98.00 | 4.60 | 3, 000 | 652 | 15 | 14,976 | 3, 254 |
| Barricaded lgloo | 125.0 | 0.87 | 108.8 | 4.77 | 3,950 | 828 | 13 | 15,893 | 3,335 |
| Earth Covered Igloo | 147.0 | 1.20 | 176.3 | 5.60 | 6,000 | 1,071 | 13 | 19,023 | 3,398 |
| Barge | 150.8 | 0.79 | 119.0 | 4.91 | 17,920 | 3,650 | 20 | 15, 383 | 3,131 |
| Unbarricaded Mag. $50 \times 150$ | 173.1 | 0.79 | 136.9 | 5.15 | 1,800 | 350 | 25 | 15,510 | 3, 015 |
| Frame Bldg, and RR Cars | 177.6 | 0.58 | 103.0 | 4.69 | 7,920 | 1,691 | 15 | 15,268 | 3,262 |
| Magazine | 179.2 | 0.58 | 104.0 | 4.70 | 2,400 | 511 | 8 | 16,241 | 3,492 |
| Unbarricaded Building | 193.0 | 0.42 | 81.05 | 4.32 | 5,280 | 1,220 | 15 | 13,921 | 3,221 |
| Freight Cars | 200.0 | 1.00 | 200.0 | 5.85 | 5, 280 | 903 | 5 | 21,193 | 3,630 |
| Unbarricaded Mag. | 225.5 | 0.79 | 178.2 | 5.62 | 7.920 | 1,410 | 8 | 19,879 | 3, 539 |
| Unbarricaded Dump | 270.0 | 1.00 | 270.0 | 6.46 | 2,400 | 372 | 30 | 19,650 | 3,046 |
| Steamer | 300.0 | 0.79 | 237.0 | 6.19 | 15,840 | 2,560 | 20 | 20, 118 | 3,257 |
| Underground Bunkers | 350.0 | 1.00 | 75.0 | 4.21 | 1,500 | 356 | 24 | 12, 189 | 2,894 |
| Steamer | 400.0 | 1.25 | 500.0 | 7.93 | 9,000 | 1,135 | 20 | 26, 663 | 3,366 |
| Steamer | 544.0 | 1.10 | 599.0 | 8.42 | 4,500 | 535 | 20 | 28, 498 | 3,389 |
| Barricaded Bldg. $500 \times 25$ | 550.0 | 1.17 | 643.5 | 8.63 | 7,300 | 846 | 12 | 30,533 | 3,544 |
| Unbarricaded Mag. | 587.5 | 0.58 | 341.0 | 6.99 | 13,200 | 1,890 | 15 | 23, 872 | 3,424 |
| Steamer | 400.0 | 1.00 | 400.0 | 7.36 | 3,900 | 530 | 20 | 24,532 | 3,336 |
| Unbarricaded Mag. | 800.0 | 1.00 | 800.0 | 9.28 | 5, 280 | 569 | 15 | 32,495 | 3,508 |
| Building | 1,000.0 | 1.00 | 1,000.0 | 10.00 | 5,280 | 528 | 15 | 35, 185 | 3,527 |
| Steamer | 1,000.0 | 1.00 | 1,000.0 | 10.00 | 6,600 | 660 | 20 | 34, 395 | 3,447 |
| Steamer | 730.0 | 0.42 | 306.5 | 6.74 | 5,280 | 784 | 20 | 22, 192 | 3,297 |
| Steamer | 2, 280.0 | 0.42 | 958.0 | 9.85 | 11,500 | 1,168 | 20 | 33, 857 | 3,443 |
| Steamer | 2,136.0 | 1.00 | 2,136.0 | 12.88 | 13,000 | 1,010 | 20 | 45, 193 | 3,518 |
| Steamer | 2,600.0 | 1.00 | 2,600.0 |  | 18,480 |  | 20 | 50, 528 |  |
| Barricade | 2,670.0 | 1.00 | 2,670.0 | 13.86 | 4,346 | 314 | 20 | 48, 924 | 3,536 |
| Open Pile | 4,500.0 | 0.42 | 1,890.0 | 12.34 | 4,920 | 399 | 70 | 35,737 | 2,898 |

## SECRET

APPENDIX B COMPUTER PROGRAM FOR REGRESSION STUDY

OF HE EXPLOSIONS

## SECRET


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COMPUTER PROGRAM FOR REGRESSION STUDY OF HE EXPLOSIVE

## SECRET

APPENDIX C
MAXIMUM DEBRIS DISTANCE AND CRATER DIMENSIONS FOR SELECTED EXPLOSIONS

## SECRET




## SECRET

APPENDIX D<br>FRAGMENT-SIZE DISTRIBUTION AND DISPERSION<br>OF FRAGMENTS<br>FROM PANTEX ORDNANCE PLANT EVENT

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## SECRET

FRAGMENT-SIZE DISTRIBUTION FOR REINFORCED CONCRETE STRUCTURE (Pantex Ordnance Plant Event)

| $\begin{aligned} & \text { Fragment } \\ & \text { Size, } \\ & \text { (lb) } \end{aligned}$ | $\begin{gathered} \text { Number } \\ \text { of } \\ \text { Fragments } \end{gathered}$ | Percent of <br> Total Fragments | Cumulative <br> Number of Fragments | Cumulative Percent of Total Fragments | ```Weight of Fragments. (lb)``` | Percent of Total Weight | Cumulative Weight of Fragments, (1b) | Cumulative Percent of Total Weight |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1/16 | 56 | 0.180 | 56 | 0.18 | 3.5 | 0.0041 | 3.5 | 0.004 |
| 1/8 | 906 | 2.916 | 962 | 3.10 | 113.25 | 0.133 | 16.75 | 0.137 |
| $1 / 4$ | 18, 138 | 58.374 | 19,100 | 61.47 | 4,534.5 | 5. 324 | 4,651.25 | 5.461 |
| $1 / 3$ | 12 | 0.039 | 19, 112 | 61.51 | 4.0 | 0.005 | 4,655. 25 | 5.466 |
| 1/2 | 40 | 0.129 | 19,152 | 61.64 | 20.0 | 0.023 | 4,675.25 | 5.490 |
| 3/4 | 5,858 | 18.853 | 25,010 | 80.49 | 4,393,5 | 5.159 | 9,068.75 | 10.64 |
| 1 | 32 | 0.103 | 24,042 | 80.59 | 32.0 | 0.038 | 9.100 .75 | 10.69 |
| 1-1/4 | 1 | 0.003 | 25, 043 | 80.60 | 1.25 | 0.001 | 9.102 .0 | 10.69 |
| 1-1/2 | 10 | 0.032 | 25,053 | 80.63 | 15.0 | 0.018 | $9,117.0$ | 10.71 |
| 2 | 2, 502 | 8.052 | 27,555 | 88.68 | 5,004,0 | 5.876 | 14,121.0 | 16.58 |
| 2-1/2 | 8 | 0.026 | 27, 563 | 88.71 | 20.0 | 0.023 | 14,141.0 | 16.60 |
| 3 | 2,680 | 8.625 | 30, 243 | 97.33 | 8,040.0 | 9.440 | 22,181.0 | 26.04 |
| 3-1/3 | 1 | 0.003 | 30, 244 | 97.34 | 3.33 | 0.004 | 22,184.0 | 26.05 |
| 3-1/2 | 14 | 0.045 | 30,258 | 97.38 | 49.0 | 0.058 | 22,233.0 | 26.11 |
| 4 | 12 | 0.039 | 30,270 | 97.42 | 48.0 | 0.056 | 22,281.0 | 26. 16 |
| 5 | 2 | 0.006 | 30, 272 | 97.43 | 10.0 | 0.012 | 22, 291.0 | 26.17 |
| 6 | 13 | 0.042 | 30, 285 | 97.47 | 78.0 | 0.092 | 22,369.0 | 26.27 |
| 7 | 1 | 0.003 | 30, 286 | 97.47 | 7.0 | 0.008 | 22, 376. 0 | 26.27 |
| 8 | 3 | 0.010 | 30,289 | 97.48 | 24.0 | 0.028 | 22.400 .0 | 26. 30 |
| 9 | 1 | 0.003 | 30,290 | 97.48 | 9.0 | 0.011 | 22,409.0 | 26.31 |
| 10 | 7 | 0.023 | 30, 297 | 97.51 | 70.0 | 0.082 | 22,479.0 | 26.39 |
| 12 | 106 | 0.341 | 30,403 | 97.85 | 1,272.0 | 1.493 | 23,751.0 | 27.89 |
| 14 | 13 | 0.042 | 30,416 | 97.89 | 182.0 | 0.214 | 23,933.0 | 28. 10 |
| 15 | 2 | 0.006 | 30,418 | 97.90 | 30.0 | 0.035 | 23,963.0 | 28. 14 |
| 18 | 9 | 0.029 | 30,427 | 97.92 | 162.0 | 0.190 | 24,125.0 | 28. 33 |
| 20 | 1 | 0.003 | 30,428 | 97.93 | 20.0 | 0.023 | 24, 145.0 | 28. 35 |
| 40 | 4 | 0.129 | 30,432 | 97.94 | 160.0 | 0.188 | 24, 305.0 | 28. 54 |
| 50 | 1 | 0.003 | 30,433 | 97.94 | 50.0 | 0.059 | 24,355.0 | 28.60 |
| 60 | 240 | 0.772 | 30,673 | 98.72 | 14,400.0 | 16.908 | 38.755.0 | 45.51 |
| 70 | 17 | 0.055 | 30,690 | 98.77 | 1,190.0 | 1.397 | 39,945.0 | 46.90 |
| 75 | 41 | 0.132 | 30.731 | 98.90 | 3,075.0 | 3.610 | 43,020.0 | 50.51 |
| 80 | 1 | 0.003 | 30,732 | 98.91 | 80.0 | 0.094 | 43, 100.0 | 50.61 |
| 90 | 248 | 0.798 | 30,980 | 99.70 | 22,320.0 | 26.208 | 65,420.0 | 76.82 |
| 95 | 2 | 0.006 | 30,982 | 99.71 | 190.0 | 0.223 | 65,610.0 | 77.04 |
| 100 | 6 | 0.019 | 30,988 | 99.73 | 600.0 | 0.705 | 66, 210.0 | 77.74 |
| 150 | 1 | 0.003 | 30,989 | 99.73 | 150.0 | 0.176 | 66, 360.0 | 77.92 |
| 180 | 81 | 0.261 | 31,070 | 99.994 | 14,580.0 | 17.120: | 80,940.0 | 95.04 |
| 225 | 1 | 0.003 | 31,071 | 99.997 | 225.0 | 0.264 | $81,165.0$ | 95.30 |
| 4000 | 1 | 0.003 | 31,072 | 100.000 | 4,000.0 | 4.697 | 85,165 | 100.00 |

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fragmentation data on exploned dry sandstone blocks

| Test: <br> Shothole Diameter: <br> Diameter of Explosive: | B <br> $3 / 4$ in. <br> 5/8 in. |  |  | Block Dimensions: $18-\mathrm{in}$ diameter $\times 30-\mathrm{in}$. long <br> Explosive: XL Hawkite <br> Weight of Explosive: 20 z |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Size Range of Fragments | Average Diameter, (in.) |  |  | Average <br> Weight per <br> Fragment, <br> (1b <br> 62. | Number of Fragments in Class | $\qquad$ | Cumulative <br> Percent of Total <br> Fragments | Minimum <br> Fragment Weight in Class, (1b) |
| Over 8 inch <br> 4-8 inch <br> 2-4 inch <br> 1-2 inch <br> 1/2-1 inch <br> $1 / 4-1 / 2$ inch <br> 1/16-1/4 inch | $(12)$ 6 3 1.5 0.72 0.37, | 57, 400 <br> 44,700 <br> 43, 800 <br> 31,800 <br> 20, 500 <br> 11,100 12,300 <br> 12,300 | 126.5 98.5 96.5 70.1 45.2 24.5 | 62.9 7.86 0.983 0.1229 0.01535 0.00192 | $\begin{array}{r} 2.013 \\ 12.51 \\ 98.15 \\ 571.0 \\ 2,945.0 \\ 12.760 .0 \end{array}$ | $\begin{array}{r} 2.013 \\ 14.52 \\ 112.7 \\ 684.0 \\ 3,629.0 \\ 16.389 .0 \end{array}$ | 0.0124 0.0887 0.688 4.175 22.18 100.0 | 18.68 <br> 2.332 <br> 0.2916 <br> 0.03644 <br> 0.00455 0.00057 |
| 500 microns $-1 / 16$ inch <br> 251-500 microns <br> 124-251 microns <br> 66-124 microns <br> 20-66 microns <br> 10-20 microns <br> 5-10 microns <br> 2-5microns <br> Under 2 microns |  | $\begin{array}{r} 3,330 \\ 5,920 \\ 5,820 \\ 2,030 \\ 1,060 \\ 200 \\ 159 \\ 122 \\ 68 \end{array}$ |  |  |  |  |  |  |
| Total Weight of Recovered Material |  | 240,309 | 530.0 |  |  |  |  |  |
| Total Weight of Material above 1/4-in. Fragment Size |  | $\begin{gathered} 209,300 \\ (87.1 \%) \end{gathered}$ | 461.3 |  |  |  |  |  |
| Original Volume of Material (less shothole) cu.in. |  | 7612 |  |  |  |  |  |  |
| Computed Density of Material, $\mathrm{lb} / \mathrm{cu}$. in. |  |  | 0.0696 |  |  |  |  |  |

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FRAGMENTATION DATA ON EXPLODED DRY SANDSTONE BLOCKS

| Test: <br> Shothole Diameter: <br> Diameter of Explosive: | 3/4 in. <br> $5 / 8 \mathrm{in}$. |  |  | Block Dimensions: $18-\mathrm{in}$. diameter $\times 30-\mathrm{in}$. <br> Explosive: Colexg <br> Weight of Explosive: 1 oz |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Size Range of Fragments | Average Diameter, (in.) | Total Class Weight, (gm) | Total Class Weight, ( 1 b ) | Average Weight per Fragments, (lb) | Number of <br> Fragments in Class | Cumulative <br> Number of <br> Fragments | Cumulative Percent of Total Fragments | Minimum Fragment Weight in Class,(1b) |
| Over 8 inch <br> 4-8 inch <br> 2-4 inch <br> 1-2 inch <br> 1/2-1 inch <br> 1/4-1/2 inch <br> 1/16-1/4 inch <br> 500 microns - $1 / 16$ inch <br> 251-500 microns <br> 124-251 microns <br> 66-124 microns <br> 20-66 mirrons <br> 10-20 microns <br> 5-10 microns <br> 2-5 microns <br> Under 2 microns | $\begin{aligned} & (12) \\ & 6 \\ & 3 \\ & 1.5 \\ & 0.75 \\ & 0.375 \end{aligned}$ | $\begin{array}{r} 94,400 \\ 60,100 \\ 38,600 \\ 25,200 \\ 6,690 \\ 4,170 \\ 3,490 \\ 1,560 \\ 2,230 \\ 1,810 \\ 567 \\ 262 \\ 74 \\ 67 \\ 42 \\ 14 \end{array}$ | 208.0 132.5 85.1 55.5 14.7 9.2 | $\begin{aligned} & \hline 62.8 \\ & 7.85 \\ & 0.981 \\ & 0.1226 \\ & 0.0153 \\ & 0.0019 \end{aligned}$ | $\begin{array}{r} 3.32 \\ 16.89 \\ 86.7 \\ 452.0 \\ 964.0 \\ 4,835.0 \end{array}$ | $\begin{array}{r} 3.32 \\ 20.21 \\ 106.9 \\ 559.0 \\ 1,523.0 \\ 6,358.0 \end{array}$ | $\begin{gathered} 0.0522 \\ 0.318 \\ 1.68 \\ 8.78 \\ 23.96 \\ 100.0 \end{gathered}$ | 18.6 <br> 2. 325 <br> 0.2906 <br> 0.0363 <br> 0.0045 <br> 0.00056 |
| Total Weight of Recovered Material |  | 239, 276 | 528.0 |  |  |  |  |  |
| Total Weight of Material above $1 / 4$-in. Fragment Size |  | $\begin{aligned} & 229,160 \\ & (95.8 \%) \end{aligned}$ | 505.0 |  |  |  |  |  |
| Original Volume of Material (less shothole), cu.in. |  | 7,612 |  |  |  |  |  |  |
| Computed Density of Material, lb/cu.in. |  |  | 0.0694 |  |  |  |  |  |

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FRAGMENIATOV DATA OVEXHLODED DF: SANDSTONE RLOCKS

| Size Range of Fragments | Average Diameter (in.) | Total <br> Class <br> Weight, ( gm ) | Total Class Weight, <br> ( 16 ) | Average Weight per Fragment, (1b) | Number of Fragments in Class | Cumulative Number of Fragments | Cumulative Percent of Total Fragments | Minimum Fragment Weight in Class,(1b) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Over 8 inch | (12) | 102, 000 | 225.0 | 66.2 | 3.40 | 3.40 | 0.058 | 19.6 |
| 4-8 inch | 6 | 55, 000 | 121.2 | 8.28 | 14.65 | 18.05 | 0.308 | 2.45 |
| 2-4 inch | 3 | 46,000 | 101.5 | 1.035 | 98.1 | 116.15 | 1.982 | 0.3066 |
| 1-2 inch | 1.5 | 29,900 | 66.0 | 0.1293 | 511.0 | 627.0 | 10.7 | 0.0384 |
| 1/2-1 inch | 0.75 | 7,600 | 16.75 | 0.0162 | 1,036.0 | 1,663.0 | 28.3 | 0.0048 |
| 1/4-1/2 inch | 0.375 | 3, 860 | 8.51 | 0.00202 | 4,205.0 | 5,868.0 | 100.0 | 0.0006 |
| 1/16-1/4 inch |  | 3,180 |  |  |  |  |  |  |
| 500 microns - $1 / 16$ inch |  | 1,440 |  |  |  |  |  |  |
| 251-500 microns |  | 1,510 |  |  |  |  |  |  |
| 124-251 microns |  | 1,560 |  |  |  |  |  |  |
| 66-124 microns |  | 495 |  |  |  |  |  |  |
| 20-66 microns |  | 260 |  |  |  |  |  |  |
| 10-20 microns |  | 66 |  |  |  |  |  |  |
| 5-10 microns |  | 51 32 |  |  |  |  |  |  |
| 2-5 microns <br> Under 2 microns |  | 32 18 |  |  |  |  |  |  |
| Total Weight of Recovered Material |  |  |  |  |  |  |  |  |
|  |  | 252,972 | 557.2 |  |  |  |  |  |
| Total Weight of Material above !/4 in. Fragment Size |  |  |  |  |  |  |  |  |
|  |  | 244, 360 | 539.0 |  |  |  |  |  |
|  |  | (96.6\%) |  |  |  |  |  |  |
| Original Volume of Material (Less shothole), cu.in. |  | 7,616 |  |  |  |  |  |  |
| Computed Density of Material, $\mathrm{lb} / \mathrm{cu}$. in. |  |  | 0.0731 |  |  |  |  |  |

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FRACMENTA:OVDAPAONEXFLODED DEY SANDSTOVE BLOCKS

| Test: <br> Shothole Diameter: <br> Diameter of Explosive: | $\begin{aligned} & G \\ & 1-1 / 4 \mathrm{in} . \\ & 1-1 / 8 \mathrm{in} . \end{aligned}$ |  | Block Dimensions: 18-in <br> Explosive: <br> Weight of Explosive: |  |  | diameter $\times$ 30-in. long <br> Rounkol <br> 1 oz |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Size Range of Fragments | Average Diameter, (in.) |  |  | Average Weight per Fragment, (1b) | Number of Fragrents in Class | $\qquad$ | Cumulative <br> Percent of Total <br> Fragments | Minimum <br> Fragment Weight in Class, (1b) |
| Over 8 inch <br> 4-8 inch <br> 2-4 inch <br> 1-2 inch <br> 1/2-1 inch <br> 1/4-1/2 inch <br> 1/16-1/4 inch <br> 500 microns $-1 / 16$ inch <br> 251-500 microns <br> 124-251 microns <br> 66-124 microns <br> 20-66 microns <br> 10-20 microns <br> 5-10 microns <br> 2-5 microns <br> Under 2 microns | $\begin{aligned} & (12) \\ & 6 \\ & 3 \\ & 1.5 \\ & 0.75 \\ & 0.375 \end{aligned}$ | $\begin{array}{r} 58,400 \\ 115,000 \\ 35,900 \\ 21,700 \\ 9,430 \\ 3,330 \\ 2,690 \\ 1,330 \\ 1,470 \\ 1,430 \\ 389 \\ 189 \\ 45 \\ 23 \\ 21 \\ 12 \end{array}$ | 128.8 25.6 79.1 47.8 20.8 7.3 | $\begin{gathered} 64.9 \\ 8.25 \\ 1.031 \\ 0.129 \\ 0.016 \\ 0.002 \end{gathered}$ | 1.98 30.74 76.7 371.0 $1,290.0$ $3,670.0$ | $\begin{gathered} 1.98 \\ 32.72 \\ 109.0 \\ 480.0 \\ 1,770 \\ 5,440.0 \end{gathered}$ | $\begin{gathered} 0.0368 \\ 0.602 \\ 2.004 \\ 8.83 \\ 32.55 \\ 100.0 \end{gathered}$ | $\begin{aligned} & 19.54 \\ & 2.445 \\ & 0.3056 \\ & 0.0382 \\ & 0.00475 \\ & 0.00059 \end{aligned}$ |
| Total Weight of Recovered Material |  | 251,359 | 554.0 |  |  |  |  |  |
| Total Weight of Material above $1 / 4$ in. Fragment Siz: |  | $\begin{aligned} & 243,760 \\ & (97.0 \%) \end{aligned}$ | 537.4 |  |  |  |  |  |
| Original Volume of Material (less shothole), cu.in. |  | 7,488 |  |  |  |  |  |  |
| Computed Density of Material, lb/cu.in. |  |  | 0.0730 |  |  |  |  |  |

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FRAGMESTAT:ON DATA ON EXPLODED DRY SANDSTONE BLOCKS

| Test: <br> Shothole Diameter: <br> Diameter of Explosive: | $\begin{aligned} & \mathrm{K} \\ & 1-1 / 4 \mathrm{in} . \\ & 5 / 8 \mathrm{in} . \end{aligned}$ |  |  | Block Dimensions: 18-in <br> Explosive: <br> Weight of Explosive: |  | ```diameter x 30-in. long Colex 1 oz``` |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Size Range of Fragments | Average Diameter, (in.) | Total Class Weight, $(\mathrm{gm})$ | $\begin{aligned} & \hline \text { Total } \\ & \text { Class } \\ & \text { Weight, } \\ & (\mathrm{lb}) \end{aligned}$ | $\begin{aligned} & \text { Average } \\ & \text { Weight per } \\ & \text { Fragment, } \\ & \text { (lb) } \end{aligned}$ | Number of Fragments in Class |  | $\begin{gathered} \hline \text { Cumulative } \\ \text { Percent of } \\ \text { Total } \\ \text { Fragments } \end{gathered}$ | Minimum Fragment Weight in Class, (1b) |
| Over 8 inch <br> 4-8 inch <br> 2-4inch <br> 1-2 inch <br> 1/2-1 inch <br> $1 / 4-1 / 2$ inch <br> 1/16-1/4 inch | $\begin{aligned} & (12) \\ & 6 \\ & 3 \\ & 1.5 \\ & 0.75 \\ & 0.375 \end{aligned}$ | $\begin{array}{r} 119,000 \\ 53,300 \\ 30,600 \\ 14,300 \\ 4,080 \\ 1,810 \\ 1,470 \end{array}$ | 262.1 117.5 67.5 31.5 9.0 4.0 | $\begin{aligned} & 59.6 \\ & 7.45 \\ & 0.931 \\ & 0.1163 \\ & 0.01456 \\ & 0.00182 \end{aligned}$ | $\begin{gathered} 4.395 \\ 15.79 \\ 72.45 \\ 271.0 \\ 618.0 \\ 2,191.0 \end{gathered}$ | $\begin{array}{r} 4.40 \\ 20.19 \\ 92.64 \\ 364.0 \\ 982.0 \\ 3.173 .0 \end{array}$ | 0.139 0.636 2.92 11.47 30.92 100.0 | 17.67 2.208 0.276 0.0345 0.00431 0.00054 |
| 500 microns - $1 / 16$ inch <br> 251-500 microns <br> 124-251 microns <br> 66-124 microns <br> 20-66 microns <br> 10-20 microns <br> 5-10 microns <br> 2-5 microns <br> Under 2 microns |  | $\begin{array}{r} 588 \\ 652 \\ 597 \\ 174 \\ 118 \\ 34 \\ 20 \\ 12 \\ 8 \end{array}$ |  |  |  |  |  |  |
| Total Weight of Recovered Material |  | 226,763 | 499.5 |  |  |  |  |  |
| Total Weight of Material above $1 / 4-\mathrm{in}$. Fragment Size |  | $\begin{aligned} & 223,090 \\ & (98.4 \%) \end{aligned}$ | 491.6 |  |  |  |  |  |
| Original Volume of Material (less shothole), cu.in. |  | 7,588 |  |  |  |  |  |  |
| Computed Density of Material, lb/cu. in. |  |  | 0.0659 |  |  |  |  |  |

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# CONCRETE FRAGMENT WEIGHTS AND DIMENS:ONS FOR 1/24-SCALE SHIELDED REACTOR MODELS 

Model Nu. 3
Event Simulated, 730 lb 「NI', 1 masec, w/Pyrocore Vessel $3 / 4$ Full of Water

| Piece <br> Number | Individual Weipht, |  | Group Weight, |  | Height, (in.) | $\begin{aligned} & \text { Length, } \\ & \text { (in.) } \end{aligned}$ | $\begin{aligned} & \text { Wicith, } \\ & (\text { in. }) \end{aligned}$ | Dist. of Fragments from Original Model Position, If $>10 \mathrm{ft}$ |  | Reinarlis |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1 b | oz | 1b | Oz |  |  |  | - It | in. |  |
| 44 | 39 | 15 | 39 | 15 | 6 |  | 12 |  |  | Ring, 12-in. diam. with 4-1/2-in. diam. hole |
| 43 | 11 | 2 | 11 | 2 | 7 | 9 | 4 |  |  |  |
| 1 | 8 | 3 |  |  | 5 | 8 | 4 |  |  |  |
| 4 | 8 | 1 | 16 | 4 | 6 | 6 | 4 |  |  |  |
| 2 | 5 | 7 |  |  | 6 | 7 | 4 |  |  |  |
| 6 | 5 | 7 |  |  | 6 | 4 | 4 |  |  |  |
| 5 | 5 | 4 |  |  | 4 | 5 | 4 |  |  |  |
| 8 | 5 | 1 | 21 | 3 | 6 | 5 | 4 |  |  |  |
| 30 | 4 | 5 | 4 | 5 | 6 | 4 | 4 |  |  |  |
| 16 | 3 | 11 |  |  | 4 | 5 | 4 |  |  |  |
| 3 | 3 | 9 |  |  | 4-3/4 | 4 | 4 |  |  |  |
| 15 | 3 | 4 |  |  | 6 | 6 | 4 |  |  |  |
| 32 | 3 | 1 | 13 | 9 | 5-1/4 | 4 | 4 |  |  |  |
| 12 | 2 | 15 |  |  | 4 | 5 | 4 |  |  |  |
| 9 | 2 | 14 |  |  | 4 | 4 | 4 |  |  |  |
| 7 | 2 | 13 |  |  | 4-1/2 | 4 | 4 |  |  |  |
| 11 | 2 | 8 |  |  | 2-1/2 | 5-1/2 | 4 |  |  |  |
| 37 | 2 | 5 |  |  |  | 3 | 4 |  |  |  |
| 24 | 2 | 4 | 15 | 11 | 4-1/2 | $3-1 / 2$ | 4 |  |  |  |
| 23 | 1 | 15 |  |  | 4 | $3-1 / 2$ | 3 |  |  |  |
| 21 | 1 | 14 |  |  | 5 | $2-1 / 2$ | 4 | 14 | 5 |  |
| 17 | 1 | 13 |  |  | 4 | 4 | 3 |  |  |  |
| 10 | 1 | 12 |  |  | 5 | 2-1/2 | 4 |  |  |  |
| 20 | 1 | 9 |  |  |  |  | 2-1/2 |  |  |  |
| 22 | 1 | 9 |  |  | 2-1/2 | $2-1 / 2$ | 2 |  |  |  |
| 27 | 1 | 9 |  |  | 2 | 4 | 4 |  |  |  |
| 28 | 1 | 7 |  |  | 4 | 2-1/2 | 4 |  |  |  |
| 31 | 1 | 7 |  |  | 4-3/4 | 2-1/2 | 4 |  |  |  |
| 13 | 1 | 5 |  |  | 5-3/4 | 2 | 4 |  |  |  |
| 14 | 1 | 5 |  |  | 2-1/2 | 4 | 4 |  |  |  |
| 29 | 1 | 5 |  |  | 2 | $2-1 / 2$ | 4 |  |  |  |
| 19 | 1 | 4 |  |  | 3 | 3 | 3 |  |  |  |
| 33 | 1 | 4 |  |  | 4-3/4 | $2-1 / 2$ | 4 |  |  |  |
| 18 | 1 | 3 | 22 | 9 | 4 | $2-1 / 2$ |  |  |  |  |
| 40 |  | 15 |  |  | 4 | 2 | 2-3/4 |  |  |  |
| 26 |  | 14 |  |  | 2 | 3-1/2 | 2-1/2 |  |  |  |
| 36 |  | 13 |  |  | 4-1/2 | 1-3/4 |  |  |  |  |
| $+1$ |  | 12 |  |  | 4-1/2 | 1-1/2 | 2 |  |  |  |
| 25 |  | 11 |  |  | 2-1/2 | 1-1/2 |  |  |  |  |
| 34 |  | 10 |  |  | 3/4 | 2 | $2-1 / 2$ |  |  |  |
| 38 |  | 10 |  |  | 4 | 1-1/2 | 3 |  |  |  |
| 35 |  | 9 |  |  | 3-1/2 | 2 |  |  |  |  |
| 39 |  | 7 |  |  | 4-1/2 | 1 | 2-3/4 |  |  |  |
| 42 |  | 7 |  |  | 1-1/2 | 4 |  |  |  |  |
| 45 |  | 6 |  |  | 1-1/2 | 2-1/4 | 1-1/2 | 18 |  |  |
| 46 |  | $1 / 4$ |  |  | 3/4 | 1 | 1/4 | 19 | 6 |  |
| - | 4 | 1-3/4 | 11 | 4 |  |  |  |  |  | Approximately <br> 50 small <br> chunks of conceret. |

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CONCRETE FRAGMENT WEIGHTS AND DIMENE.ONS Model No. 5 FOR 1/24-SCALE SHIELDED REACTOR MODELS Event Simulated, 510 lb TNT, $1 \mathrm{msec}, \mathrm{w} /$ Pyrocore Vessel Full of Water

| Piece Number | Individual Weight, |  | Group Weight, |  | Height, (in.) | $\begin{aligned} & \text { Length, } \\ & \text { (in.) } \end{aligned}$ | Width, (in.) | Remarks |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | lb | oz | 1b | oz |  |  |  |  |
| 9 | 32 |  | 32 |  | 4 | 15 | 10 |  |
| 16 | 23 | 12 |  |  | 4 | 15 | 10 |  |
| 18 | 23 | 12 | 47 | 8 | 4 | 15 | 8 |  |
| 17 | 17 |  | 17 |  | 4 | 12 | 6 |  |
| 2 | 8 | 6 | 8 | 6 | 4 | 10 | 4 |  |
| 6 | 6 | $\stackrel{2}{14}$ | 6 5 | $\stackrel{2}{14}$ | 4 | 8 10 | 3 |  |
| 8 1 | 5 4 | 14 14 | 5 | 14 | 4 | 10 6 | 3 4 |  |
| 21 | 4 | 14 |  |  | 4 | 6 | 4 |  |
| 7 | 4 | 14 |  |  | 4 | 4 | 6 |  |
| 22 | 4 | 8 |  |  | 2 4 | 4 | 4 |  |
| 9 3 | 4 4 4 | 4 | 27 | 8 | 4 4 | 4 | 4 |  |
| 5 | 3 | 12 |  |  | 4 | 6 | 4 |  |
| 4 | 3 | 6 |  |  | 4 | 6 | 3 |  |
| 10 | 3 | 4 | 10 | 6 | 4 | 4 | 6 |  |
| 11 | 2 | 4 | 2 | 4 | 4 | 4 | 3 |  |
| 20 | 1 | 12 4 |  |  | 4 | 4 |  |  |
| 14 |  | 12 | 3 |  | 4 | 3 | $2^{1-1 / 2}$ |  |
| 13 |  | 12 |  |  | 3 | 2 |  |  |
| 23 |  | 6 |  |  | 3 | 3 | 1/2 |  |
| 15 24 |  | 8 | 3 | 8 | 1-1/2 | 1-1/2 | 1/2 |  |
|  | 1 | 8 |  | 8 |  |  |  | 40 small of concrete |

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## CONCRETE FRAGMENT WEEGHTS AVD [)!ME*VSONS FOR 1,24-SCALE SHIELDED REACTOR MODE: S

Model No. 14
Event Simulated, 160 lb TNT, 」msec, w/Pyrocore Vessel 3/4 Full of Water

| Piece Number | Individual Weight, |  | Group Weight, |  | Height,(in.) | $\begin{aligned} & \text { Length, } \\ & \text { (in.) } \end{aligned}$ | Width, (in.) | Remarks |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1b | OZ | 1b | Oz |  |  |  |  |
| 1 | 140 |  | 140 |  | 18-1/2 | 12 | 12 | Cylinder, 12-in. diam. with 4-1/4-in. diam. hole |
| 2 | 19 | 14 | 19 | 14 | 3-1/2 | 12 | 12 |  |

Model No. 15
Event Simulated, 150 lb TNT, l msec, w/Pyrocore Vessel 3/4 Full of Water

| Piece <br> Number | Individual Weight, |  | Group Weight, |  | Height, (in.) | Length, (in.) | Width, (in.) | Remarks |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1b | Oz | 1 b | Oz |  |  |  |  |
| 15 | 142 |  | 142 |  | 18 |  | 12 | Cylinder, with 4-1/4-in. diam. hole at center |
| 1 | 3 | 8 |  |  | 2 | 6 | 5 |  |
| 2 | 3 | 8 | 7 |  | 2-1/4 | 7 | 5 |  |
| 3 | 1 | 15 |  |  |  | 5-1/4 |  |  |
| 6 | 1 | 8 |  |  | 1-1/2 | 7-1/2 | 3-1/4 |  |
| 5 | 1 | 6 |  |  | 2 | 4-1/2 | 3 |  |
| 7 | 1 | 6 |  |  | 1-3/4 | 5 | 2-1/2 |  |
| 4 | 1 | 5 |  |  | 1-3/4 | 6 | 2-1/4 |  |
| 8 | 1 |  | 8 | 8 | 1-3/4 | 5-3/4 | $2-1 / 2$ |  |
| 9 |  | 13 |  |  | 2 | 2-3/4 | $\therefore-3 / 4$ |  |
| 11 |  | 8 |  |  | 1-1/2 | 4 | 2-1/2 |  |
| 12 |  | 6 |  |  | 1-1/2 | 3-1/2 |  |  |
| 10 |  | 5 |  |  | 1-1/2 | 3-1/4 |  |  |
| 13 |  | 5 |  |  | 1-3/4 |  | 1-1/2 |  |
| - |  | 7 | 2 | 12 |  |  |  | 10 small chunks of concrete |

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CONCRETE FRAGMENT WEIGHTS AND D:MENS:ONS Model No. 16 FOR 1/ट4-SCALE SHIELDED REACTOR MODELS

Event Simulated, 210 b rNT, 1 maec, w/MDF


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## CONCRETE FRAGMENT WEIGHTS AND DIMENSIONS

 FOR 1/24-SCALE SHIELDED REACTOR MODELSModel No. 17
Event Simulated, 160 lb TNT, 1 msec , w/Pyrocore Vessel Full of Water

| Piece Number | Individual Weight, |  | Group Weight, |  | Height, (in.) | $\begin{gathered} \text { Length, } \\ \text { (in.) } \end{gathered}$ | Width,(in.) | Remarks |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | lb | oz | lb | oz |  |  |  |  |
| 1 | 145 |  | 145 |  | 18-1/2 | 12 | 12 | Cylinder, 12-in. diam. with 4-1/4-in, diam. hole |
| 3 | 12 | 4 | 12 | 4 | 3 | 12 | 6 |  |
| 4 | 4 | 1 | 4 | 1 | 2-1/2 | 6-1/2 | 5-3/4 |  |
| 6 | 1 | 11 | 1 | 11 | 2-1/4 | 6 | 2-1/2 |  |
| 2 |  | 7 |  |  | 1-1/2 | 2-1/2 | 2-1/4 |  |
| 7 |  | 4 |  |  | 3/4 | 2 | 2-1/2 |  |
| 5 |  | 2 |  |  |  | 2 | 3/4 |  |
| 9 |  | 1-1/2 |  |  | 1/2 | 2 | 1-3/4 |  |
| 8 |  | 1 |  | 15-1/2 | 1/2 | 2-1/2 | 1-1/4 |  |

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## CONCRETE FRAGMENT WEIGHTS AND DIMENSIONS FOR 1/24-SCALE SHIELDED REACTOR MODELS

Mori-1 No. 19
Event Simulated, 210 lb [NT, 1 mese, w/MDF


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APPENDIX G<br>CONCRETE FRAGMENT WEIGHTS AND DIMENSIONS FOR 1/12-SCALE SHIELDED REACTOR MODELS

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CONCRETE FRAGMENT WEIGHTS AND DIMENSIONS FOR 1/12-SCALE SHIELDED REACTOR MODELS

| Energy Source: | Pyrocore |
| :--- | :--- |
| Period: | 1 msec |
| Water in Vessel: | $100 \%$ |


| Model No. | A | B | C | D | E | F |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Event Simulated |  |  |  |  |  |  |
| TNT (lb) | 150 | 150 | 400 | 400 | 650 | 650 |
| Megawatt-Seconds | 280 | 280 | 750 | 750 | 1230 | 1230 |
| Model Wt. (lb) | $\begin{aligned} & 1308 \\ & \text { (avg.) } \end{aligned}$ | $\begin{aligned} & 1308 \\ & \text { (avg.) } \end{aligned}$ | 1292 | 1311 | 1286 | 1342 |
| \% Fragments |  |  |  |  |  |  |
| $<11 \mathrm{~b}$ | -- | -- | 1.5 | 0.9 | 1.3 | 1.4 |
| 1-5 lb | -- | -- | 1.4 | 0.5 | 3.4 | 1.4 |
| 5-10 lb | -- | -- | 1.3 | 1.7 | 6.4 | 3.2 |
| 10-15 lb | -- | -- | 1.9 | 1.1 | 11.4 | 6.1 |
| 15-20 lb | -- | -- | 1.4 | 3.9 | 2.6 | 3.6 |
| 20-25 lb | 2 | -- | 8.7 | 13.2 | 10.5 | 8.0 |
| $>25 \mathrm{lb}$ | 98 | 100 | 83.9 | 78.6 | 64.4 | 76.1 |
| No. Fragments |  |  |  |  |  |  |
| $<1 \mathrm{lb}$ | --* | -- | 31(1616) | 22(1050) | 46(120)* | 47(99) |
| 1-5 lb | -- | -- | 10 | 4 | 21 | 8 |
| 5-10 lb | -- | -- | 2 | 3 | 11 | 6 |
| 10-15lb | -- | -- | 2 | 1 | 11 | 7 |
| 15-20 1b | -- | -- | 1 | 3 | 2 | 3 |
| 20-25 lb | 1 | -- | 5 | 8 | 6 | 5 |
| $>25 \mathrm{lb}$ | 4 | 6 | 12 | 14 | 20 | 20 |
| INo. Fragments Excluding Fines | 5 | 6 | 63 | 55 | 117 | 96 |

*Figures in parentheses indicate "Fines".

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## APPENDIX H <br> TRAJECTORY OF AN AIR PARTICLE <br> DURING THE POSITIVE PHASE OF A NUCLEAR BLAST

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## APPENDIX H <br> TRAJECTORY OF AN AIR PARTICLE DURING THE POSITIVE PHASE OF A NUCLEAR BLAST

Using the 1957 "The Effects of Nuclear Weapons" (Ref. 20), the velocity-time history of an air particle may be found at any fixed point on the ground surface for a surface burst of any yield. If we know the velocity of an air particle $u_{1}$ at one location, we can predict its new location after a short time interval $\Delta t_{1}\left(e . g .\right.$, old location $\left.+u_{1} \Delta T_{1}\right)$. Repeating this procedure enables us to establish the horizontal distance-time relationship of the particle. Because so many factors enter into the determination of the particle velocity as presented in reference 20, an otherwise straightforward integration becomes a rather involved bookkeeping problem. The following outline indicates the steps involved in the numerical integration:
I. Select
A. Weapon Yield
B. Overpressure at front of forest. $\mathrm{P}_{1}$
II. Compute (Time: $T=0$ )
A. Shock Velocity $U_{1}$ for $p_{1}$ (see Fig. 3. 80, Ref. 20)
B. Particle Velocity $u_{1}$ for $p_{1}$ (see Fig. 3. 80, Ref. 20)
C. Distance from ground zero (GZ) to front of forest: $x_{0}$

1. Fig. 3. 94a; l-KT surface burst; Ref. 20
2. Scaling Law; Eq. 3. 86. 1 Ref. 20
III. Select $T=\Delta T_{1}$
IV. Compute
A. Distance traveled by particle in $\Delta \mathrm{T}_{1} \Delta \mathrm{x}_{1}=\mathrm{u}_{1} \Delta \mathrm{~T}_{1}$
B. Distance traveled by shock front in $T_{1} \Delta y_{1}=U_{1} \Delta T_{1}$
C. Time gap between shock front and partıcle:

$$
\mathrm{t}_{1}=\left(\mathrm{U}_{1}-\mathrm{u}_{1}\right) \quad \Delta \mathrm{T}_{1} / \mathrm{U}_{1}
$$

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D. Overpressure at distance $\mathrm{x}_{\mathrm{o}}+\Delta \mathrm{x}_{1}$ from GZ: $\mathrm{p}_{2}^{\prime}$

1. Fig. 3. 94a; 1-KT surface burst; Ref. 20
2. Scaling Law; Eq. 3. 86. 1 Ref. 20
E. Overpressure at distance $\mathrm{x}_{\mathrm{o}}+\Delta \mathrm{y}_{1}$ from GZ: $\mathrm{p}_{2}$
3. Fig. 3. 94a; 1-KT Surface burst; Ref. 20
4. Scaling Law; Eq. 3. 86. 1 Ref. 20
F. Shock Velocity $U_{2}$ for $p_{2}$ (see Fig. 3. 80, Ref. 20)
G. Duration of positive phase at $\mathrm{x}_{\mathrm{o}}+\Delta \mathrm{x}_{1}: \mathrm{t}_{1+}$
5. Fig. 3. 96; surface burst; Ref. 20
6. Scaling from Art. 3. 88, Ref. 20
H. Overpressure behind shock front $p\left(t_{1}\right)$
7. Compute $t_{1} / t_{1+}$
8. Compute $p\left(t_{1}\right)$ from Eq. 3. 82.1 (Ref. 20) using $P_{2}^{\prime}$
I. Particle Velocity $u_{2}$ for $\mathrm{p}\left(\mathrm{t}_{1}\right)$; (Fig. 3.80 Ref. ${ }^{20}$ )
V. Select $\left(T=\Delta T_{2}\right)$
VI. Compute
A. $\Delta \mathrm{x}_{2}=\mathrm{u}_{2} \Delta \mathrm{~T}_{2}$
B. $\Delta y_{2}=U_{2} \Delta T_{2}$
C. $\mathrm{t}_{2}=\left[\Delta \mathrm{y}_{1}+\Delta \mathrm{y}_{3}-\left(\Delta \mathrm{x}_{1}+\Delta \mathrm{x}_{2}\right)\right] / \mathrm{U}_{2}$
D. $P_{3}^{\prime}$ at $x_{0}+\Delta x_{1}+\Delta x_{2}$ from $G Z$
E. $p_{3}$ at $x_{0}+\Delta y_{1}+\Delta y_{2}$ from $G Z$

The curves in the figures referred to in the outline were all fit with analytical expressions and the entire procedure was programmed for the UNIVAC 1105. Figure H-l shows the detailed flow diagram of the program

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Figure H-1 Detaıled Flow Diagram For The Partıcle Velocity Integration

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## APPENDIX J <br> THE VULNERABILITY OF ANTENNA SYSTEMS

The severity of the debris hazard to antenna systems is only one phase of debris problems associated with hardened sites, but it is used to describe the estimating procedures that can be applied to debris problems in general, using data collected in this report.

The approach used in this study made estimates in several manners and noted their consistency. First, estimates of the maximum range of antenna vulnerability were based on the correlation of maximum debris distance and equivalent yield observed in Chapter Two. Debris distribution data from the Pantex Ordnance Plant event were then scaled up to nuclear yields. This was done by scaling up the ground ranges according to the cube-root-of-yield scaling, although it may have been more appropriate to also scale up the total volume of fragments in a manner which accounted for the total material volume of the nuclear crater. The result was perhaps optimistic, resulting in fewer fragments. Next, debris environment was estimated independently by scaling up DANNY BOY findings to high yields. This was done by scaling both ground ranges and fragment densities according to cube-root-of-yield scaling. A third estimate of debris environment was based on use of the hydrodynamic model of crater formation developed by Brode and Bjork at RAND Corporation ${ }^{*}$. An analytical solution to the debris environment was obtaned by considering peak velocities from the hydrodynamic model to be initial fragment velocities, and following the trajectories of the ejecta from the crater to ultimate impact with the ground. Results of these three approaches are consistent in predicting severe debris environments under likely hardening criteria.

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There are two potential sources of debris hazards to antenna systems: throwout material from the crater and loose material or broken structure picked up and transported by the blast winds. The second source of debris can be eliminated by clearing an area around the antenna. No simple defense against crater throwout debris is apparent. This problem is therefore restricted to prediction of the debris hazard resulting from the crater throwout alone.

A comprehensive method of analysis of debris effects emanating from nuclear detonations has yet to be developed. To a great extent the character of the hardened sites is responsible for this state of affairs. The hardened missile silo with a reinforced concrete cover as the only exposed component was considered relatively invulnerable to debris damage. The sole problem was to predict the weight of debris on the closure after an attack so that sufficient power could be designed into the closure operating mechanism. Communications systems, however, pose entirely different problems. Components of antenna systems, for example, may be quite vulnerable to debris damage because of their electrical essentials. Thus, the hardened antenna design must be based on such criteria as density, size, and energy of debris particles to be expected at the antenna location.

Some studies which border on, or are corollary to, the problem have been performed. These include studies of missiles from accidental explosions (Chapter Two), studies of debris distribution from nuclear tests (JANGLE U and DANNY BOY), and analytical studies dealing with the crater formation problem (RAND). Each of these is used to estimate the debris hazard for the hypothetical antenna structure.

It should be emphasized that at present, it is only possible to assess the debris problem very roughly. Our intention is to objectively review the available data and to compare predictions based on different source material. Hopefully, the results will be consistent. We do not contend that this study completely settles the debris question for hardened sites. We do feel, however, that it is the most reasonable approach to the problem within the current state-of-the-art, -- i. e., short of further nuclear testing.

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## J-1

## Estimates Based on High-Explosive Debris Data

The Armed Services Explosive Safety Board data or 206 HE detonations ranging in magnitude from $8-1 \mathrm{~b}$ of tetrytol to $9,000,000 \mathrm{lb}$ of ammonium nitrate were collected (Chapter Two). A statistical analysis relating the maximum missile distance to the weight of explosive (TNT equivalent) was made. On a $\log -\log$ plot a linear regression line to best fit the data was found to be of the form,

$$
\begin{equation*}
\log _{10} D_{M}=2.950+0.322 \log _{10} \mathrm{~W} \tag{J-1}
\end{equation*}
$$

where

$$
\begin{aligned}
& \mathrm{D}_{\mathrm{M}}=\text { maximum missile distance, ft } \\
& \mathrm{W}=\text { equivalent weight of } \mathrm{TNT}, \mathrm{lb} .
\end{aligned}
$$

Similarly, a quadratic regression line to best fit the data was found to be:

$$
\begin{equation*}
\log _{10} D_{M}=2.960+0.347 \log _{10} W-0.016\left(\log _{10} W\right)^{2} \tag{J-2}
\end{equation*}
$$

Note that Eq. ( $\mathrm{J}-1$ ) indicates one should scale according to $\mathrm{w}^{0.322^{\circ}}$ : A standard error was found to be $2.47 \mathrm{D}_{\mathrm{M}}$. Equations ( $\mathrm{J}-1$ ) and (J-2) were applied to the antenna vulnerability problem using an equivalence factor of 0.50 to relate nuclear yield to TNT equivalent; the results are presented in Table J-l.

Table J-1
LIMIT OF ANTENNA VULNERABILITY BASED ON HE DATA

| Weapon <br> Yield, <br> (MT) | Linear <br> Maximum <br> Debris <br> Distance, <br> (miles) | Range of Maximum <br> Debris Distance for <br> One Standard Error <br> (miles) | Maximum <br> Debris <br> Distance, <br> (miles) | Range of Maximum <br> Debris Distance for <br> One Standard Error <br> (miles) |
| :---: | :---: | :---: | :---: | :---: |
| 5 | 19.4 | 7.9 to 48.0 | 6.4 | 2.6 to 15.7 |
| 10 | 24.2 | 9.8 to 59.9 | 6.8 | 2.8 to 16.8 |
| 20 | 30.3 | 12.3 to 75.0 | 7.6 | 3.1 to 18.9 |
| 50 | 40.7 | 16.5 to 100.7 | 8.4 | 3.4 to 20.8 |
| 100 | 50.9 | 20.6 to 121.0 | 9.1 | 3.7 to 22.4 |

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It is interesting to note that a data point corresponding to $5,200,000 \mathrm{lb}$ of TNT (2. $60-\mathrm{KT}$ nuclear weapon) was included and that the maximum missile distance for this point was 3.5 miles. Thus, we extrapolated four cycles from six cycles of data.

The relationship between the high explosion debris problem and the debris emanating from the crater of a nuclear weapon is certainly questionable. The HE debris generally results from buildings and equipment in the immediate vicinity of the explosions. The mechanism by which this debris is formed is different from mechanisms by which a crater is formed. Nuclear detonation is accompanied by rather substantial winds whereas in the HE explosion these winds are essentially absent. Nevertheless, it is desirable to take into consideration the large body of data that is available for the HE debris problem, particularly in view of the scarcity of nuclear test data. As a matter of fact, since essentially all debris associated with an HE detonation originates near the point of detonation, one can argue that this debris is, in fact, similar in origin to throwout debris.

The second item of interest is the distribution of debris outward from the point of the explosion. A detailed debris study which included complete descriptions and final locations of the debris was carried out at Pantex Ordnance Plant (Ref. 1l). The explosive was 2000 lb of TNT detonated in a reinforced concrete bunker. For application to the antenna vulnerability problem, all reported ranges were scaled up by the cube-root-of-yield law. The total volume of fragments at the Pantex study was estimated to be 1000 cu ft , whereas the crater volume for, say a $20-\mathrm{MT}$ weapon surface burst is about $3.7 \times 10^{9} \mathrm{cu} \mathrm{ft}$. Rather than using the ratio of $3.7 \times 10^{9} / 10^{3}$ to scale the number of fragments, we used the more optimistic (resulting in fewer fragments) cube-root-of-yield factor (215 for this case). The resulting missile density as a function of ground range is shown on Fig. J-1. The dashed portions of the curve are extrapolated to expected antenna locations. These results make it apparent that debris problems are critical for antenna systems.

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Fig': J-1 F: agmen! Distritution ficum 20-MT Weapron Based on Pantex Test

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## J-2 Estimates Based on Nuclear Test Results

Armour Research Foundation (Ref. 2) recently completed an experimental study of crater throwout from an underground nuclear detonation; these results are derived completely from that work. This study was part of Project 7 DANNY BOY, a 0.43-KT nuclear device buried at a depth of 110 ft in basalt on the Buckboard Mesa at the Nevada iest site. The expected crater zone was salted with a variety of objects which were located after the blast. In addition, the natural debris was tabulated in certain areas. Of particular interest here are the natural debris radial distribution charts reproduced from reference 2 in Fig. J-2, J- 3, and J-4. Areas I, II, and III in these figures correspond to different orientations with respect to ground zero. The results must be scaled up to the cases of interest here. The ground ranges are reasonably scaled by the cube-root-of-yield law. Scaling of the fragment density expressed as fragments-per-square-foot requires some discussion. If the total number of fragments is assumed proportional to the crater volume, then the number of fragments scale directily as the weapon yield (crater diameter and depth each scale as the cube root of yield). The area over which these fragments are distributed scales as length square or as yield to the two-thirds power. Therefore, the fragment density scales as $W / W^{2 / 3}$ or as cube-root of yield. The results oi Fig. J-2, J-3, and J-4 for the three sectors are averaged and scaled up to a l-KT weapon as a standard. The result is shown on Fig. !. • This is then used to determine debris density as function of range for any weapon yield in kilotons.

This has been done for an antenna system locition of 4.0 miles for selected weapon yields. Results are shown in Table J 2.

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Figure J-2 Throwout Lensity, Area I, DANNY BOY Event:

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Figure J-3 Throwout Density, Area II, DANNY BOY Event

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Figure.J-t Throwout Density. Area III, DANNY BOY Evert

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FigureJ-5 Debris Density for 1-KT Weapon Basfed on DANNY BOY Results

Table J-2
DEBRIS DENSITY FOR ANTENNA SYSTEM LOCATION
AT 4.0 MILES GROUND RANGE
(Based on extrapolation of DANNY BOY data)

| Weapon Yield, (MT) | Debris Density, (fragments/sq ft) |  |  |
| :---: | :---: | :---: | :---: |
|  | 1 in . to 6 in . Equivalent Diameter ( 0.0003 to $0.065 \mathrm{ft}^{3}$ ) | 6 in . to 12 in. Equivalent Diameter ( 0.065 to $0.523 \mathrm{ft}^{3}$ ) | 12 in . to 18 in. Equivalent Diameter (0. 523 to $1.77 \mathrm{ft}^{3}$ ) |
| 10 | 0. 86 | 0.00037 | 0.023 |
| 20 | 43. 0 | 0.171 | 0.610 |
| 50 | 1,260.0 | 41.3 | 11.5 |
| 100 | 11,400.0 | 1,410.0 | 72. 9 |

Based on the assumptions made with regard to scaling, it is obvious from Table J-2 that debris would pose a major problem. Note that particles of the size shown in Table $J-2$ were found in the DANNY BOY test. Thus, in applying the results of Table $J-2$ to antenna vulnerability problems, all three sizes of particles must be considered simultaneously. No attempt was made here to "scale" particle size which probably varies with both weapon yield and soil type.

These results are more severe than those predicted from HE data (Fig. J-2). Recall, however, that the number of fragments was scaled up by only the optimistic cube-root-of-yield scaling. Therefore. we feel the Table J-2 results are more significant.

J-3 Estimates Based on Analytical Studies
A completely analytical treatment of this problem depends primarily on a mathematical model to treat the crater formation problem. Such an analysis gives initial velocity vectors of material or particles leaving the crater, and the subsequent motion of these particles can be followed by means of standard trajectory analysis. We use a crater model devised at the RAND Corporation and then compute debris density at ranges from ground zero. The RAND model is first discussed, the trajectory analysis

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is then presented, and finally, results are presented which are applicable to antenna systems.

## J-3.1 Application of the RAND Crater Model

Brode and Bjork studied the formation of.a crater resulting from a 2 -MT weapon surface burst, assuming the material in the crater zone to be rock (tuff). Their interest was primarily in the early period of the crater formation when the pressure acting on the rock medium is very much greater than the shear strength of the rock. They, therefore, assumed the hydrodynamic model valid and numerically integrated the appropriate field equations. Pressure and velocity fields in the crater zone are presented. The velocity fields are reproduced here as Fig. J-6 through J-9.

The RAND model is formulated in Eulerian coordinates so that the velocity vectors represent the velocity of the mass currenily at the point in space indicated by the base of the vector. It is therefore not possible from the available data to rigorously follow the motion of a specific mass of crater material. Rather we used these data by assuming that the peak velocity at each point in the crater is the initial velocity of the mass at that point.

A grid was established to roughly cover the crater and the data of Fig. J-6 through J-9 were used to determine peak velocities at each point of grid. The grid is shown in Fig. J- 10 with grid points identified. The velocity vectors which are most severe from the viewpoint of throwout were selected; the resulting velocities are shown in Table J-3.

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Figure $J-6$ Velocity Field in Crater ( $T=0.1026 \mathrm{msec})$

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Figure - 7 Velocity Field in Crater $(T=52.49 \mathrm{msec})$

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Figure 7.9 Velocaty Field in Crater ( $T=105 \mathrm{msec})$

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Figure J-10 Crater Divisions

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Table J-3
INITIAL TABLE VELOCITY

| Grid Point (see <br> Fig. J-10 for Lo- <br> cation in Crater) | Initial Horizontal <br> Velocity, <br> (meters/msec) | Initial Vertical <br> Velocity, <br> (meters/msec) |
| :---: | :---: | :---: |
|  | 0.20 | 0.30 |
| 1 | 0.13 | 0.53 |
| 2 | 0.16 | 0.31 |
| 3 | 0.08 | 0.20 |
| 4 | 0.24 | 0.36 |
| 5 | 0.40 | 0.20 |
| 6 | 0.12 | 0.20 |
| 7 | 0.09 | 0.14 |
| 8 | 0.03 | 0.10 |
| 9 | 0.01 | 0.08 |
| 10 | 0.10 | 0.10 |
| 11 | 0 | 0 |
| 12 | 0 | 0 |
| 13 | 0 | 0 |
| 14 | 0 | 0 |
| 15 |  |  |

The velocity data in Table J-3 were used as the initial velocities in computing debris particle trajectories. Because the hydrodynamic model does not predict debris fragment size, trajectories were computed for a range of particle sizes.

## J-3. 2 Trajectory Analysis

Consider the motion of a particle through a medium such that the drag force acting on the particle $1 s$ proportional to the square of the relative velocity between the particle and the air. It, is assumed that the vertical and horizontal motion of the debris particle are decoupled. This is true if the center of pressure of the particle coincides with its centroid for all orientations so that no rotation occurs.

The equations of motion are then

$$
\begin{align*}
& \ddot{x}=\frac{k a \rho}{2}(u-\dot{x})^{2}  \tag{J-3}\\
& \ddot{y}=\frac{k^{\prime} a \rho}{2} \dot{y}^{2}+g \tag{J4}
\end{align*}
$$

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## where

$\mathrm{x}=$ horizontal coordinate
$y=$ vertical coordinate, positive downward
( ; = differentiation with respect to time
$a=$ aerodynamic coefficient, $\frac{\text { (projected area } x \text { drag coefficient) }}{\text { mass }}$
$\rho=$ mass density of air
$u=$ air particle velocity
$\mathrm{g}=\mathrm{gravitational}$ constant
$k=\left\{\begin{array}{l}+1 ; u>x \\ -1 ; u \leqslant x\end{array}\right.$
$k^{\prime}=\left\{\begin{array}{l}+1 ; y \leqslant 0 \\ -1, y>0\end{array}\right.$

These are Riccati-type nonlinear differential equations, and can be linearized by a simple transformation of coordinates.

The horizontal equation of motion ( $\mathrm{J}-3$ ) can be linearized by the substitution,

$$
x=-\frac{2}{k a \rho} \frac{\dot{s}}{s}
$$

Note that $x=-\frac{2}{k a \rho} \log _{e} s$,
thus,

$$
\begin{equation*}
\ddot{s}+k a \rho u \dot{s}+\frac{k^{2} a^{2} \rho^{2}}{4} u^{2} s=0 \tag{J-6}
\end{equation*}
$$

Equation ( $\mathrm{J}-6$ ) is a linear differential equation with variable coefficients because of the variation of particle velocity $u$ and air density $\rho$ with time. This could be numerically integrated but, because of our interest in relatively large times, the computation time would be prohibitive, and the resulting cumulative error undoubtedly sizable. Equation (J-6) is therefore solved by assuming that $u$ and $\rho$ are constant over an interval of time. The solution can then be extended in time by matching initial
conditions after each interval of time, and then changing $u$ and $\rho$ to a new constant value for the next interval. If $u$ is constant, Eq. (J-6) has a solution

$$
\begin{equation*}
s=\left(C_{1}+C_{2} t\right) \exp \left[-\frac{u a p^{k}}{2} t\right] \tag{J-7}
\end{equation*}
$$

where

$$
C_{1} \text {, and } C_{2} \text { are constants of integration. }
$$

The vertical equation-of-motion Eq. J-4 can be linearized by the substitution,

$$
\begin{equation*}
y=-\frac{z}{k^{\prime} a \rho} \frac{\dot{z}}{z} \tag{J-8}
\end{equation*}
$$

Then

$$
\begin{equation*}
y=-\frac{2}{k^{\top} \alpha \rho} \log _{e} z \tag{J-9}
\end{equation*}
$$

Then reduced equation then becomes,

$$
\begin{equation*}
\ddot{z}+k^{\prime} \frac{a_{\rho} \rho g}{2} z=0 \tag{J-10}
\end{equation*}
$$

Recalling that $k^{\prime}=+1$ for $y \leqslant 0$ and $k^{\prime}=-1$ for $\dot{y}>0$, Eq. $J-10$ has the solution,

$$
\begin{align*}
& z=C_{3} \cos \sqrt{\frac{a p g}{2}} t+C_{4} \sin \sqrt{\frac{a p g}{2}} t  \tag{J-11}\\
& \text { for } \dot{y} \leqslant o \text { (i. e., on way up) }
\end{align*}
$$

and

$$
\begin{aligned}
& z=C_{5} \exp \left[\sqrt{\frac{a \rho g}{2}}\right] t+C_{6} \exp \left[-\sqrt{\frac{a \rho g}{2}} t\right] \\
& \text { for } y>o \text { (i. e., on way down). }
\end{aligned}
$$

Equations ( $\mathrm{J}-11$ ) are used to compute the total time of flight for the particles, and then the total horizontal distance traveled is determined from the value of $s$, as computed from Eq. ( $J-7$ ) at the time the particle hits the ground.

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## J-3. 3 Numerical Results

The simultaneous solution of Eq. ( $\mathrm{J}-6$ ) and ( $\mathrm{J}-11$ ) was carried out on the UNIVAC 1105 digital computer. The initial velocities presented in Table J-3 were used as initial conditions and solutions were obtained for l-in., 6-in., and 12-in. diameter particles.

The analytic forms of the weapon parameters used in the computer program are taken from reference 20. The actual values as taken from the computer program for a $20-\mathrm{MT}$ weapon are plotted in Fig. J-11. For lack of better data, the close-in values for overpressure shock velocity and particle velocity were taken to be constant from ground zero out to a scaled ground range of 100 ft .

Solutions were obtained for particle velocity and air density assumed to be constant in $0.25-\mathrm{sec}$ and $0.1-\mathrm{sec}$ intervals. The numerical results differed by less than 2 percent so the computer runs were finally made using the $0.25-\mathrm{sec}$ interval. Numerical results for flight time, horizontal distance traveled and final velocity are given in Table J-4.

These data were then converted to fragment density values.
Consider the crater to be broken up into annular rings as shown in Fig. J-10. The material in each of the three horizontal layers was first distributed over the impact area. This was done by assuming that the material from each ring is spread at constant depth over a radial distance equal to the difference in the computed maximum trajectory distance of adjacent points. The fragment density in that region is then given as

$$
\begin{equation*}
N=\frac{V}{\frac{4}{3} \pi r_{f}^{3}} \frac{1}{\pi\left(r_{o}^{2}-r_{i}^{2}\right.} \tag{J-12}
\end{equation*}
$$

## where

$\mathrm{N}=$ fragment density
$\mathrm{V}=$ volume in crater as given in Fig. J-10
$r_{f}=$ fragment radius
$r_{o}=$ outer radial distance for ring of interest
$r_{i}=$ inner radial distance for ring of interest.

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Figure J-11 Weapon Parameter for 20-MT Weapon
Table J-4
TRAJECTGRY ANALYSIS FOR DATA FOR 20•MT WEAPON

| Point in Crater (See Fig. J-10 for location) | One-Inch Particles |  |  | Six-Inch Particles |  |  | Twelve-Inch Particles |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Final Location, (miles) | $\begin{aligned} & \text { Final } \\ & \text { Velocity, } \\ & \text { (fps) } \end{aligned}$ | Flight Time, (sec) | Final Location, (miles) | Final <br> Velocity, <br> (fps) | Flight Time, (sec) | Final Location, (miles) | $\begin{aligned} & \text { Final } \\ & \text { Velocity, } \\ & \text { (fps) } \end{aligned}$ | Flight Time, (sec) |
| 1 | 2. 9 | 1943 | 12.9 | 5.1 | 709 | 25.7 | 7.0 | 507 | 33.2 |
| 2 | 3.5 | 2131 | 14.4 | 5.5 | 723 | 29.4 | 7.5 | 504 | 38. 4 |
| 3 | 2.9 | 1972 | 12.9 | 5.1 | 710 | 25.9 | 7.0 | 510 | 33.5 |
| 4 | 2.5 | 1611 | 11.8 | 4.6 | 705 | 23.1 | 6.6 | 532 | 29.6 |
| 5 | 3.0 | 2069 | 13.3 | 5.2 | 711 | 26.9 | 7.3 | 498 | 34.9 |
| 6 | 2. 4 | 1580 | 11.8 | 4.6 | 696 | 23.1 | 7.0 | 505 | 29.6 |
| 7 | 2.5 | 1607 | 11.8 | 4.7 | 704 | 23.1 | 6.6 | 529 | 29.6 |
| 8 | 2.2 | 1336 | 10.9 | 4.4 | 702 | 20.8 | 6.2 | 553 | 26. 3 |
| 9 | 2.1 | 1137 | 10.0 | 4.0 | 703 | 18.7 | 5.9 | 585 | 23.3 |
| 10 | 2.0 | 1137 | 10.0 | 4.0 | 703 | 18.7 | 5.9 | 586 | 23.2 |
| 11 | 2.0 | 1132 | 10.0 | 4.0 | 701 | 18.7 | 5.9 | 583 | 23.2 |

The effects of the three layers were then added.
The results of Table J-4 were then transposed to fragment density-distance curves and, for consistency, scaled down to a l-KT weapon (both the range and fragment density were scaled as the cube root of yield). The result is plotted on Fig. J-12. It should be emphasized that the particle sizes on Fig. J-12 are not predicted by the model. Figure J-12 represents three solutions assuming that all of the crater material breaks up into the same size particles. Therefore, only one of the three curves in Fig. J-12 should be used at one tıme. This can be compared with the DANNY BOY results in Fig. J-5. The dashed line drawn in Fig. J- 12 is from the DANNY BOY results, assuming that all debris fragments break up into one-inch radus particles. It is not surprising that the RAND model gives more severe results because of the assumed hydrodynamic behavior of the soil and the neglect. of in-flight collisions between debris particles.

An interesting result was obtained while studying the trajectories of the debris leaving the crater. The total flight distance was found to be independent of the initial horizontal velocity component for the range of horizontal velocities predicted by the RAND model. In other words, the horizontal motion is determined completely by the blast winds for large weapons. The inttial vertical velocity, of course, determines the time of flight upon which the total throw of a particle is very dependent. Based on this brief analysis, it appears that for a megaton-yield weapon, the surface burst (or perhaps partially buried burst) results in the most severe debris problem.

## J-4 General Consistency of Results

Crater throwout debris was studied with particular emphasis on predicting its severity to antenna systems. The state-of-the-art is such that a completely reliable evaluation of this problem was not possible. The objective here was to examine avallable experimental data and analyses .. to establish bounds on the magnitude of the problem for antenna systems. First, the voluminous data which exist on debris resulting from high explosive detonation mastudied. Results are presented in Fig. J-l and Table J-1. It can be seen that the maximum missile distances predicted


Figure J-12 Debrıs Density for 1-KT Weapon Based on Brode's Model of Crater Formation
are beyond likely antenna locations, and that at expected antenna locations the debris density is quite high. Secondly, the results of the DANNY BOY nuclear test are extrapolated to make predictions for antenna systems. It is assumed that both distance and fragment density scale as the cube root of yield. The resulting fragment density-distance curve for a l-KT weapon is shown in Fig. J-5. This can be readily scaled up to real antenna situations and again a severe debris problem is predicted. Third, a completely analytical solution was obtained by following the motion of a particle from the crater to impact with the ground. Brode's hydrodynamic model was used for the crater formation phase of the motion. The resulting debris density-distance relationships, which indicated a relatively severe debris problem for expected antenna locations, are shown in Fig. J-12.

Therefore, while no conclusive evaluation of the problem was possible, all available data point to a very critical debris problem for antenna systems. There are two possible approaches to the design of antenna systems insofar as the debris problem is concerned.

First, locations could be restricted to those where the debris density was at a specified low level. Secondly, some degree of hardness could be provided against debris particles. At first look this would seem a most challenging task. The debris will have terminal velocities very close to the particle velocity, which can well be over 1000 fps . At these velocities it seems apparent that even quite small debris particles would be capable of damaging any fragile critical elements.

Certainly much has yet to be learned about the crater throwout problem. The soil type, depth of burial, and yield must all have some effect. None of this is understood at present. The results presented here are all based on a rock-like crater material. The debris particle size must certainly be a function of soil type. A sandy material should produce very small fragments that would tend to sand blast rather than fracture antenna elements. Also, as can be seen from Table J-4 sand would not be transported as far as a soil which breaks into large fragments.

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The extremes of depth of burial were included in this study. DANNY BOY represents a completely buried shot at optimum depth* whereas the Brode crater model is for a surface burst. As remarked earlier, it appears that the surface burst will give the more critical debris problem for megaton range weapons.

The dependence on weapon yield (scaling) represents perhaps the most important unknown. Most of our results depend on an assumed cube-root-of-yield scaling law, although the analytical results were only scaled up from a two-megaton weapon.

Perhaps the most surprising result of this study is the consistency of the results of the three methods studied when so many unknown factors exist. This leads to some degree of confidence in the predictions of the study.

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[^0]:    Figure 2.50 Fragment-Size Distribution at Various Ground Ranges
    for a Reinforced Concrete Structure

[^1]:    Note: $\quad x_{\text {max }}=$ maximum contour line distance

[^2]:    * Taken in atmosphere.

[^3]:    (2) Photographic observation of a series of tests involving underwater and bottom bursts of $300-1 \mathrm{~b}$ HE charges showed plumes with decidedly vertical trajectories, (Ref.. 31). This behavior is shown in Fig. 7. 15 (from Ref. 31) and 7.16 (from Ref. 20).
    (3) Full-scale nuclear test experience has shown that if the depth of the underwater burst is not too great, the bubble remains intact until it rises to the water surface. At this point debris is expelled into the steam and fission gases (Ref. 29). Such debris would have predominantly vertical trajectories as shown in Fig. 7. 16.

    The problem of far-flung debris from undervater bursts is probably a hazard to field troops in cases of shallow streams only -- where the initial gas bubble breaks the surface and provides ejecta entering the atmosphere at all angles and with high initial velocities. Under these conditions the stream-bed charge placed either on or somewhat beneath the bottom -- on the basis of energy partitioning -- can conceivably produce ejecta which enters the atmosphere at initial velocities greater than that from the comparable surface or shallow-burried burst on land.

    An analytic solution for initial velocities of the ejecta for either of the cases cated is beyond the scope of this portion of this investigation, and in fact, may be well beyond the scope of current knowledge. Experimentation to date can only contribute to rudimentary estimates of debris behavior from stream-bed charges.

    For explosions where the gas bubble does not initially break the surface or approach the surface, we can probably neglect maximum debris distance in determining a safe line for personnel. In these cases radioactive spray from the plume itself, the condensation cloud and the base surge may extend farther than the ejected bottom material.

[^4]:    FRAGMENTATION DATA ON EXPLODED DRY SANDSTONE BLOCKS

[^5]:    APPENDIX $F$
    CONCRETE FRAGMENT WEIGHTS AND DIMENSIONS
    FOR 1/24-SCALE SHIELDED REACTOR MODELS

[^6]:    * H. L. Brode and R. L. Bjork, "Craterıng from a Megaton Surface Burst', Paper L, Proceedings of the Geophysical Laboratory Lawrence Radiation Laboratory Cratering Symposium, UCRL-6438, University of California, October 1961.

[^7]:    * Optimum depth of burial is that depth which produces maximum volume of apparent crater. It results in ejecta having trajectories with pronounced vertical components.

